

Guide to the expression of uncertainty in measurement<br>
— Part 5: Examples<br>
Guide pour l'expression de l'incertitude de mesure — Partie 5:<br>
JCGM GUM-5:2024-12-06<br>
JCGM GUM-5:2024-12-06 Guide to the expression of uncertainty in measurement — Part 5: Examples

Guide pour l'expression de l'incertitude de mesure — Partie 5: Exemples

JCGM GUM-5:2024-12-06

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# **Contents**







# **Foreword**

<span id="page-5-0"></span>**Toreword**<br> **Streward**<br>
The 1997 a Joint Committ[ee](https://www.ilac.org/) for Giuldes in Metrology (JCGM), chaired by the Director of the<br>
nurear International des Poids et Metures (BIPM), was created by the seven international<br>
significations t In 1997 a Joint Committee for Guides in Metrology [\(JCGM\)](https://www.bipm.org/en/committees/jc/jcgm/), chaired by the Director of the Bureau International des Poids et Mesures (BIPM), was created by the seven international organizations that had originally in 1993 prepared the 'Guide to the expression of uncertainty in measurement' and the 'International vocabulary of basic and general terms in metrology'. The JCGM assumed responsibility for these two documents from the Technical Advisory Group 4 of the International Organization for Standardization (ISO/TAG4).

The Joint Committee is formed by the [BIPM](https://www.bipm.org/) with the International Electrotechnical Commission [\(IEC\)](https://www.iec.ch/), the International Federation of Clinical Chemistry and Laboratory Medicine [\(IFCC\)](https://www.ifcc.org/), the International Laboratory Accreditation Cooperation (ILAC), the International Organization for Standardization [\(ISO\)](https://www.iso.org/), the International Union of Pure and Applied Chemistry [\(IUPAC\)](https://www.iupac.org/), the International Union of Pure and Applied Physics [\(IUPAP\)](https://www.iupap.org/), and the International Organization of Legal Metrology (OIML).

JCGM has two Working Groups. Working Group 1, 'Expression of uncertainty in measurement', has the task to promote the use of the 'Guide to the expression of uncertainty in measurement' and to prepare documents for its broad application. Working Group 2, 'Working Group on International vocabulary of basic and general terms in metrology', has the task to revise and promote the use of the 'International vocabulary of basic and general terms in metrology' (the 'VIM').

In 2008 the JCGM made available a slightly revised version (mainly correcting minor errors) of the 'Guide to the expression of uncertainty in measurement', labelling the document 'JCGM 100:2008'.

In 2017 the JCGM rebranded the documents in its portfolio that have been produced by Working Group 1 or are to be developed by that Group: the whole suite of documents is now known as the 'Guide to the expression of uncertainty in measurement' or 'GUM', and is concerned with the evaluation and expression of measurement uncertainty, as well as its application in science, trade, health, safety and other societal activities.

This part of the suite contains a collection of examples illustrating the methods for uncertainty evaluation described in the GUM. An overview of the suite is given in Annex [A.](#page-122-0)

This document has been prepared by Working Group 1 of the JCGM, and has benefited from detailed reviews undertaken by member organizations of the JCGM and National Metrology Institutes.

# <span id="page-6-0"></span>**Introduction**

The GUM suite of documents (parts listed in Annex [A\)](#page-122-0) is concerned with the evaluation of uncertainty in measurement and how uncertainty is reported and used. Compact examples are given throughout the GUM suite that relate to specific areas of application and are intended to illustrate particular concepts. This document contains larger examples and case studies intended to illustrate the principles presented in the GUM suite for evaluating and reporting measurement uncertainty.

The examples are chosen to illustrate good practice in the evaluation of measurement uncertainty using the methods in the GUM, such as the GUF or the MCM or alternative methods not (yet) covered by the guidance in the suite. In some examples, methods for uncertainty evaluation are compared to highlight aspects of their suitability for the presented case.

**Introduction**<br>
The GIMM saide of decuments (parts listed in Annex A) is concerned with the evaluation of<br>
mechanism on increasing the committee of the transition of the committee in the committee of a policism of a singl The context and measurement in the examples are described to the extent relevant; simplifications can have been made in the interest of presenting a concise case. The data used in the examples is, where possible, taken from real measurements. These data are not necessarily the best, representing the state-of-the-art or in any other way to be interpreted as such. Measurement procedures and practices are presented in a concise format and not intended to enable reproducing the measurement. Where appropriate, references are provided to documents describing these procedures or practices.

Examples [1](#page-9-0) [\(Measurement of pH: linear interpolation\)](#page-9-0), 2 [\(Determination of](#page-21-0) benzo[a][pyrene\)](#page-21-0), [5](#page-39-0) [\(Greenhouse gas emission inventories\)](#page-39-0), 9 [\(Calibration of a torque](#page-65-0) [measuring system\)](#page-65-0), [10](#page-73-0) [\(Conformity assessment of total suspended particulate matter in](#page-73-0) [air\)](#page-73-0), [11](#page-80-0) [\(Effect of considering a 2D image as a set of pixels on a computed quantity\)](#page-80-0) and [12](#page-89-0) [\(Between-bottle homogeneity of reference materials\)](#page-89-0) were taken from a compendium of examples [[1](#page-126-1)].

Examples [6](#page-48-0) [\(Simple linear measurement models\)](#page-48-0), 7 [\(Calibration of weights: second-order](#page-54-0) [effects in uncertainty evaluation\)](#page-54-0), 8 [\(Gauge block calibration\)](#page-59-0) and [16](#page-109-0) [\(Comparison loss](#page-109-0) [in microwave power meter calibration\)](#page-109-0) were taken from JCGM 101:2008 [[12](#page-126-2)]. Example [13](#page-95-0) [\(Measurement of Celsius temperature using a resistance thermometer\)](#page-95-0) was taken from JCGM 102:2011 [7].

Examples [3](#page-27-0) [\(Relative molecular mass of glucose\)](#page-27-0) and [4](#page-30-0) [\(Gravimetric mixture preparation](#page-30-0) [and the calculation of composition\)](#page-30-0) were provided by A.M.H. van der Veen (VSL, The Netherlands). Example 14 [\(Activity of a radioactive source corrected for decay\)](#page-101-0) was provided by C. Michotte (BIPM). Example [15](#page-104-0) [\(Breaking force of steel wire rope\)](#page-104-0) was provided by S. Sidney (NLA, South Africa).

All examples were edited prior to inclusion in this document.

Results are generally reported in the manner described in JCGM 101:2008 [[12](#page-126-2)]. However, more than the recommended one or two significant decimal digits are often given here to facilitate comparison of the results obtained from the various approaches.

JCGM WG1 Committee draft

# **Guide to the expression of uncertainty in measurement — Part 5: Examples**

# <span id="page-8-0"></span>**0 Scope**



**0.1** This document gives a number of examples, which are worked out in considerable detail to illustrate the principles presented in the Guide to the expression of uncertainty in measurement (GUM) for evaluating and reporting uncertainty in measurement. Together with the (shorter) examples included in the other parts of the GUM, they should enable users of the GUM to put these principles into practice in their own work.

**Guide to the expression of uncertainty in**<br> **Committed Committee Committee Committee** which are worked on the considerable<br>
11 This decument gives a number of examples, which are worked on the considerable<br>
detail to ill **0.2** Because the examples are for illustrative purposes, they have by necessity been simplified (but not over-simplified). Moreover, because they and the numerical data used in them have been chosen mainly to demonstrate the principles of the GUM, neither they nor the data should necessarily be interpreted as describing real measurements, let alone the state of the art for a particular measurement. While the data are used as given, to prevent rounding errors, additional digits are retained in intermediate calculations than are usually shown. Thus the stated result of a calculation involving several quantities may differ slightly from the result implied by the numerical values given in the text or tables for these quantities.

- **0.3** The examples are used to demonstrate, among other:
	- Using knowledge about a quantity to obtain an estimate of the quantity and its associated uncertainty;
	- Propagation of uncertainty using the law of propagation of uncertainty (LPU) as in JCGM 100:2008 and JCGM 102:2011;
	- Propagation of distributions using analytical and numerical methods, including the Monte Carlo method (MCM), as in JCGM 101:2008 and JCGM 102:2011;
	- Applying law of propagation of uncertainty (LPU) and MCM for a univariate (scalar) and multivariate (vector) measurand;
	- Applying LPU and MCM for independent and dependent input quantities;
	- Using MCM to validate the results provided by LPU;
	- Determining a coverage interval or a coverage region;
	- Using measurement uncertainty in conformity assessment as in JCGM 106:2012;
	- Reporting measurement results.

**0.4** An overview of the parts of the GUM suite is given in Annex [A.](#page-122-0) The conventions and notation applied in this document are summarised in Annex [B.](#page-123-0) The symbols used throughout the document are explained in the text and in Annex [C.](#page-125-0)

# <span id="page-9-0"></span>**1 Measurement of pH: linear interpolation**

# <span id="page-9-1"></span>**1.1 Preamble**

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1.1 Prasmble<br>
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1.1 Prasmet teament of everyoint and multi-point interpolation of calibration<br>
4.1 A general uncertainties associated with the data propagated using t **1.1.1** A generic treatment of two-point and multi-point interpolation of calibration data is given with uncertainties associated with the data propagated using the law of propagation of uncertainty (LPU) and its generalization to vector measurands. The approach is applied to the measurement of hydrogen ion activity (pH). Such measurement is one of the most common in chemistry, although correlations associated with the input quantities in the measurement model are rarely taken into account. The treatment given follows common practice, which tends to give an optimistically small evaluation of the uncertainty associated with an estimated hydrogen ion activity (pH) value. A way of taking correlation into account in one typical instance is given but its implementation is problematical because of the difficulty in quantifying the correlation.

**1.1.2** The VIM concept of calibration [11, Definition 2.39] as constituting two stages is used. Here, the first stage involves fitting to measured data a function that describes the relationship of a response (dependent) variable *Y* to a stimulus (independent) variable *X*. The second stage involves using this relationship to determine the value *x* of a stimulus given the value *y* of the response. Uncertainties in both the stimulus and response variables are handled in both stages and propagated using the LPU in JCGM 100:2008 [[6](#page-126-5)] and its generalization to vector measurands in JCGM 102:2011 [7].

**1.1.3** Two scenarios are considered. One, a single party accesses the calibration data set and provides the required interpolated value. In doing so, the party may or may not determine the calibration parameters explicitly. Two, one party has access to the calibration data set, delivering the calibration parameters to a second party, which in turn provides the interpolated value.

**1.1.4** The approach is applied to the measurement of pH in which up to three two-point interpolations are required and uncertainties are tracked through the calculation.

## <span id="page-9-2"></span>**1.2 pH of a test solution**

**1.2.1** pH, the negative logarithm to base 10 of the activity of hydrogen ion in a solution, is arguably the most measured quantity in chemistry [[19](#page-127-0)]. The pH is a measure of the acidity of a solution. The pH of a solution is generally measured using a pH-sensitive glass electrode and an AgCl reference electrode. The potential difference of such an electrode system is proportional to pH and forms the basis of pH measurement.

**1.2.2** In 2002, the International Union of Pure and Applied Chemistry (IUPAC) issued a recommendation for revision of the pH scale based on the concept of a primary reference measurement procedure for pH [[15](#page-127-1)]. The use of an electrochemical (Harned) cell fulfils the criteria for a primary reference measurement procedure so that a pH value thus obtained is traceable to the International System of Units (SI), here the SI measurement unit 1.

**1.2.3** A solution, the pH of which is measured by such a cell at the highest metrological level, may be classified as a primary measurement standard and can be used to as-

**1.2.4** There are several approaches to pH measurement involving the use of 1-point, 2-point and multi-point calibration, least-squares regression, and with or without temperature correction. Here the 2-point calibration approach, with and without temperature correction, is used. Two topics are considered: (a) when the temperature of the test solution matches that of the standard (reference) solutions and (b) when this is not the case.

## <span id="page-10-0"></span>**1.3 Specification of the measurand**

rials (CRMs) to calibrate pH meters for routine use.

The measurand is the pH of a solution being calibrated. More generally, the measurand is the interpolated independent or dependent variable obtained from a relationship between those variables derived from data representing values of the variables. Intermediate measurands, when required, are the parameters describing the relationship.

### <span id="page-10-1"></span>**1.4 Measurement model**

sign pH values to coller solutions. Sinnalnd solutions are sold as certified reference inter-<br>tiga (CRMs) to calibrate pH meters for coutine use.<br>
1.2.4 There are several approaches to pH metasurement involving the use of There are two stages involved in calibration [11, definition 2.39]: (i) determine a calibration curve from calibration data and (ii) use that calibration curve. Because of the relative simplicity of two-point and multi-point interpolation as considered here, it may be preferable when circumstances permit to combine the stages into a single-stage model. Such a model avoids having to deal with intermediate correlation associated with the calibration curve parameters that are estimated in the first stage and used in the second. Operating in two stages corresponds to the use of a multi-stage model [10, Clause 8.4] and is necessary when the construction and use of the calibration model are carried out by different parties.

### <span id="page-10-2"></span>**1.4.1 Generic approach to two-point calibration**

**1.4.1.1** Two calibration points  $(X_1, Y_1)$  and  $(X_2, Y_2)$  are given that bracket  $X_0$ , an *X*-value for which *Y*<sup>0</sup> , the corresponding *Y* -value, is required under the assumption that the *Y* -value lies on the straight line joining the calibration points (see figure [1.1\)](#page-11-0).

**1.4.1.2** By similar triangles, with  $\delta X = X_2 - X_1$  and  $\delta Y = Y_2 - Y_1$ ,

<span id="page-10-3"></span>
$$
\frac{Y - Y_1}{X - X_1} = \frac{Y_2 - Y_1}{X_2 - X_1} = \frac{\delta Y}{\delta X}.
$$
\n(1.1)

A common representation of a straight-line calibration function, which is used here, is

$$
Y = a + bX, \tag{1.2}
$$

where *a* is the intercept on the *Y* -axis and *b* is the gradient [[61](#page-130-0)].

NOTE The form (1.2) is used in the straight-line calibration standard ISO/TS 28037 [[61](#page-130-0)] and will be familiar to many end users.

<span id="page-10-4"></span>**1.4.1.3** In the *single-stage model*, a single party has estimates of  $(X_1, Y_1)$ ,  $(X_2, Y_2)$  and  $X_0$ , and obtains an estimate of  $Y_0$ , the *Y* -value on the line corresponding to  $X_0$  (figure [1.1\)](#page-11-0). In doing so, the party may or may not determine *a* and *b* explicitly. The measurement model is specified by the description of the provision of  $Y_0$ .



<span id="page-11-2"></span><span id="page-11-1"></span><span id="page-11-0"></span>Figure 1.1: Two-point calibration using a linear interpolation function

**1.4.1.4** In the *two-stage model*, however, one party has estimates of  $(X_1, Y_1)$  and  $(X_2, Y_2)$ , and provides estimates of *a* and *b* (intermediate measurands) to the second party using the straightforwardly verified

$$
b = \frac{Y_2 - Y_1}{X_2 - X_1}.
$$
\n
$$
a = Y_1 - bX_1.
$$
\n(1.3)

The second party (possibly identical to the first party) has estimates of  $a$ ,  $b$  and  $X_0$ , and obtains the estimate of  $Y_0$  using the expression

$$
Y_0 = a + bX_0. \tag{1.5}
$$

The measurement model is again described by the process to provide  $Y_0$ . The resulting expression

<span id="page-11-3"></span>
$$
Y_0 = Y_1 + b(X_0 - X_1),\tag{1.6}
$$

and formula [\(1.3\)](#page-11-1) constitute the measurement model with  $Y_0$  as the measurand.

NOTE When  $X_1$  and  $X_2$  are far from the origin, that is,  $|X_2 - X_1| \ll \max(|X_1|, |X_2|)$ , an alternative form may be numerically more stable. One such form is given by working with a transformed *X*-variable

$$
\widetilde{X} = X - X_1.
$$

Using expressions (1.2) and [\(1.4\)](#page-11-2), the calibration function can then be expressed as

$$
Y = Y_1 + b\tilde{X},\tag{1.7}
$$

which is evaluated at the value  $X_0$  of the independent variable. The resulting expression

$$
Y_0 = Y_1 + b(X_0 - X_1) = Y_1 + \frac{Y_2 - Y_1}{X_2 - X_1}(X_0 - X_1),
$$

and formula (1.3) constitute the measurement model with  $Y_0$  as the measurand. It is accepted that such a transformation is not always appropriate.

The form of interpolation considered here is forward interpolation. Inverse interpolation, when the stimulus value  $X_0$  corresponding to a response value  $Y_0$  is required, can also be carried out (for treatments see [[61,](#page-130-0) [63](#page-130-1)]) but is not required here. The roles of *X* and *Y* can be interchanged when permitted by the context.

### <span id="page-12-0"></span>**1.4.2 Generic approach to multi-point calibration**

**1.4.2.1** Multi-point calibration is the treatment in clauses [1.4.1](#page-10-2) and 1.4.1.3 extended to an arbitrary number of points. In these clauses, a straight-line segment joining two of the calibration points serves as the calibration function. When there are *m* calibration points ( $m \geq 2$ ), with strictly increasing stimulus values, the points are joined pairwise by successive straight-line segments, the overall construction being a piecewise-linear function or first-degree spline [[26](#page-128-0)], acting as the calibration function. For each interval between pairs of successive points, the treatment of clauses 1.4.1 and 1.4.1.3 can be applied directly to the appropriate segment of the piecewise-linear function.

NOTE When  $m = 2$  the calibration function is a single straight-line segment is naturally monotonic, a necessary condition. For  $m > 2$ , the ordered points may not form a monotonic sequence, a situation not considered here [[63](#page-130-1)].

**1.4.2.2** Alternatively, straight-line fitting by least squares can be used taking reported uncertainties associated with the calibration data into consideration [61]. Polynomial interpolation or polynomial regression can also be used [63].

### <span id="page-12-1"></span>**1.4.3 Metrological extension**

The form of in[te](#page-10-4)rpolation consi[d](#page-10-4)ered here is forward interpolation, there is a properties applied to the spatial to the formula the traft interpolation of the spatial to the spatial to the spatial to the spatial to the sp The measurement model implied by two-point calibration is the algorithm to provide  $Y_0$ given *X*<sup>0</sup> . The data will generally have associated uncertainties arising from a Type A evaluation of uncertainty [[6,](#page-126-5) clause 4.2] especially following an analysis of repeated observations. Often there will also be uncertainties obtained from a Type B evaluation and associated covariances arising from common measurement effects [[6,](#page-126-5) clause 4.3]. Such covariances should be handled also to avoid producing invalid statements of uncertainty associated with predicted *Y* -values. The calibration data considered here are assumed to have independent errors.

### <span id="page-12-2"></span>**1.5 pH measured at a stipulated temperature**

**1.5.1** When the cell potential is measured at one of the stipulated temperatures on a reference material certificate for standard solutions, the following process is applicable. A correction approach is used to provide the pH of a test solution [[121](#page-135-0)] in which pH<sub>x</sub>, the pH of a test solution X, is given by using a cell twice to measure potential  $E_{\rm X}$  in X and potential  $E_{\mathcal{S}}$  in a standard solution S:

<span id="page-12-3"></span>
$$
pH_X = pH_S + \frac{E_X - E_S}{k}.
$$
\n(1.8)

In expression (1.8),  $\text{pH}_{\text{S}}$  is the pH of S, and

$$
k = \frac{RT \ln 10}{F},
$$

where *R* is the gas constant, *T* the temperature in kelvin and *F* the Faraday constant.

**1.5.2** Two other approaches are bracketing methods, which are generally more accurate and used here. Use is made of the reference material certificates for the standard solutions, which give pH values and associated standard uncertainties at stipulated temperatures.

It is assumed the temperature lies within the range (here, 5 ◦C to 50 ◦C) covered by the certificates for the Standard Reference Materials (SRMs) used here, namely, potassium dihydrogen phosphate (186-I-g) and disodium dihydrogen phosphate (186-II-g) [82]. Those certificates contain tables of pH values for temperatures from 5 ◦C to 50 ◦C at a spacing of  $5^{\circ}$ C.

<span id="page-13-4"></span><span id="page-13-3"></span>**1.5.3** The potential  $E_X$  of the test solution X is measured at temperature  $T_X$ , one of the stipulated temperatures in clause [1.5.2.](#page-13-1) Likewise, the potentials  $E_{\text{S}_1}$  and  $E_{\text{S}_2}$  are measured of two cells with standard solutions  $\mathsf{S}_1$  and  $\mathsf{S}_2$  such that the values of  $E_{\mathsf{S}_1}$  and  $E_{\mathsf{S}_2}$ bracket  $E_{\rm X}$  and are as near as possible to it [23]. The pH of S<sub>1</sub> and S<sub>2</sub> at temperature  $T_{\rm X}$ , namely,  $\rm pH_{\it S_{1},\it T_{X}}$  and  $\rm pH_{\it S_{2},\it T_{X}},$  are taken from reference material certificates.

<span id="page-13-1"></span>J.5.2 The order approaches are brocketing methods, which are generally more aperture<br>significant and the methods in the method in the method of the signification<br>of used bere, the is made of the neferance material certifi **1.5.4** Assuming its validity between the points  $(E_{S_1}, pH_{S_1,T_X})$  and  $(E_{S_2}, pH_{S_2,T_X})$ , linear interpolation is used to provide the pH value pH<sub>X</sub> corresponding to potential  $E_{\rm X}$ . The output quantity, the measurand, generically  $Y_0$ , is p $H_X$ , the pH of the test solution. The input quantities in the measurement model are  $E_{S_1},$   $E_{S_2},$   $\rm pH_{S_1,T_X},$   $\rm pH_{S_2,T_X}$  and  $E_X,$  corresponding respectively to  $X_1, X_2, Y_1, Y_2$  and  $X_0$  in the generic approach.

**1.5.5** In an extended measurement model [10, clause 9], account is taken of further influences. In this case, main effects are pH instrument calibration, instrument resolution and interpolation model. The uncertainties associated with the first two effects are provided by the instrument manual and inspection of the output display. Incorporating correction terms to account for these effects,

<span id="page-13-2"></span>
$$
E_{\text{X,corr}} = E_{\text{X}} + \delta E_{\text{res,X}} + \delta E_{\text{cal}},
$$
  
\n
$$
E_{\text{S1,corr}} = E_{\text{S1}} + \delta E_{\text{res,S}_1} + \delta E_{\text{cal}},
$$
  
\n
$$
E_{\text{S2,corr}} = E_{\text{S2}} + \delta E_{\text{res,S}_2} + \delta E_{\text{cal}}.
$$
  
\n(1.9)

Interpolated pH and correction quantities in expressions [\(1.9\)](#page-13-2) relating to cell potential are assumed independent. Clause 1.10.1 contains a discussion of the validation of results.

## <span id="page-13-0"></span>**1.6 pH measurement accounting for temperature**

When the cell potential is not measured at one of the temperatures specified on a reference material certificate for standard solutions, the procedure in clause [1.5](#page-12-2) can be extended as follows. A bracketing approach as in clause [1.5](#page-12-2) is used in which three linear interpolations are carried out to estimate pH for solution X and can be described by the following procedure, which constitutes an algorithmic description of the measurand model.

Measure the temperature  $T_X$  of the test solution X.

Determine the closest bracketing temperatures  $T_1$  and  $T_2$  in the SRM certificates such that  $T_X \in [T_1, T_2].$ 

Measure the potential  $E_{\rm X}$  in solution X and potentials  $E_{\rm S_1}$  and  $E_{\rm S_2}$  in standard solutions  ${\rm S_1}$ and  $S_2$ .

Obtain the certified pH value  $\text{pH}_{S_1,T_1}$  at temperature  $T_1$  for standard solution  $S_1$  from the certificate for S<sub>1</sub>. Likewise, obtain the certified pH value pH<sub>S2,  $T_1$ </sub> at temperature  $T_1$  for standard solution  $S_2$  from the certificate for  $S_2$ .

Repeat this process for temperature  $T_2$ , that is, obtain the pH value pH $_{\mathrm{S}_1,T_2}$  for solution  $\mathrm{S}_1,$ and the pH value pH $_{S_2,T_2}$  for solution  $S_2$ .

Apply three steps of linear interpolation:

- 1. Interpolate between the points  $(E_{S_1}, \text{ pH}_{S_1,T_1})$  and  $(E_{S_2}, \text{ pH}_{S_2,T_1})$  to give the pH value pH $_{\text{X}, T_1}$  at temperature  $T_1$  corresponding to potential  $E_{\text{X}}.$
- 2. Interpolate between the points  $\left(E_{S_1},\text{ pH}_{S_1,T_2}\right)$  and  $\left(E_{S_2},\text{ pH}_{S_2,T_2}\right)$  to give the pH value pH $_{\text{X},T_2}$  at temperature  $T_2$  corresponding to potential  $E_{\text{X}}.$
- 3. Interpolate between the points  $\left(T_1, \text{ pH}_{\text{X}, T_1}\right)$  and  $\left(T_2, \text{ pH}_{\text{X}, T_2}\right)$  so obtained to give the pH value pH $_{\text{X}, T_\text{X}}$  at temperature  $T_\text{X}$  corresponding to potential  $E_\text{X}.$

The generic treatment in clause [1.4.1](#page-10-2) is thus applied three times to implement these three stages of two-point interpolation.

The measurement model is specified by the above algorithmic description where the input quantities in the model are  $T_X$ ,  $E_X$ ,  $E_{S_1}$ ,  $E_{S_2}$ ,  $pH_{S_1,T_1}$ ,  $pH_{S_2,T_1}$ ,  $pH_{S_1,T_2}$  and  $pH_{S_2,T_2}$ .

#### <span id="page-14-0"></span>**1.7 Uncertainty propagation**

Obtain the certified pH value  $pE_{1,1}$  at imaginature  $T_1$  for standard dutting  $S_1$  for the certified pH value  $pE_{1,2}$ ,  $\pi$  imaginature  $T_1$  for certified obtion  $S_2$  from the certified of  $S_3$ . Here  $PE_{2,1}$  a **1.7.1** Uncertainty propagation in this clause is based on the assumption that the input quantities  $T_X$ ,  $E_X$ ,  $E_{S_1}$ ,  $E_{S_2}$ ,  $pH_{S_1,T_1}$ ,  $pH_{S_2,T_1}$ ,  $pH_{S_1,T_2}$  and  $pH_{S_2,T_2}$ , that is, the measured potentials and the pH values for the standard solutions, to the procedure in clause [1.6](#page-13-0) are independent. This assumption is consistent with IUPAC recommendations for pH measurement [[15](#page-127-1)]. There, procedures are given for accounting for input standard uncertainties based on the correlation-free variant of the law of propagation of uncertainty in [[6,](#page-126-5) clause 5.1]. In practice, input quantities are likely to be correlated and account should be taken of that fact. See the discussion in clause [1.10.3.](#page-18-1)

<span id="page-14-3"></span>**1.7.2** Consider first the *two-stage model*. For the first stage, the inputs are the calibration data *x*<sup>1</sup> , *y*<sup>1</sup> , *x*<sup>2</sup> and *y*<sup>2</sup> and their associated standard uncertainties. The outputs are the calibration parameters  $\hat{a}$  and *b* and their associated covariance matrix  $V_{[\hat{a}, \hat{b}]}$ . The model is bivariate (two output quantities): is bivariate (two output quantities):

$$
b = \frac{Y_2 - Y_1}{X_2 - X_1} = \frac{\delta Y}{\delta X}, \qquad a = Y_1 - bX_1.
$$
 (1.10)

For the uncertainty propagation, [[7,](#page-126-3) formula (3)] is applied to obtain the output covariance matrix for the model parameters *a* and *b*:

<span id="page-14-2"></span><span id="page-14-1"></span>
$$
\boldsymbol{V}_{[\hat{a},\hat{b}]} = \begin{bmatrix} u^2(\hat{a}) & u(\hat{a},\hat{b}) \\ u(\hat{a},\hat{b}) & u^2(\hat{b}) \end{bmatrix} = \boldsymbol{C}_1 \boldsymbol{V}_{in} \boldsymbol{C}_1^\top.
$$
\n(1.11)

In formula (1.11),  $u(\hat{a}, \hat{b})$  denotes the covariance associated with  $\hat{a}$  and  $\hat{b}$ ,  $V_{in}$  denotes the covariance matrix of the input parameter values, namely, the  $4 \times 4$  diagonal matrix with diagonal entries  $u^2(x_1)$ ,  $u^2(y_1)$ ,  $u^2(x_2)$  and  $u^2(y_2)$ , and

$$
\mathcal{C}_1 = \frac{1}{\delta x} \begin{bmatrix} -\hat{b}x_2 & x_2 & \hat{x}_1 & -x_1 \\ \hat{b} & -1 & -\hat{b} & 1 \end{bmatrix}
$$
 (1.12)

is the sensitivity matrix containing the first-order partial derivatives of  $X_1$ ,  $Y_1$ ,  $X_2$  and  $Y_2$ with respect to *a* and *b*, evaluated at their estimates  $x_1$ ,  $y_1$ ,  $x_2$ ,  $y_2$  and *b*.

**1.7.3** In the second stage, the inputs are the outputs from the first stage together with  $x_0$  and  $u(x_0)$ . The measurement model is

$$
Y_0 = a + bX_0.
$$

Applying uncertainty propagation [[7,](#page-126-3) Formula (3)] once more, the output standard uncertainty  $u(y_0)$  is given by

<span id="page-15-0"></span>
$$
u^2(y_0) = \mathbf{C}_2 \mathbf{V}_{[\hat{a}, \hat{b}]} \mathbf{C}_2^{\top} + \hat{b}^2 u^2(x_0),
$$
\n(1.13)

where  $C_2$  is the  $1 \times 2$  sensitivity matrix

<span id="page-15-1"></span>
$$
C_2 = \begin{bmatrix} 1 & x_0 \end{bmatrix}.
$$
 (1.14)

<span id="page-15-4"></span>**1.7.4** The *single-stage model* can be obtained by combining the two stages above. The substitution of formula [\(1.11\)](#page-14-1) into expression (1.13) yields

<span id="page-15-3"></span>
$$
u^2(y_0) = \mathbf{C}_2 \mathbf{C}_1 \mathbf{V}_{in} \mathbf{C}_1^\top \mathbf{C}_2^\top + \widehat{b}^2 u^2(x_0).
$$

Setting

$$
q = \frac{x_0 - x_1}{\delta x},\tag{1.15}
$$

the use of expressions [\(1.12\)](#page-14-2) and [\(1.14\)](#page-15-1) gives

$$
C_2C_1 = \frac{1}{\delta x} \left[ \begin{array}{cc} -\hat{b}(x_2 - x_0) & x_2 - x_0 & -\hat{b}(x_0 - x_1) & x_0 - x_1 \end{array} \right]
$$
  
=  $\left[ \begin{array}{cc} -\hat{b}(1 - q) & 1 - q & -\hat{b}q & q \end{array} \right].$ 

Hence, using

$$
V_{\text{in}} = \left[ \begin{array}{cc} u^2(x_1) & & \\ & u^2(y_1) & \\ & & u^2(x_2) \\ & & & u^2(y_2) \end{array} \right],
$$

$$
\mathbf{C}_2 \mathbf{C}_1 \mathbf{V}_{\text{in}} = \left[ -\widehat{b}(1-q)u^2(x_1) \quad (1-q)u^2(y_1) \quad -\widehat{b}qu^2(x_2) \quad qu^2(y_2) \quad \right].
$$

Thus,

s the sensitivity matrix containing the first-order partial derivatives of 
$$
X_1
$$
,  $Y_1$ ,  $X_2$  and  $Y_2$   
with respect to a and b, evaluated at their estimates  $x_1$ ,  $y_1$ ,  $x_2$ ,  $y_2$  and  $\hat{b}$ .  
1.7.3 In the second stage, the inputs are the outputs from the first stage together with  
 $v_0$  and  $u(x_0)$ . The measurement model is  
 $Y_0 = a + bX_0$ .  
Applying uncertainty propagation [7, Formula (3)] once more, the output standard uncer-  
ainry  $u(y_0)$  is given by  
 $u^2(y_0) = C_2V_{[2, \bar{b}]}\nC_2^T + \hat{b}^2u^2(x_0)$ ,  
where  $C_2$  is the 1 × 2 sensitivity matrix  
 $C_2 = \begin{bmatrix} 1 & x_0 \end{bmatrix}$ .  
1.7.4 The single-stage model can be obtained by combining the two stages above. The  
substitution of formula (1.11) into expression (1.13) yields  
 $u^2(y_0) = C_2C_1V_{in}C_1^TC_2^T + \hat{b}^2u^2(x_0)$ .  
1.7.4 The single-stage model can be obtained by combining the two stages above. The  
substitution of formula (1.11) into expression (1.13) yields  
 $u^2(y_0)$   
 $C_2C_1 = \frac{1}{\delta x} \begin{bmatrix} -\hat{b}(x_2 - x_0) & x_2 - x_0 & -\hat{b}(x_0 - x_1) & x_0 - x_1 \end{bmatrix}$   
 $= \begin{bmatrix} -\hat{b}(1-q) & 1-q & -\hat{b}q & q \end{bmatrix}$ .  
Hence, using  
 $V_{in} = \begin{bmatrix} u^2(x_1) & u^2(y_1) & u^2(x_2) & 0 \\ 0 & u^2(y_2) & u^2(y_2) \end{bmatrix}$ ,  
 $C_2C_1V_{in}C_1^T = \frac{1}{\delta x} \begin{bmatrix} -\hat{b}(1-q)u^2(x_1)(-\hat{b}x_2) + (1-q)u^2(y_1)x_2 - \hat{b}qu^2(x_2)\hat{b}x_1 \\ -qu^2(y_2)x_1 - \hat{b}^2(1-q)u^$ 

which is used in expression [\(1.13\)](#page-15-0) to yield, after simplification,

<span id="page-15-2"></span>
$$
u^{2}(y_{0}) = \hat{b}^{2}(1-q)^{2}u^{2}(x_{1}) + (1-q)^{2}u^{2}(y_{1}) + \hat{b}^{2}q^{2}u^{2}(x_{2}) + q^{2}u^{2}(y_{2}) + \hat{b}^{2}u^{2}(x_{0}).
$$
 (1.16)

The result (1.16) can also be confirmed from first principles.

### <span id="page-16-0"></span>**1.8 pH estimation at a stipulated temperature and associated uncertainty evaluation**

**1.8.1** The process in clauses [1.5.3](#page-13-3) and [1.5.4](#page-13-4) is followed. Consider the estimation of pH at a specific temperature. Values of potential in the test and standard solutions  $S_1$  and  $S_2$ are as follows:

$$
\widehat{E}_{\text{X}} = -1.875 \text{ mV},
$$
\n $\widehat{E}_{\text{S1}} = 6.15 \text{ mV},$ \n $\widehat{E}_{\text{S2}} = -26.35 \text{ mV},$ 

each of which was the average of 4 repeated observations. pH values at 25 °C for S<sub>1</sub> and S<sub>2</sub> from [[82](#page-132-0)] are

$$
\widehat{pH}_{S_1,25\degree C} = 6.8640, \qquad \widehat{pH}_{S_2,25\degree C} = 7.4157.
$$

Formulæ [\(1.3\)](#page-11-1) and [\(1.6\)](#page-11-3) yield  $\widehat{PH}_X = 7.0002$ . The standard uncertainties associated with the input quantities are given in [82] and are as follows:

$$
u(\widehat{E}_{\rm X}) = 0.0250 \,\text{mV}, \qquad u(\widehat{E}_{\rm S1}) = 0.0289 \,\text{mV}, \qquad u(\widehat{E}_{\rm S2}) = 0.0289 \,\text{mV},
$$

$$
u(\widehat{\rm PH}_{\rm S_1,25^{\circ}C}) = 0.0051, \qquad u(\widehat{\rm PH}_{\rm S_2,25^{\circ}C}) = 0.0051.
$$

The propagation of uncertainty carried out in accordance with expressions [\(1.15\)](#page-15-3) and  $(1.16)$  yields  $u(pH_X) = 0.0041$ .

1.8 **PH** estimation at a stipulated temperature and associated uncertainty<br>
1.8.1 The process in chause 1.5.3 and 1.5.4 is followed. [C](#page-13-2)onsider the esti[m](#page-13-2)ation of pit<br>
1.8.1 The process in chause 1.5.3 and 1.5.4 is followed. **1.8.2** These results relate to the basic measurement model for pH. The extended model would work with the corrected quantities in (1.9) rather than the uncorrected quantities. The estimates of all correction terms in the extended model are taken as zero.  $\delta E_{\text{cal}}$  appears in three of expressions [\(1.9\)](#page-13-2), so seemingly inducing correlation. However, this quantity is eliminated when the corrected quantities are used rather than the original. This effect can be seen mathematically by substituting  $E_{\rm X,corr}$ ,  $E_{\rm S1,corr}$  and  $E_{\rm S2,corr}$  as the 'new'  $X_0$ ,  $X_1$ and  $X_2$ , respectively, from expressions  $(1.9)$  into expressions  $(1.3)$  and  $(1.6)$ .

**1.8.3** The extended model would deliver the same estimate and standard uncertainty as the basic model to the number of digits reported.

### <span id="page-16-1"></span>**1.9 pH estimation accounting for temperature and associated uncertainty evaluation**

**1.9.1** The process in clause 1.6 is followed. The measured temperature of the test solution is  $\widehat{T}_X = 23.7$  °C. The bracketing pair  $T_1 = 25$  °C and  $T_2 = 20$  °C is selected.

Potential measurement gives

 $\widehat{E}_X = -1.875 \text{ mV}, \qquad \widehat{E}_{S_1} = 6.15 \text{ mV}, \qquad \widehat{E}_{S_2} = -26.35 \text{ mV},$ 

each of which was the average of 4 repeated observations. pH values at 25 ℃ and 20 ℃ for  $S_1$  and  $S_2$  from [82] are

$$
\widehat{\text{pH}}_{S_1,T_1}=6.8640, \quad \widehat{\text{pH}}_{S_2,T_1}=7.4157, \\ \widehat{\text{pH}}_{S_1,T_2}=6.8796, \quad \widehat{\text{pH}}_{S_2,T_2}=7.4323.
$$

Linear interpolation between the points

$$
(\widehat{E}_{S_1}, \widehat{\text{pH}}_{S_1, T_1}) \equiv (6.15 \,\text{mV}, \, 6.8640) \quad \text{and} \quad (\widehat{E}_{S_2}, \widehat{\text{pH}}_{S_2, T_1}) \equiv (-26.35 \,\text{mV}, \, 7.4157)
$$

gives the pH value  $\widehat{pH}_{X,T_1} = 6.8681$  at temperature  $\widehat{T}_1 = 25\text{ °C}$  corresponding to potential  $\widehat{E}_{\rm X} = -1.875$  mV.

Likewise, linear interpolation between

 $(\widehat{E}_{S_1}, \widehat{\text{pH}}_{S_1,T_2}) \equiv (6.15 \text{ mV}, 6.8796) \text{ and } (\widehat{E}_{S_2}, \widehat{\text{pH}}_{S_1,T_2}) \equiv (-26.35 \text{ mV}, 7.4323)$ 

gives the pH value  $\widehat{pH}_{X,T_2}$  = 7.4200 at temperature  $\widehat{T}_2 = 20$  °C corresponding to potential  $\widehat{E}_{\rm X} = -26.35$  mV.

Linear interpolation between the points  $(T_1, \text{ pH}_{X,T_1})$  and  $(T_2, \text{ pH}_{X,T_2})$  gives the pH value  $\widehat{\text{pH}}_X$  = 7.0109 at temperature  $\widehat{T}_X = 23.7 \text{ °C}$  corresponding to potential  $\widehat{E}_X = -1.875 \text{ mV}$ .

Associated standard uncertainties are



The above standard uncertainties associated with standard pH solutions are given in [[82](#page-132-0)].

The application of the method of clause 1.6 gives

 $\widehat{\text{pH}}_{\text{X}} = 7.0109$ ,

$$
u(\widehat{\text{pH}}_X) = 0.0041.
$$

## <span id="page-17-0"></span>**1.10 Interpretation of results**

### <span id="page-17-1"></span>**1.10.1 Validation of results**

gives the pH value  $\hat{p}H_{\chi} = 6.8681$  at temperature  $\hat{T}_1 = 25$  °C corresponding to parential<br>  $\hat{T}_2 = -1.875$  mV,<br>  $\hat{Z}_3 = -1.875$  mV,<br>  $\hat{Z}_4 = -1.875$  mV,<br>  $\hat{Z}_5 = -2.635$  mV,<br>  $\hat{Z}_6 = -2.635$  mV,<br>  $\hat{Z}_7 = -2.635$  mV, **1.10.1.1** A check of the numerical accuracy of linear interpolation was made for the example in clause [1.6.](#page-13-0) Since NIST certificate [[82](#page-132-0)] gives pH values for the standard solutions considered at temperature values from 5 ◦C to 50 ◦C in steps of 5 ◦C, cubic interpolation based on [[63](#page-130-1)] was carried out using temperature values 15 ◦C, 20 ◦C, 25 ◦C and 30 ◦C (two values immediately on either side of 23.7 °C), and the corresponding pH values for  $S_1$  given in the certificate. The interpolated value at 23.7 ◦C was 6.8677 compared with 6.8681 from linear interpolation. The magnitude of the difference, 0.0004, between these values is a factor of ten smaller than the standard uncertainty associated with the obtained pH value. A comparable result was obtained for  $S_2$  and for the other linear interpolations carried out. Thus, linear interpolation is considered adequate in this example.

**1.10.1.2** In a study by Damasco et al. [[36](#page-129-0)] it was reported that a Monte Carlo method (MCM) applied to primary pH measurement gave similar results to the LPU [[6](#page-126-5)]. The work of Wiora and Wiora [120] came to the same conclusion. To test these conclusions for the data used here, the MCM of JCGM 101:2008 [[12](#page-126-2)] was applied to the example in clause [1.9.](#page-16-1) The input quantities were modelled by normal distributions with means equal to the input estimates and standard deviations equal to the associated standard uncertainties. For 10<sup>6</sup> Monte Carlo trials, exactly the same result was delivered as in that clause to the number of decimal places stated.

## <span id="page-18-0"></span>**1.10.2 Rounding**

**1.10.2.1** The instrument display gave results in volts with 3 significant decimal places. Assume a rounding error in the last digit, that is, in the interval  $\pm 0.0005$  V. Characterizing resolution by a rectangular distribution over this interval, the consequent resolution stan-p dard uncertainty  $E_{\rm{res}}$  applying to all potential readings is  $0.0005$  V/  $\sqrt{3}$  = 0.000 29 V. This standard uncertainty is some one hundredth of the above potential standard uncertainties and so is negligible.

**1.10.2.2** The standard uncertainties associated with the interpolated value  $pH_X$  are scarcely influenced by the uncertainties associated with the pH values of the standard solutions. As an instance, if the latter standard uncertainties are replaced by zero in the ex-ample in clause [1.9,](#page-16-1)  $u(pH_X)$  becomes 0.0040 (originally 0.0041), implying that the further repeated observations of the three potentials would do much to reduce  $u(pH_X)$ , assuming the repeated observations are genuinely independent.

## <span id="page-18-1"></span>**1.10.3 Correlation**

**1.10.3.1** It is emphasized that the treatment given regards all input quantities as independent. Independence is a common assumption in general in pH uncertainty evaluation. This assumption is often made implicitly (see [43, 50, 71, 114], for instance), but has adverse consequences in that evaluated pH uncertainties can be optimistically small. To obtain more valid results, covariance effects should be be quantified and incorporated.

**1.10.3.2** Laboratories that follow IUPAC recommendations [15] will not take correlation into consideration and so might be reporting unreasonably small measurement uncertainties. There seems to be little relevant literature available on pH measurement on obtaining correlations associated with input quantities. If such correlations were available, they could be accounted for by applying the provisions in JCGM 100:2008 [[6,](#page-126-5) clause 5.2]. It is noted that correlation issues are discussed in [[119](#page-135-2)], but they relate to correlations induced by the choice of parametrization rather than being associated with input quantities.

1.10.2 Rounding<br>
1.10.2 Rounding<br>
1.10.2 The instru[m](#page-130-2)ent display gave results in volts with 3 significant desirnal aloss.<br>
Assume a rounding error in the last digit, that is, in the interval =0.0005V. Characterizing<br>
Assum **1.10.3.3** In terms of the pH certificates used here, it would appear that covariance between pH at different temperatures, and probably between pH values for different materials, can be deduced. The certificates give two uncertainties, one (as in clauses [1.8](#page-16-0) and [1.9\)](#page-16-1) for an SI-traceable value and one, much smaller, that omits the uncertainty associated with the Bates-Guggenheim convention [4]. The Bates-Guggenheim conventional uncertainty, as given on the certificates, could hence reasonably be taken as at least an approximate covariance when using the SI-traceable values. Doing so gives the correlation very close to unity, and indeed that (or any respectably high value) is suggested as a generally conservative treatment. See clause [1.10.3.6](#page-20-0) for a relevant discussion.

**1.10.3.4** To indicate the effect of ignoring correlation, an exercise is carried out in which full correlation is present between the pH values for the standards. Assume that all quantities are independent apart from these two pH values, which are accorded a correlation of unity. For this case of perfect correlation [[21,](#page-127-3) [29,](#page-128-1) [44](#page-129-2)], the standard uncertainties associated with these quantities must be identical, which indeed they are, being equal to 0.0051. In terms of the generic notation of clause [1.7.2,](#page-14-3) the input covariance matrix  $V_{in}$  is no longer diagonal but has covariance  $u^2(y_1) = u^2(y_2)$  in off-diagonal positions (2, 4) and (4, 2):

<span id="page-19-0"></span>
$$
V_{\text{in}} = \begin{bmatrix} u^2(x_1) & & u^2(y_1) & u^2(y_1) \\ & u^2(y_1) & u^2(y_1) & u^2(y_1) \end{bmatrix}.
$$
 (1.17)

Noting that  $u^2(y_1) = u^2(y_2) = u(y_1, y_2)$  in the fully correlated case, where  $u(y_1, y_2)$  is the covariance associated with  $y_1$  and  $y_2$ , by applying a similar treatment to that in clause 1.7.4  $\,$ but using the covariance matrix [\(1.17\)](#page-19-0) gives

<span id="page-19-1"></span>
$$
u^{2}(y_{0}) = \hat{b}^{2}(1-q)^{2}u^{2}(x_{1}) + (1-q)^{2}u^{2}(y_{1}) + 2q(1-q)u(y_{1}, y_{2}) + \hat{b}^{2}q^{2}u^{2}(x_{2}) + q^{2}u^{2}(y_{2}) + \hat{b}^{2}u^{2}(x_{0}).
$$
\n(1.18)

The only difference is that the terms

$$
(1-q)^{2}u^{2}(y_{1})+q^{2}u^{2}(y_{2})=[(1-q)^{2}+q^{2}]u^{2}(y_{1})
$$
\n(1.19)

in the uncorrelated treatment [expression (1.16)] are replaced by

<span id="page-19-2"></span>
$$
(1-q)^{2}u^{2}(y_{1})+2q(1-q)u(y_{1}, y_{2})+q^{2}u^{2}(y_{2})
$$
  
= [(1-q)^{2}+2q(1-q)+q^{2}]u^{2}(y\_{1}) (1.20)

in the fully correlated case [expression (1.18)]. Since expression (1.20) simplifies (exactly) to  $u^2(y_1)$ , expression [\(1.18\)](#page-19-1) becomes

$$
u^{2}(y_{0}) = \hat{b}^{2}(1-q)^{2}u^{2}(x_{1}) + u^{2}(y_{1}) + \hat{b}^{2}q^{2}u^{2}(x_{2}) + \hat{b}^{2}u^{2}(x_{0})
$$
\n(1.21)

in the fully correlated case. By applying a similar treatment to that in subclause [1.7](#page-14-0) on the general single-stage model but using the covariance matrix (1.17) gives

<span id="page-19-3"></span>
$$
u^{2}(y_{0}) = \hat{b}^{2}(1-q)^{2}u^{2}(x_{1}) + (1-q)^{2}u^{2}(y_{1}) + 2q(1-q)u(y_{1}, y_{2}) + \hat{b}^{2}q^{2}u^{2}(x_{2}) + q^{2}u^{2}(y_{2}) + \hat{b}^{2}u^{2}(x_{0}).
$$
\n(1.22)

Noting that  $u^2(y_1) = u^2(y_2) = u(y_1, y_2)$  in the fully correlated case, where  $u(y_1, y_2)$  is the covariance associated with  $y_1$  and  $y_2$ , expression [\(1.22\)](#page-19-3) becomes

<span id="page-19-4"></span>
$$
u^{2}(y_{0}) = \hat{b}^{2}(1-q)^{2}u^{2}(x_{1}) + u^{2}(y_{1}) + \hat{b}^{2}q^{2}u^{2}(x_{2}) + \hat{b}^{2}u^{2}(x_{0}).
$$
\n(1.23)

liagonal but has covariance  $u^2(y_1) = u^2(y_2)$  in off diagonal positions (2,4) and (4, 2)<br>  $V_{\text{in}} = \begin{bmatrix} u^2(x_1) & u^2(y_1) & u^2(y_2) \\ u^2(y_1) & u^2(y_1) \end{bmatrix}$ . (1.4)<br>
(oting that  $u^2(y_1) = u^2(y_2) = u(y_1, y_2)$  in the fully correlat **1.10.3.5** In the application of expression [\(1.23\)](#page-19-4) to the data in subclause [1.7,](#page-14-0) the standard uncertainty associated with  $pH_X = 7.0109$  becomes  $u(pH_X) = 0.0051$  (compared with 0.0041 when correlation between the pH standards is disregarded). Unsurprisingly, this value of  $u(pH_X)$  is the same as the (identical) values of the 'input' standard uncertainties *u*(pH<sub>S<sub>1</sub>,25 $\circ$ c) and *u*(pH<sub>S<sub>2</sub>,25 $\circ$ c). Thus there is no reduction in uncertainty in the fully</sub></sub> correlated case.

It is observed that the standard uncertainties  $u( {\rm pH}_{\rm S1,25^\circ C} )$  and  $u( {\rm pH}_{\rm S2,25^\circ C} )$  make comparatively large contributions compared with those for the measured potential values. The five standard uncertainty contributions [the square roots of the successive terms on the right side of expression [\(1.16\)](#page-15-2)] in the case where correlation is ignored are

 $-0.0004$ ,  $0.0001$ ,  $0.0038$ ,  $0.0013$ ,  $-0.0004$ 

to four decimal places. The values are to be compared with the corresponding four values from expression [\(1.23\)](#page-19-4), namely,

 $-0.0004$ ,  $0.0001$ ,  $0.0051$ ,  $-0.0004$ 

in the correlated case, confirming that perfectly correlated standard uncertainty components are combined additively (see [[6,](#page-126-5) clause 5.2.2, note 1]):

 $0.0038 + 0.0013 = 0.0051$ .

The situation is compounded in clause 1.9 where four (rather than two) pH values and two temperature values are involved.

<span id="page-20-0"></span>**1.10.3.6** Independence of the pH standards is not a valid general assumption. A treatment such as given in  $[29, \text{section } 4.1]$  $[29, \text{section } 4.1]$  $[29, \text{section } 4.1]$  is suggested, that is, to work with a common correlation coefficient  $\rho$  associated with the input pH values. The basic change would be that the off-diagonal terms of  $V_{\text{in}}$  in formula (1.17) would become  $\rho u^2(y_1)$ . (The case  $\rho = 0$ yields the uncorrelated case and  $\rho = 1$  the case of perfect correlation.) A value for  $\rho$ might be obtained on technical grounds by examining uncertainty budgets (to see the relative contribution from the Bates-Guggenheim convention, for instance) or some other means such as employing expert judgment.

**1.10.3.7** A further, chemical, issue is that the NIST standards are not solutions; they are solids that must be weighed, mixed and fully dissolved in high-purity water. Buffer solutions are not particularly sensitive to minor dilution problems, but the preparation just mentioned will add further variation. Atmospheric CO $_2$ , for example, can shift the pH values, especially in neutral and alkaline pH test samples ( $pH \ge 7$ ). Although a laboratory would exercise care in measuring secondary solutions, the analysis given here omits these important handling effects.

to four decisions (123), mannels, are to be compared with the corresponding famefalues<br>from respection (1.23), mannels,  $-0.0004$ ,<br> $-0.0004$ ,  $-0.0004$ ,  $-0.0004$ ,  $-0.0004$ ,<br>tin the correlated case, confirming that perfe

# <span id="page-21-0"></span>**2 Determination of benzo[a]pyrene**

# <span id="page-21-1"></span>**2.1 Preamble**

**2.1.1** The aim of this example is to show the uncertainty evaluation for determination of low levels of benzo[a]pyrene (BaP), an important polycyclic aromatic hydrocarbon (PAH) for ambient air monitoring. A comparison between the results obtained according to the JCGM 100:2008 uncertainty framework (GUF) [[6](#page-126-5)] and the Monte Carlo method (MCM) for the propagation of distributions [[12,](#page-126-2) [7](#page-126-3)] is made and discussed.

**2.1.2** Determination of low levels of PAHs is an important issue since they are ubiquitous toxic contaminants that can be present in all the environmental compartments even at trace levels. Among the various PAHs, BaP is classified as a carcinogenic agent and is listed in the current European legislation [[42](#page-129-3)] as a marker of the carcinogenic risk for the whole class of PAHs in ambient air.

**2.1.3** This example aims at comparing the results obtained by application of the GUF [[6](#page-126-5)] and the MCM for propagation of distributions [12] to real data sets derived from the quantification of a low mass of BaP spiked on filters commonly used for airborne particulate matter sampling.

**2.1.4** The description of the analytical method to quantify BaP in ambient air can be found in [[99](#page-133-0)], whereas details on the uncertainty evaluation, not explicitly reported in the present example, can be found in [[103](#page-133-1)].

# <span id="page-21-2"></span>**2.2 Specification of the measurand**

**2.** Determination of benzo[a]pyrene<br>
2.1 Preamble<br>
2.1 Preamble<br>
2.1 The sim of this example is to show the uncertainty evaluation for determinations<br>
of wherein of benzo[a]pyrene (BaP), an important polycyclic aromate h **2.2.1** A glass fibre filter (Pall & Whatman) having diameter of 47 mm, a type of filter commonly used for the sampling of airborne particulate matter, was spiked with the certified reference material (CRM) NIST SRM 2260a, containing 36 PAHs in an organic solution. The spiked filter was extracted by Soxhlet, following the extraction procedure described in [[99](#page-133-0)]. The same filter was subsequently extracted a second time, thus obtaining a diluted sample. The mass of BaP present in the two extracts was determined using a gas chromatograph coupled with a mass spectrometer (GC-MS).

NOTE Spiking is an analytical procedure by which a known mass (or amount of substance) of one or more compounds is added to serve as a reference point.

**2.2.2** The measurand in the present example is the mass  $m<sub>E</sub>$  of BaP contained in a nominal volume of 1 µL of the second extract.

## <span id="page-21-3"></span>**2.3 Measurement model**

**2.3.1** The quantification of the mass of BaP contained in 1 uL of the second extract was performed according to the Internal Standard method described in EN 15549 [[40](#page-129-4)]. An aliquot of the NIST SRM 2270, containing deuterated benzo[a]pyrene (BaP-d<sub>12</sub>), was added to the solution in order to obtain a concentration of BaP- $d_{12}$  equal to 0.2455  $\mu$ g mL<sup>-1</sup>, to be used as the internal standard. Then,  $1 \mu$ L aliquots of the solution were repeatedly (three times) injected in the GC-MS. The quotient of the peak area corresponding to the analyte and that corresponding to the internal standard was used to determine the mass  $m<sub>E</sub>$  of BaP present in the injected volume of the extracted sample, according to the model

<span id="page-22-2"></span>
$$
m_{\rm E} = f \frac{A_{\rm E}}{A_{\rm ISE}} m_{\rm ISE},\tag{2.1}
$$

where  $f$  is the GC-MS calibration factor,  $A_{\rm E}$  is the mean area (expressed in arbitrary units, a.u.) of the chromatographic peak corresponding to BaP in the extract, whereas  $m_{\text{ISE}}$ and  $A_{\text{ISE}}$  are the mass (ng) and the mean chromatographic area (a.u.) for the internal standard in the extract (ISE).

analyte and that corresponding to the internal standard was used to determine the mass<br>of the present in the injered volume of the extracted sample, according to the Toda<br>me of the present in the injered volume of the ext **2.3.2** The calibration factor *f* was obtained as the arithmetic mean of three calibration factors corresponding to three reference solutions at different BaP concentrations. Details of the calibration procedure were reported in [103]. In the evaluation of the uncertainty associated with *f* (characterized by 9 degrees of freedom), covariance terms between the three factors were taken into account: they were due to the same mass of the internal standard used in the calibration model for each of the factors and to the same CRM used for preparing the three necessary reference solutions. For the same reason,  $f$  and  $m_{\text{ISF}}$ , as input quantities of measurement model (2.1), were correlated because of the use of the same internal standard both in the calibration and in the analysis.

**2.3.3** The mass  $m_{\text{ISE}}$  and the associated uncertainty of the internal standard were obtained from its calibration certificate. The degrees of freedom was considered to be very large, so that it did not contribute to the effective degrees of freedom of the output uncertainty.

**2.3.4** Mean chromatographic peak areas  $A_{\text{E}}$  and  $A_{\text{ISE}}$  of BaP and BaP- $d_{12}$  were evaluated as the arithmetic means of the three repetitions of the area measurement of the relevant chromatographic peaks. Their standard uncertainties were calculated as the standard de-viations of those means [[6,](#page-126-5) clause 4.2.3] (hence, having two degrees of freedom). A strong linear relationship was observed between the areas of the BaP and those of the ISE in the same run; hence a corresponding covariance term for the two mean areas was calculated according to [[6,](#page-126-5) clause 5.2.3]. The estimates of the means, standard uncertainties, and covariances associated with the aforementioned input quantities are reported in table [2.1.](#page-22-3)

Quantity	Estimate	Std. uncertainty	Deg. of freedom	Covariance
	0.616	0.017		
$m_{\rm{ISE}}$	$0.2455$ ng	$0.0036$ ng	$\infty$	$-3.3 \times 10^{-5}$ ng
$A_{\rm E}$	85 114 a.u.	9564 a.u.		
$A_{\mathrm{ISE}}$	917546 a.u.	44 492 a.u.		$-2.03 \times 10^8$ (a.u.) <sup>2</sup>

<span id="page-22-3"></span>Table 2.1: Estimate, standard uncertainty, degrees of freedom and covariance of the input quantities in model (2.1)

### <span id="page-22-0"></span>**2.4 Uncertainty propagation**

## <span id="page-22-1"></span>**2.4.1 GUF**

**2.4.1.1** For calculating the uncertainty associated with the estimate  $\widehat{m}_{\rm E} = 0.014$  ng of the measurand, both the GUF [[6](#page-126-5)] and the MCM for the propagation of probability distributions

[7, 12] were applied and compared.				
some 25 % smaller.			Table 2.2: Uncertainty budget for the mass $m_{\rm E}$	<b>2.4.1.2</b> Applying the law of propagation of uncertainty (LPU) to model (2.1), taking into account the uncertainty and covariance contributions of the input quantities reported in table (2.1), the resulting standard uncertainty $u(\widehat{m}_{F})$ was 0.002 ng. Table 2.2 reports the uncertainty budget for the mass $\hat{m}_{\rm E} = 0.014$ ng of BaP of the sample obtained with the second extraction (the first two columns repeat part of the information already available in table 2.1). Specifically, in addition to the uncertainty contributions, the table gives the contributions to the uncertainty budget arising from the covariances. Accounting for covariance contributions gives $u^2(\widehat{m}_E) = 4.1$ ng <sup>2</sup> whereas disregarding them gives 3.1 ng <sup>2</sup> ,
Quantity $(x_i)$	$u(\widehat{x}_i)$		$c_i = \partial \widehat{m}_{\rm E} / \partial x_i$	$c_i^2 u^2(\hat{x}_i) / \text{ng}^2$
$\mathbf f$ $m_{\rm{ISE}}$ $A_{\rm E}$ $A_{\rm ISE}$	$1.7~\times 10^{-2}$ 3.6 $\times 10^{-3}$ 9.6 $\times 10^3$ 4.4 $\times 10^4$	ng a.u. a.u.	$2.28 \times 10^{-2}$ ng $5.71 \times 10^{-2}$ $1.65 \times 10^{-7}$ ng(a.u.) <sup>-1</sup> $-1.53 \times 10^{-8}$ ng(a.u.) <sup>-1</sup>	$1.50 \times 10^{-7}$ $4.23 \times 10^{-8}$ $2.50 \times 10^{-6}$ $4.53 \times 10^{-7}$
	$u(\widehat{x}_i,\widehat{x}_j)$		$c_i c_j$	$2c_i c_j u(\hat{x}_i, \hat{x}_j)/ng^2$
$f, m_{\text{ISE}}$ $A_{\text{E}}$ , $A_{\text{ISE}}$	$-3.3 \times 10^{-5}$ ng		$1.30 \times 10^{-3}$ ng $-2.03 \times 10^8$ (a.u.) <sup>2</sup> $-2.52 \times 10^{-15}$ ng <sup>2</sup> (a.u.) <sup>-2</sup> $1.02 \times 10^{-6}$	$-8.59 \times 10^{-8}$
$u^2(\widehat{m}_{\rm E})$				4.1 $\times 10^{-6}$
their corresponding degrees of freedom:	$v_{\text{eff}} = \frac{u^4(\widehat{m}_{\text{E}})}{\sum_i c_i^4 u^4(\widehat{x}_i)/v_i} \approx 5.2.$			<b>2.4.1.3</b> The effective degrees of freedom $v_{\text{eff}}$ for $u(\widehat{m}_E)$ was calculated using the Welch- Satterthwaite formula [6, eqn. (G2.b)], which was applied to the input uncertainties and (2.2)
tities is available [118]:	$v_{\text{eff},\rho} = \frac{u^4(\widehat{m}_{\text{E}})}{\sum_i \sum_i r_{ij}^2 c_i^2 c_i^2 u^2(\widehat{x}_i) u^2(\widehat{x}_j) / v_i}.$			This evaluation, however, neglects accounting for the covariances between the input quan- tities. A generalization of the Welch-Satterthwaite formula for use with correlated quan- (2.3)
are				The two non-trivial correlation coefficients between the estimates of the input quantities
			$r(\widehat{f}, \widehat{m}_{\text{ISE}}) = \frac{u(\widehat{f}, \widehat{m}_{\text{ISE}})}{u(\widehat{f})u(\widehat{m}_{\text{ISE}})} = \frac{-3.3 \times 10^{-5}}{1.7 \times 10^{-2} \times 3.6 \times 10^{-3}} = -0.54,$ $r(\widehat{A}_{\text{E}}, \widehat{A}_{\text{ISE}}) = \frac{u(\widehat{A}_{\text{E}}, \widehat{A}_{\text{ISE}})}{u(\widehat{A}_{\text{E}})u(\widehat{A}_{\text{ISE}})} = \frac{-2.03 \times 10^{8}}{9.6 \times 10^{$	whereas other correlation coefficients are either 0, as for $r(\hat{f}, \hat{A}_{\text{ISE}})$ , or 1, as for $r(\hat{f}, \hat{f})$ .
				The calculation of the effective degrees of freedom using equation (2.3) yields $v_{eff, \rho} \approx 4.8$ .

<span id="page-23-0"></span>Table 2.2: Uncertainty budget for the mass  $m_F$ 

$$
v_{\rm eff} = \frac{u^4(\widehat{m}_{\rm E})}{\sum_i c_i^4 u^4(\widehat{x}_i)/v_i} \approx 5.2.
$$
 (2.2)

<span id="page-23-1"></span>
$$
\nu_{\text{eff},\rho} = \frac{u^4(\widehat{m}_{\text{E}})}{\sum_{i} \sum_{j} r_{ij}^2 c_i^2 c_j^2 u^2(\widehat{x}_i) u^2(\widehat{x}_j) / v_i}.
$$
\n(2.3)

$$
r(\widehat{f}, \widehat{m}_{\text{ISE}}) = \frac{u(\widehat{f}, \widehat{m}_{\text{ISE}})}{u(\widehat{f})u(\widehat{m}_{\text{ISE}})} = \frac{-3.3 \times 10^{-5}}{1.7 \times 10^{-2} \times 3.6 \times 10^{-3}} = -0.54,
$$
  

$$
r(\widehat{A}_{\text{E}}, \widehat{A}_{\text{ISE}}) = \frac{u(\widehat{A}_{\text{E}}, \widehat{A}_{\text{ISE}})}{u(\widehat{A}_{\text{E}})u(\widehat{A}_{\text{ISE}})} = \frac{-2.03 \times 10^{8}}{9.6 \times 10^{3} \times 4.4 \times 10^{4}} = -0.48,
$$

**2.4.1.4** Coverage factors of a Student *t*-distribution with integer degrees of freedom are given in [[6,](#page-126-5) table G.2]. Coverage factors for non-integer values can be recovered by means of common statistical software.  $v_{\text{eff}} = 5.2$  (ignoring the correlations) and  $v_{\text{eff,}\rho} = 4.8$ (including the correlations) correspond to 95 % coverage factors  $k = 2.5$  and  $k = 2.6$ , respectively. Using the latter coverage factor, the expanded uncertainty  $U = k u(\widehat{m}_{\rm E})$  at a<br>05.04 equators probability is 0.0052.ps 95 % coverage probability is 0.0052 ng.

### <span id="page-24-0"></span>**2.4.2 MCM**

**2.4.2.1** The MCM for propagation of probability distributions of the input quantities was applied in order to approximate the distribution for the measurand  $m_{\rm E}$ , the mass of BaP in the extract. For this purpose, suitable probability distributions were assigned to the input quantities in model [\(2.1\)](#page-22-2), according to the criteria prescribed in [12, 7].

**2.4.2.2** Since the available information about  $f$  and  $m_{\text{ISE}}$  were estimates and their associated covariance matrix, a bivariate normal distribution was assigned to these quantities [[7,](#page-126-3) Clause 6.4.8]. Hence, the mean of this bivariate normal distribution is a vector  $[f, \widehat{m}_{\text{ISE}}]$ and its covariance matrix is

$$
V = \begin{bmatrix} u^2(\widehat{f}) & u(\widehat{f}, \widehat{m}_{\text{ISE}}) \\ u(\widehat{f}, \widehat{m}_{\text{ISE}}) & u^2(\widehat{m}_{\text{ISE}}) \end{bmatrix},
$$

whose components are available in table 2.1. Note that such a treatment is a simplification as it ignores the fact that  $u(f)$  is supported by a small degrees of freedom. One way to jointly model  $f$  using the  $t_{9}$ -distribution and  $m_{\rm ISE}$  using the normal distribution is with the use of copulas [[90](#page-133-2)].

<span id="page-24-1"></span>**2.4.2.3** The  $N = 2$  quantities  $A<sub>E</sub>$  and  $A<sub>ISE</sub>$  were estimated using  $n = 3$  repeated indications  $(A_{E_i}$  and  $A_{ISE_i}$ , respectively). Therefore, according to [7, Clause 5.3.2], a scaled and shifted bivariate *t*-distribution with one degree of freedom (*ν* = *n* − *N*) was assigned to them. The expectation of the distribution is the vector  $[\widehat{A}_{E}, \widehat{A}_{ISE}]^T$  and its scale matrix  $S/n$  is

$$
\pmb{S} = \frac{1}{\nu} \Bigg[ \begin{matrix} \sum_{i=1}^3 (\widehat{A}_{E_i} - \widehat{A}_{E})^2 & \sum_{i=1}^3 (\widehat{A}_{E_i} - \widehat{A}_{E}) (\widehat{A}_{ISE_i} - \widehat{A}_{ISE}) \\ \sum_{i=1}^3 (\widehat{A}_{E_i} - \widehat{A}_{E}) (\widehat{A}_{ISE_i} - \widehat{A}_{ISE}) & \sum_{i=1}^3 (\widehat{A}_{ISE_i} - \widehat{A}_{ISE})^2 \end{matrix} \Bigg],
$$

whose components are  $S_{1,1} = 548\,861\,202(a.u.)^2$ ,  $S_{1,2} = S_{2,1} = -1\,220\,621\,757(a.u.)^2$ , and  $S_{2,2} = 11\,877\,338\,582$  (a.u.)<sup>2</sup>. Note that for  $\nu = 1$ , the mean and the covariance matrix of the *t*-distribution are not defined. In contrast, a coverage region for the distribution can always be determined [7, Clause 5.5.2, Note 1].

2.4.14 Coverage fieldo[m](#page-22-3) of a Student *i*-distribution with integral express of flexible and the content for the product of a Student *Committ[ee](#page-126-2) degrees* of ferminal method in the committee degree of  $\alpha$  formula and  $v_{x,$ **2.4.2.4** The propagation of the probability distributions for the input quantities through measurement model [\(2.1\)](#page-22-2) was implemented in the R environment [[95](#page-133-3)] by applying the R functions rmvnorm and rmvt available in the "mvtnorm" package [[49](#page-130-3)]. For each input quantity,  $M = 10<sup>7</sup>$  values were drawn. Since only positive values of the measurand are feasible, the joint input probability density functions were numerically truncated at zero by disregarding negative values drawn during the MCM simulation [[12,](#page-126-2) Clause 9.4.2.1.1, Note], thus obtaining corresponding BaP mass values smaller in number than *M*. The number of values retained, however, was about 9  $\times$  10<sup>6</sup>, sufficient to provide a reliable numerical approximation for the measurand distribution. From the MCM distribution, the shortest 95 % coverage interval was obtained and reported in table [2.3.](#page-25-1)

NOTE The rejection of unphysical MCM values for the input distributions is just an approximate and somewhat crude procedure to take into account the constraint  $m<sub>E</sub> > 0$ . A rigorous treatment using Bayesian inference [[113](#page-134-1)] would naturally include the constraint as prior knowledge about the measurand. Also see [[12,](#page-126-2) Scope, NOTE 2].

### <span id="page-25-0"></span>**2.5 Reporting the result**

**2.5.1** Figure [2.1](#page-25-2) shows the results for  $m_E$  determined using GUF (bell-shaped curve) and MCM (frequency distribution shown as a histogram). The endpoints of the coverage intervals are also reported in table [2.3.](#page-25-1)



<span id="page-25-2"></span>Figure 2.1: Results for  $m_E$  determined using GUF (bell-shaped curve) and MCM (frequency distribution shown as a histogram).

Table 2.3: Estimate, associated standard uncertainty and 95 % coverage interval for  $m_E$  as provided by GUF; median and shortest 95 % coverage interval from MCM

<span id="page-25-1"></span>

Method	$\widehat{m}_{\rm E}/\mathrm{ng}$	Median/ng $u(\widehat{m}_{\rm F})/ng$		$95\%$ C.I./ng
<b>GUF</b>	0.014	__	0.002	[0.009, 0.019]
<b>MCM</b>	__	0.014		$\lceil \approx 0, 0.032 \rceil$

**2.5.2** In table 2.3, neither the mean nor the standard deviation of the MCM distribution are given, as they are unreliable. This is due to the presence in the input distributions of a bivariate for which these moments are not defined (see [2.4.2.3](#page-24-1) and [2.6.1\)](#page-26-1). By contrast, the median and the endpoints of coverage intervals, being quantiles (which always exist), are meaningful. It is worth noting that the MCM median (0.014 ng) is close to the GUF estimate.

## <span id="page-26-1"></span><span id="page-26-0"></span>**2.6 Interpretation of results**

**2.6** Interpretation of results<br>**2.6** Interpretation of results<br>**2.6.1** When applying MCM, the measurand estimate and the associated uncertainty at<br>easily taken as the mean and the standard deviation of the simulated our<br> **2.6.1** When applying MCM, the measurand estimate and the associated uncertainty are usually taken as the mean and the standard deviation of the simulated output results, according to [[12,](#page-126-2) eqns. (16) and (17)]. Nonetheless, [[12,](#page-126-2) Clause 6, Note 2] states that in some special circumstances, such as when one of the input quantities has been assigned a Student *t*-distribution with fewer than three degrees of freedom, the expectation and standard deviation of the output quantity might not exist and the above-cited equations (16) and (17) in JCGM 101:2008 might not then provide meaningful results. A coverage interval for the measurand can, however, be formed, since the simulated output distribution is meaningful. This is exactly the situation of the present example, for which, in fact, plausible estimates and corresponding standard uncertainties are those obtained within the GUF, as reported in table [2.3,](#page-25-1) whereas feasible coverage intervals are those provided by MCM, as discussed shortly.

**2.6.2** From figure [2.1,](#page-25-2) it is evident that the two approaches give quite different results in terms of coverage intervals. Although the assumed output distribution in the GUF is a Student *t*-distribution with few degrees of freedom, hence leading to a large coverage factor for the calculation of the corresponding expanded uncertainty, the MCM coverage interval is 3.2 times larger than that obtained in the GUF. Moreover, it appears asymmetric with respect to the measurand estimate because of the left-censoring of the values generated because of the physical constraint. This is a clear example of those situations in which the Central Limit Theorem does not hold, since the probability density function (PDF) for the output quantity is neither a normal distribution nor a scaled and shifted Student *t*distribution.

**2.6.3** The Eurachem/CITAC guidelines on uncertainty evaluation for analytical measurements [[39](#page-129-5)] recommend truncating the coverage interval at zero whenever a negative lower limit is found for a non-negative quantity.

**2.6.4** Summarizing, this example is a clear case in which blind adherence either to the approach in JCGM 100:2008 [6] or to the MCM in [[12,](#page-126-2) [7](#page-126-3)] would be dangerous. Careful considerations of estimates, standard uncertainties and coverage intervals are always needed according to the specific problem under study.

# <span id="page-27-0"></span>**3 Relative molecular mass of glucose**

## <span id="page-27-1"></span>**3.1 Preamble**

**3.1.1** This example demonstrates the uncertainty evaluation in a calculation using reference data, in this instance standard atomic weights. The case shows how to model ignorance about the isotopic composition of the elements in the material and how to propagate the measurement uncertainty using the law of propagation of uncertainty (LPU).

**3 Relative [m](#page-134-2)olecular mass of glucose**<br> **3.[1](#page-133-4) Preamble**<br> **3.1 This example demonstrates the uncertainty evaluation in a calculation using reference data, in this increase standard atomic weights. The case shows how to infi 3.1.2** Relative molecular masses of substances, also known as "molecular weights", are frequently used in science, for example, to study the thermodynamics of the conversion of glucose. Relative molar masses can be calculated from the standard atomic weights of the elements present in the molecule, appreciating the coefficients describing the atomic composition. For glucose, whose chemical formula is  $\rm C_6H_{12}O_6$ , these coefficients take the values 6 for carbon, 12 for hydrogen, and 6 for oxygen. Relative atomic and molecular masses are functions of the isotopic composition of the elements present in the material(s) at hand [[14](#page-127-4)]. The standard atomic weights apply to "normal materials", that is, materials that have not been artificially altered in isotopic composition and are a relevant source of the element. Such materials include laboratory reagents [117].

**3.1.3** In this example, the isotopic composition of the elements in the material is unavailable. The calculation of the relative molecular mass of glucose will therefore be based on the assumption that the isotopic compositions of the elements in the molecule lie somewhere in the interval of values of isotopic compositions for normal materials. For elements with documented substantial variation in isotopic composition, the standard atomic weights are given in the form of an interval (see table 3.1).

<span id="page-27-3"></span>Table 3.1: Lower (*a*) and upper (*b*) bounds for the standard atomic weight intervals of hydrogen, carbon and oxygen [93]

<span id="page-27-4"></span>

### <span id="page-27-2"></span>**3.2 Measurement model**

The relative molecular mass of glucose *M* is given by

$$
M = \sum_{i=1}^{3} \nu_i A_i,
$$
 (3.1)

where

*M* denotes the relative molecular mass (molecular weight),

*νi* the atomic composition coefficient of element *i* in the molecule,

*Ai* the standard atomic weight (relative atomic mass) of element *i*.

### <span id="page-28-0"></span>**3.3 Uncertainty evaluation**

**3.3.1** The measurement model [\(3.1\)](#page-27-4) is linear because the atomic composition coefficients are known exactly. In this case, the LPU can be applied to give an exact expression for the squared standard uncertainty associated with  $\widehat{M}$  as follows:

$$
u^{2}(\widehat{M}) = \sum_{i=1}^{3} \nu_{i}^{2} u^{2}(\widehat{A}_{i}).
$$
\n(3.2)

**3.3.2** For the above calculations, an estimate of the relative atomic mass of each element is required with an associated standard uncertainty. The measurement uncertainty should account for at least (a) the knowledge about the isotopic composition of the materials considered, and (b) the uncertainty associated with the measurement of the masses of the isotopes by mass spectroscopy.

3.3. Uncertrainty evaluation<br>
3.3. The measurement model (3.1) is linear because the atomic composition configuration<br>
or a boson exactly, in this case, the LPU can be applied to give an exact expression for the<br>
squared **3.3.3** For each element, both the estimate and the standard uncertainty can be obtained from a Type B evaluation of standard uncertainty using the atomic-weight intervals given in table [3.1.](#page-27-3) As the isotopic composition with respect to the elements H, C, and O can lie anywhere in the respective atomic-weight intervals, i.e., there is no specific information about the values within this interval. In this case, a rectangular (uniform) distribution is assigned to each input quantity *A<sup>i</sup>* [6, clause 4.3.7]. For this probability distribution, the estimate of  $A_i$  and the standard uncertainty associated with this estimate are

$$
\widehat{A}_i = \frac{a_i + b_i}{2}, \qquad u(\widehat{A}_i) = \frac{b_i - a_i}{2\sqrt{3}}.
$$

The resulting estimates and standard uncertainties associated with these input quantities are given in table [3.2.](#page-28-2)

Table 3.2: Estimates and standard uncertainties for the standard atomic weights of the elements found in glucose

<span id="page-28-2"></span>

	Element	$A_i$	$v_i$	$u(A_i)$	$v_i u(\hat{A}_i)$
	hydrogen	1.00798	12	0.00008	0.00094
2	carbon	12.0106	h	0.00058	0.00346
3	oxygen	15.9994	n	0.00021	0.00128

The estimate of the relative molecular mass of glucose is calculated as

$$
\widehat{M} = 6 \times 12.0106 + 12 \times 1.00798 + 6 \times 15.9994 = 180.1557.
$$

The associated standard uncertainty is calculated from the contributions  $v_i u(A_i)$  given in the last column of table [3.2.](#page-28-2) The standard uncertainty is

$$
u(\widehat{M}) = (0.00094^2 + 0.00346^2 + 0.00128^2)^{1/2} = 0.0038.
$$

### <span id="page-28-1"></span>**3.4 Discussion**

**3.4.1** This example highlights how an estimate and an associated standard uncertainty can be evaluated from an atomic-weight interval. The probability density function chosen depends on the assumptions made, and these usually depend on the information at hand.

<span id="page-29-0"></span>Estimates and standard uncertainties thus derived for the three elements in question are independent of the isotopic composition and valid for any material that conforms with the description of a *normal material* of International Union of Pure and Applied Chemistry (IUPAC) [[22,](#page-127-5) [117](#page-134-2)].

Si[m](#page-132-1)ilanes and standard unnear<br>interior and standard method in the develop formula control and we detect of the the conforming term<br>dependent of the isotopic composition and will do the avery metric in the description of t **3.4.2** The use of the rectangular distribution underlines the lack of detailed information concerning the isotopic composition [[89,](#page-132-1) [110](#page-134-3)] and is in agreement with earlier reports on the standard atomic weights [[115,](#page-134-4) [116](#page-134-5)]. By no means can this choice of probability distribution be understood as depicting the actual isotopic composition. Rather, it depicts the state of knowledge about the composition, which is in this example very modest. The efforts of IUPAC in investigating the isotopic composition of elements in *normal materials* focus very much on establishing the interval end points, rather than modelling the isotopic composition of elements across materials [22]. The use of the rectangular distribution in connection with the standard atomic weights is not always sufficiently accurate to ignore completely the effects of the uncertainty associated with the relative molecular masses in a further application, for example, the calculation of the composition of a mixture expressed in amount-of-substance fractions [[109](#page-134-6)].

**3.4.3** Had the isotopic composition of the glucose been available, a more accurate value for the relative molecular mass *M* could have been obtained [89]. The estimate and standard uncertainty obtained in this example are for this application and many others in chemistry fully satisfactory and sufficiently accurate for further calculations [[110](#page-134-3)]. With increasing accuracy of measurements, there will be a growing number of applications for which the approach to calculating relative molecular masses as outlined in this example is inadequate. In such cases, a batch-specific mass-spectrometric determination of the isotopic composition of the given sample used in the application would be required to obtain information about the isotopic composition of the elements in the material.

**3.4.4** Calculations involving the standard atomic weights should be performed using the latest values published for these. At the time of writing, the standard atomic weights from 2021 [[93](#page-133-4)] were the current.

# <span id="page-30-0"></span>**4 Gravimetric mixture preparation and the calculation of composition**

## <span id="page-30-1"></span>**4.1 Preamble**

**4.1.1** This example demonstrates the uncertainty evaluation of the composition of a calibration gas mixture prepared by static gravimetry. It illustrates how data from weighing, purity analysis and molar masses are combined, and how correlations between quantities are duly taken into account. The propagation of uncertainty is carried out using the multivariate law of propagation of uncertainty (LPU) from JCGM 102:2011 [[7](#page-126-3)].

4 Gravitmetric mixture preparation and the calculation of composition<br>1961 Domain (2011)<br>21. The example demonstrates the uncertainty evaluation of the composition of a calculation<br>11. This example demonstrates the uncert **4.1.2** Fundamental to the metrology of chemical composition is the preparation of mixtures having well-known composition from pure substances. These mixtures are often used as reference materials for calibration. There are two principal ways for preparing such mixtures: gravimetric and volumetric. In this example, the gravimetric preparation of a gas mixture is considered. The mixture is prepared from two high-purity gases: carbon dioxide and nitrogen. These gases have been checked for purity, so a correction can be applied for the contents of the impurities present in the gases. The objective is to calculate the amount fractions of the abundant components and their associated standard uncertainties and covariances.

## <span id="page-30-2"></span>**4.2 Measurement model**

**4.2.1** The composition of a gravimetrically prepared gas mixture is commonly expressed in amount fractions, which are calculated as the ratio of the amount of component  $k$   $(n_k)$ in the gas mixture by the sum of the amounts of all components in the mixture. In a slightly amended notation, the measurement model of ISO 6142-1 [64] for the amount fraction of a component  $k$  ( $k = 1 \ldots q$ ) in a gravimetric mixture of  $p$  parent gases is

<span id="page-30-3"></span>
$$
y_k = \frac{n_k}{\sum_{k=1}^q n_k}, \quad \text{where} \quad n_k = \sum_{j=1}^p \frac{m_j x_{kj}}{M_{p,j}}.
$$
 (4.1)

where

 $y_k$  denotes the amount fraction of component  $k$  in the mixture,

*nk* the amount of component *k*,

*n* the (total) amount of substance in the mixture,

*mj* the mass of parent gas *j* added to the mixture,

 $x_{ki}$  the amount fraction of component *k* in parent gas *j*,

*M*P,*<sup>j</sup>* the molar mass of parent gas *j*

*p* the number of parents,

*q* the number of components in the mixture.

Expression (4.1) is an example of a multivariate measurement model. The the molar mass of parent gas *j*,  $M_{P,j}$  is determined by the chemical composition and the molar masses of

its constituent components,  $M_k$   $(k = 1 \dots q)$ :

$$
M_{P,j} = \sum_{k=1}^{q} x_{kj} M_k.
$$
 (4.2)

where

 $M_{P,j}$  denotes the molar mass of parent gas *j* 

 $x_{ki}$  the amount fraction of component *k* in parent gas *j*,

*M<sup>k</sup>* the molar mass of component *k*,

*q* the number of components in the mixture.

In the current example, there are two parent gases ( $p = 2$ ) comprising eight components  $(q = 8)$  comprised of  $e = 5$  different chemical elements.

**4.2.2** The weighing is performed using the substitution method [78, [109](#page-134-6)], where the sample and reference cylinder are weighed alternately a number of times. The weighing data are given in table [4.1.](#page-31-0) The same reference cylinder is used throughout the preparation process. Its mass is not needed, as it cancels in the expressions for the masses of the parent gases transferred into the sample cylinder.

<span id="page-31-0"></span>Table 4.1: Readings of the balance for the sample and reference cylinders prior to filling them  $(r_0)$ , after adding parent gas  $\#1$   $(r_1)$ , and after adding parent gas  $\#2$   $(r_2)$ 



is constituent components,  $M_2$  ( $k = 1...q$ ):<br>  $M_{2y}$  decores the molar mass of parent gas j<br>  $M_{2y}$  decores the molar mass of parent gas j<br>  $M_{2y}$  decores the molar mass of parent gas j<br>  $M_3$  the amount fraction of co **4.2.3** The chemical composition of the two parent gases is given in table [4.2.](#page-32-3) The amount fraction of the most abundant component is calculated by difference, that is, by subtracting the sum of the amount fractions of the other components from unity [[109,](#page-134-6) [60](#page-130-4)]. The subtraction gives rise to non-zero covariances between this amount fraction and the amount fractions of the minor components. It is noted that the standard uncertainty associated with the amount fraction nitrogen in the nitrogen gas is only marginally larger than the standard uncertainty associated with the amount fraction argon (the dominant uncertainty component).

		Carbon dioxide gas		Nitrogen gas		
	Component	$\widehat{x}$	$u(\widehat{x})$	$\widehat{x}$	$u(\widehat{x})$	
	Ar			5.0	3.0 <sub>•</sub>	
	CH <sub>4</sub>	0.10	0.06	0.008	0.005	
	CO	0.5	0.3	0.015	0.009	
	CO <sub>2</sub>	999988.0	3.4	0.010	0.006	
	H <sub>2</sub>			0.025	0.015	
	H <sub>2</sub> O	1.9 4.3	0.1 2.1	0.010 999 994.9	0.006 3.0	
	$\mathrm{N}_2$ ${\rm O}_2$	5.3	2.6	0.005	0.003	
4.3	<b>Uncertainty evaluation</b>					
4.3.1	<b>Correlations</b>					
	are done sequentially (see table 4.1) and the estimates of both parent gas masses, $m(CO_2) = r_1 - r_0$ and $m(N_2) = r_2 - r_1$ ,					
	are affected by the same $r_1$ . Last, the estimates of the impurity amount fractions in the parent gases are correlated since the amount fraction of the most abundant component is estimated by mass balance.					
	4.3.2 General considerations					
	The multivariate LPU in JCGM 102:2011 uses vectors and matrices. For the mea- surement model (4.1) of ISO 6142-1, the vector of input quantities can be considered to be partitioned into several input vectors, namely one containing the masses of the trans- ferred parent gases $m$ , another containing the molecular weights $M$ , and a third with the compositions of the parent gases $x_1, \ldots, x_p$ .					
4.3.2.1	4.3.2.2 Thus, all the input quantities can be represented by a vector $\alpha$ formed by "stack- ing" the vectors $m, M, x_1, \ldots, x_p$ :					

<span id="page-32-3"></span>Table 4.2: Chemical composition of the two parent gases expressed as amount fractions of their components in  $\mu$ mol mol<sup>-1</sup>

### <span id="page-32-0"></span>**4.3 Uncertainty evaluation**

### <span id="page-32-1"></span>**4.3.1 Correlations**

$$
m(CO_2) = r_1 - r_0
$$
 and  $m(N_2) = r_2 - r_1$ ,

### <span id="page-32-2"></span>**4.3.2 General considerations**

$$
a = \begin{bmatrix} m \\ M \\ x \end{bmatrix}, \quad x = \begin{bmatrix} x_1 \\ \vdots \\ x_p \end{bmatrix}.
$$
 (4.3)

**4.3.2.3** If there are no dependencies between the vectors  $m, M$ , and  $x$ , then  $V_a$  can be written in block matrix notation as follows:

$$
V_{\widehat{\alpha}} = \text{diag}\{V_{\widehat{m}}, V_{\widehat{M}}, V_{\widehat{x}}\}.
$$
\n
$$
(4.4)
$$

L3.2.3 If there are no dependents between the vectors  $m$ ,  $M$ , and  $x$ , then  $V_{\mu\nu}$ complementation of the covariance matrices  $V_{Rt}$ ,  $V_{\overline{S}}$ , and  $V_{\mu}$  is discussed briefly in  $V_{R} = \text{diag}(V_{Rt}, V_{\overline{S}t}, V_{\overline{S}}$ The construction of the covariance matrices  $V_{\hat{m}}$ ,  $V_{\hat{M}}$ , and  $V_{\hat{x}}$  is discussed briefly in Clauses 4.2.2.4.2.4. and 4.2.5, geometrically The mathed of construction is boost on the Clauses [4.3.3,](#page-33-0) [4.3.4,](#page-34-0) and [4.3.5,](#page-35-0) respectively. The method of construction is based on the LPU from JCGM 102:2011 [[7](#page-126-3)] considering multivariate sub-models.

### <span id="page-33-0"></span>**4.3.3 Weighing data**

**4.3.3.1** The mass of a parent gas *m<sup>j</sup>* is given by the difference between the weighings of a cylinder before and after adding the parent gas to it. Since the relationship between these two observed quantities is linear, it can be expressed in a matrix notation as follows:

$$
m = C_r r, \tag{4.5}
$$

where

*r* is a vector of mass readings of sample cylinder shown in table 4.1,

*m* a vector of the masses of the parent gases,

.

*Cr* the sensitivity matrix relating *m* to *r* .

The elements of *r* refer alternating to weighing the sample cylinder and the reference cylinder. In this example,  $m$  is a vector of length  $p = 2$ .  $C_r$  is a matrix of dimension  $p \times (p + 1)$ , which, in this case is as follows:

$$
C_r = \left[ \begin{array}{rrr} -1 & 1 & 0 \\ 0 & -1 & 1 \end{array} \right]
$$

Using the LPU from JCGM 102:2011, the covariance matrix associated with  $\hat{m}$  can be evaluated as

$$
V_{\widehat{m}}=C_rV_{\widehat{r}}C_r^{\top}.
$$

where

 $C_r$  denotes the sensitivity matrix relating  $m$  to  $r$ ,

 $V_{\widehat{r}}$  the covariance matrix associated with  $\widehat{r},$ 

 $V_{\hat{m}}$  the covariance matrix associated with  $\hat{m}$ 

**4.3.3.2** As the mass measurements  $\widehat{r}_j$  ( $j = 1 \dots p+1$ ), shown in table [4.1,](#page-31-0) can be assumed to be unconvoluted [100], the exerging equation  $V_j$  is the disconsel metrix. to be uncorrelated [109], the covariance matrix  $V_{\widehat{r}}$  is the diagonal matrix

$$
V_{\widehat{r}} = \text{diag}\left\{u^2(\widehat{r}_1), u^2(\widehat{r}_2), u^2(\widehat{r}_3)\right\},\,
$$

where  $u^2(\hat{r}_j)$  denotes the squared standard uncertainty associated with  $\hat{r}_j$ . Here we use  $u(\hat{r}) = 2 \text{mg}$  is used for all weighings of the cylinders and we take the following mean estimates of the mass measurements calculated from data in table [4.1:](#page-31-0)

 $\hat{r} = [7770.937 \text{ g} \quad 8016.471 \text{ g} \quad 8640.268 \text{ g}].$ 

### <span id="page-34-0"></span>**4.3.4 Molar masses**

**4.3.4.1** The molar masses are calculated from the standard atomic weights *A* and the subscripts in the formulæ of the molecules involved [[64,](#page-131-1) [109](#page-134-6)]. The vector *M* holding the *q* molar masses of the parent gas impurities is given by the following matrix equation:

<span id="page-34-1"></span> $M = \nu A.$  (4.6)

where

*M* denotes the vector holding the molar masses of the *q* components in the final mixture,

*ν* the matrix holding the coefficients in the *q* molecular formulæof the components,

*A* denotes the vector holding the *e* standard atomic weights of the elements occurring in the components, as shown in table 4.4.

The subscripts in the molecule formulæ , presented in a *q*×*e* matrix *ν*, are constants, where each row corresponds to a component of the prepared gas mixture and columns correspond to different chemical elements, as shown in table 4.4.

**4.3.4.2** The estimate  $\widehat{M}$  of *M* is given by  $\widehat{M} = \nu \widehat{A}$ , where  $\widehat{A}$  is the estimate of *A*. The covariance matrix associated with  $\tilde{M}$  is obtained by applying the LPU to equation [\(4.6\)](#page-34-1):

<span id="page-34-3"></span>
$$
V_{\widehat{M}} = \nu V_{\widehat{A}} \nu^{\top},\tag{4.7}
$$

where  $V_{\hat{A}} = \text{diag} \{ u^2(\hat{A}_1), \dots, u^2(\hat{A}_e) \}$ . Here we assume the independence between the standard atomic weights of the elements. This is an appropriate assumption with the standard atomic weights, although not applicable generally to atomic weights in normal materials (see clause [3.4.2\)](#page-29-0) [[89,](#page-132-1) [110](#page-134-3)].

4.34 Molar masses are calculated from the standard ato[mi](#page-35-1)c weights Anal 1<br>
4.34.1 The molar masses of chemical committee of the molecules involved [64, 109). The vector M holding the<br>
and the formulas of the parent gas imp **4.3.4.3** Here we adopt the modelling assumption of treating all standard atomic weights of the elements using a rectangular probability distribution whose bounds are provided by International Union of Pure and Applied Chemistry (IUPAC) (see table [4.3\)](#page-34-2). Recall that a rectangular distribution R(*a*, *b*) with endpoints *a* and *b* (*a < b*) has an expectation p  $(a + b)/2$  and a standard deviation  $(b - a)/\sqrt{12}$  [[6,](#page-126-5) clause 4.3.7]. Guidance on the use of atomic weights and assigning probability density functions (PDFs) is given in [[89,](#page-132-1) [110](#page-134-3)].

<span id="page-34-2"></span>

Element	Distribution	a	h		u(A)
H	R(a, b)	1.00784	1.00811	1.00798	0.00008
	R(a, b)	12.0096	12.0116	12.0106	0.0006
N	R(a, b)	14.00643	14.00728	14.00686	0.00025
$\cup$	R(a, b)	15.99903	15.99977	15.99940	0.00021
Ar.	R(a, b)	39.792	39.963	39.878	0.049

Table 4.3: Standard atomic weights of the makeup elements for components in the two parent gases

**4.3.4.4** Combining the subscripts in the molecule formulæ of the components and the atomic weights of the elements (Eqs. [4.6](#page-34-1) and [4.7\)](#page-34-3) provides molecular weight estimates which are summarised in table [4.4.](#page-35-1)

<span id="page-35-1"></span>

	$\nu(H)$	$\nu(C)$	v(N)	v(O)	v(Ar)	$\widehat{M}$	$u(\hat{M})$
Ar	$\boldsymbol{0}$	0	0	0	1	39.878	0.049
$CH_4$	4	$\mathbf 1$	0	0	0	16.0425	0.0007
CO	0	1	0	1	$\boldsymbol{0}$	28.0100	0.0006
CO <sub>2</sub>	0	$\mathbf{1}$	0	2	$\mathbf{0}$	44.0094	0.0007
H <sub>2</sub>	2	0	0	0	0	2.01595	0.00016
$H_2O$	$\overline{2}$	$\overline{0}$	$\bf{0}$	$\mathbf{1}$	0	18.0154	0.0003
$N_2$	$\mathbf 0$	0	2	0	0	28.0137	0.0005
O <sub>2</sub>	$\mathbf 0$	$\overline{0}$	0	2	0	31.9988	0.0004
owing way [7, Definition 3.20, Note 3]: $V_{\widehat{M}} = DR_{\widehat{M}}D,$							<b>I.3.4.5</b> The covariance matrix $V_{\widehat{M}}$ is related with the correlation matrix $R_{\widehat{M}}$ in the fol- (4.8)
vhere the matrix $\bm{D}$ is defined as							
$\mathbf{D} = \text{diag}\{u(\widehat{M}_1) \dots u(\widehat{M}_q)\}.$							
							Conversely, the correlation matrix is obtained from the covariance matrix as follows:
$R_{\widehat{M}} = D^{-1} V_{\widehat{M}} D^{-1}.$							(4.9)
llkanes, which have as general molecular formula $C_nH_{2n+2}$ $(n \ge 1)$ ). The upper triangular part of the correlation matrix $R_{\widehat{M}}$ is 4.3.4.6							
0 1	0	0	$\boldsymbol{0}$		0	0 0	
1		0.825	0.707	0.475	0.280	0 0	
	1		0.960 $\boldsymbol{0}$		0.280	$\boldsymbol{0}$	0.347
		$\mathbf{1}$		$\Omega$	0.481	0	0.595
				$\mathbf{1}$	0.589	0 0	
					1	0	0.808
						1 0 1	
of the second and third components, CH <sub>4</sub> and CO.							We observe a large correlation (0.960) between the third and fourth components, CO and $CO2$ , whereas the second-largest correlation (0.825) is between the molar masses
1.3.5 Composition of the parent gases							
mportant role and one can consider three distinct scenarios [109]:							<b>4.3.5.1</b> When preparing a gas mixture, the composition of the parent gases plays an

Table 4.4: Atomic composition coefficients and molar masses of the components in the two parent gases

$$
V_{\widehat{M}} = DR_{\widehat{M}}D,\tag{4.8}
$$

$$
D = \mathrm{diag}\{u(\widehat{M}_1) \ \ldots \ u(\widehat{M}_q)\}.
$$

$$
R_{\widehat{M}} = D^{-1} V_{\widehat{M}} D^{-1}.
$$
\n
$$
(4.9)
$$



### <span id="page-35-0"></span>**4.3.5 Composition of the parent gases**

- 1. The parent gas is a near-pure gas whose impurities have been characterized;
- 2. The parent gas is a gravimetrically prepared gas mixture;
- 3. All components in the parent gas have been analysed.
**4.3.5.2** In all cases, the covariance matrix  $V_{\hat{x}_j}$  of the amount fractions in each parent gas is not diagonal [[109](#page-134-0)]. Due to the fact that the amount fractions of all components in the parent gas mixtures add up to unity, these covariance matrices are always singular [2].

**4.3.5.3** The covariance matrix  $V_{\hat{x}}$  associated with the vector  $\hat{x}$  is block diagonal when working with near-pure gases, i.e.,  $V_{\hat{x}} = \text{diag}\left\{V_{\hat{x}_1}, \ldots, V_{\hat{x}_p}\right\}$ . This situation can be different for a multistage gas mixture preparation, for there exist correlations between the composition of a parent gas and the molar masses. Furthermore, there can be correlations between the composition of two or more parent gases.

**4.3.5.4** In this example we consider the scenario where all impurity components have been independently characterized and the amount fraction of the main component estimated by difference. In such a scenario, the covariance between the calculated (major) component  $x_{\text{calc}}$  and any other impurity amount fraction  $x_i$  is  $V(x_{\text{calc}}, x_i) = -u^2(x_i)$ .

#### **4.3.6 Composition of the gas mixture**

**4.3.6.1** Consider a vector *n*, containing the amounts of the components and their total amount in the mixture. That is,  $\boldsymbol{n}^{\top} = (n_1, \ldots, n_q, n_{\text{mix}})$ . From this vector, one can compute the vector *y* that is the amount fractions of the individual components in the gas mixture. The sensitivity matrices are obtained with respect to the vector *n*. The expression for  $U(\hat{n})$ is

<span id="page-36-0"></span>
$$
U_{\hat{n}} = C_a U_{\hat{a}} C_a^{\top},\tag{4.10}
$$

where *C<sup>a</sup>* can be written in block matrix notation as

 $C_a = \begin{bmatrix} C_m & C_M & C_{x_1} & \cdots & C_{x_p} \end{bmatrix}$ .

**4.3.6.2** The matrices  $C_m$ ,  $C_M$ ,  $C(x_j)$  hold the sensitivity coefficients of all input parameters. The expressions for the elements of these matrices have been derived elsewhere [[109](#page-134-0)]. The sensitivity matrix  $C_m$  of dimension  $(q + 1) \times p$  is obtained as follows:

$$
C_m = \left[ \begin{array}{cccc} \frac{x_{11}}{M_{\text{P},1}} & \frac{x_{12}}{M_{\text{P},2}} & \cdots & \frac{x_{1p}}{M_{\text{P},p}} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{x_{q1}}{M_{\text{P},1}} & \frac{x_{q2}}{M_{\text{P},2}} & \cdots & \frac{x_{qp}}{M_{\text{P},p}} \\ \frac{1}{M_{\text{P},1}} & \frac{1}{M_{\text{P},2}} & \cdots & \frac{1}{M_{\text{P},p}} \end{array} \right]
$$

where  $M_{P,j}$  is the average molar mass of a parent gas *j*, as given by equation [\(4.2\)](#page-31-0). Similarly, the sensitivity matrix  $C_M$  of dimension  $(q + 1) \times q$  is obtained as follows:

,

\n- **4.3.5.2** In all cases, the covariance matrix 
$$
V_{\bar{x}_j}
$$
 of the amount fractions in each parent gas not diagonal [109]. Due to the fact that the amount fractions of all components in the parent gas mixtures add up to unity, these covariance matrices are always singular [2].
\n- **4.3.5.3** The covariance matrix  $V_{\bar{x}}$  associated with the vector  $\hat{x}$  is block diagonal when working with near-pure gases, i.e.,  $V_{\bar{x}} = \text{diag} \{ V_{\bar{x}_j}, \ldots, V_{\bar{x}_p} \}$ . This situation can be different for a multistage gas mixture preparation, for there exist correlations between the composition of two or more parent gases.
\n- **4.3.5.4** In this example we consider the scenario where all impurity components have been independent and the amount fraction of the main component as the mean component gas and the amount fraction of the main component and by difference. In such a scenario, the covariance between the calculated (major) component  $x_{\text{calc}}$  and any other impurity amount fraction  $x_i$  is  $V(x_{\text{calc}}, x_i) = -u^2(x_i)$ .
\n- **4.3.6** Composition of the gas mixture
\n- **4.3.6.** Composition of the gas mixture
\n- **4.3.6.** Conposition of the gas mixture
\n- **4.3.6.** Conduction in block matrix notation as
\n- $C_a = \begin{bmatrix} C_m & C_M & C_{x_1} & \ldots & C_{x_p} \end{bmatrix}$ .\n .

\n
\n- **4.3.6.2** The matrices  $C_m$ ,  $C_M$ ,  $C(x_j)$  hold the sensitivity coefficients of all input parameters. The expressions for the elements of these matrices have been derived elsewhere [109]. The sensitivity matrix  $C_m$  of dimension  $(q+1) \times p$  is obtained as follows:
\n- <math display="inline</li>

.

and the sensitivity matrix  $C_{x_j}$  of dimension  $(q + 1) \times q$  is given as follows:

$$
\mathbf{C}_{x_j} = -\frac{m_j}{M_{\text{P},j}^2} \left[ \begin{array}{cccc} x_{1j}M_1 - M_{\text{P},j} & x_{1j}M_2 & \dots & x_{1j}M_q \\ x_{2j}M_1 & x_{2j}M_2 - M_{\text{P},j} & \dots & x_{2j}M_q \\ \vdots & \vdots & \vdots & \vdots \\ x_{qj}M_1 & x_{qj}M_2 & \dots & x_{qj}M_q - M_{\text{P},j} \\ M_1 & M_2 & \dots & M_q \end{array} \right]
$$

**4.3.6.3** The propagation of the uncertainty associated with the amount fraction of the components to the amount fractions can be carried out as follows. The covariance matrix  $V_{\hat{v}}$  associated with the vector  $\hat{y}$  of estimates of the amount fraction fractions of the *q* components in the mixture is calculated from the covariance matrix  $V_{\hat{n}}$  associated with  $\hat{n}$ using

<span id="page-37-0"></span>
$$
V_{\hat{y}} = C_n V_{\hat{n}} C_n^{\top}.
$$
\n(4.11)

The sensitivity matrix  $C_n$  of dimension  $q \times (q + 1)$  is given in block matrix notation by

$$
\mathbf{C}_n = \begin{bmatrix} n_{\text{mix}}^{-1} \mathbf{I} & -n_{\text{mix}}^{-2} \mathbf{n}' \end{bmatrix},\tag{4.12}
$$

where  $I$  denotes the identity matrix of dimension  $q \times q$  and  $n' = (n_1, \ldots, n_q)^\top$ . By substi-tuting equation [\(4.10\)](#page-36-0) for  $V_{\widehat{n}}$  into equation (4.11),  $V_{\widehat{y}}$  can be computed from  $V_{\widehat{a}}$ :

$$
V_{\hat{y}} = C_n C_a V_{\hat{a}} C_a^{\top} C_n^{\top}.
$$
\n(4.13)

The use of the chain rule of differentiation is implemented here by means of a matrix multiplication, which can be readily verified [109].

**4.3.6.4** The covariance matrix  $V_{\hat{v}}$  associated with a complete composition  $\hat{y}$  is singular, as the amount fractions of the components sum to unity [2].

**4.3.6.5** The results for the amount fractions for the mixture are given in table [4.5.](#page-38-0) The upper triangular part of the correlation matrix for the chemical composition of the gas mixture  $R_{\widehat{y}}$  is



The correlation matrix associated with the composition of the gas mixture shows that most correlations between pairs of amount fraction are rather weak; only the correlations between the amount fractions of argon and nitrogen (−0.50), and between carbon dioxide and nitrogen (−0.86) have appreciable influence. The joint PDF of the latter is shown in figure 4.1.

NOTE Due to round-off effects, the covariance matrix corresponding to the correlation matrix  $R_{\hat{v}}$ can appear to be non-singular. The covariance matrix  $V_{\widehat{y}}$  is however singular.



<span id="page-38-1"></span>Figure 4.1: Contour plot of the bivariate normal probability density function for  $N_2(y_7)$ and  $CO<sub>2</sub>(y<sub>4</sub>)$  amount fractions in the gas mixture with contours corresponding to 68 %, 95 %, 99 % and 99.9 % coverage probability



<span id="page-38-0"></span>

### **4.4 Discussion**

This example illustrates some of the advantages of using the multivariate form of the LPU (JCGM 102:2011). First, the LPU does not only provide the variances associated with the estimates of the various output quantities but also their covariances. Second, the multivariate modelling aids in structuring the model in a fashion that can be readily implemented. Important submodels are the formation of the covariance matrices of the transferred masses, molar masses, and the composition of the parent gases. The use of multivariate modelling is particularly useful in describing multi-stage gas mixture preparation.

# **5 Greenhouse gas emission inventories**

## <span id="page-39-1"></span>**5.1 Preamble**

## **5.1.1 General**



**5 [G](#page-130-0)reenhouse gas emission inventories**<br>
5.1 Preamble<br>
1.1.1 [C](#page-126-1)ommit<br>
1.1.1.1 Committee of the complision of consumption of equality data to provide second<br>
1.1.1.1 The complision of consumptions functions are activated pr **5.1.1.1** The compilation of country-scale greenhouse gas emissions' inventories involves combining standard emissions' factors and activity data to provide sector emissions' estimates. This example constitutes a small study of this huge, impactful topic, concentrating on agriculture, where large uncertainties arise, and the effect of correlation on the results. It also gives a brief comparison of the attitudes taken to uncertainty evaluation by the Intergovernmental Panel on Climate Change (IPCC) and the Joint Committee for Guides in Metrology (JCGM).

**5.1.1.2** Greenhouse gas (GHG) emissions' inventories are usually obtained by combining standard emissions' factors and activity data to provide sub-sector emissions' estimates. These sub-sectors are then combined to give sector estimates and further combined to estimate total emissions.

**5.1.1.3** A specific example is given relating to part of the inventory in the agriculture, land use and waste sector to provide an estimate of total emissions and the associated uncertainty.

**5.1.1.4** Current practice provided by the IPCC regarding uncertainty propagation in the area is considered in the context of the IPCC Guidelines for National Greenhouse Gas Inventories [[56](#page-130-0)]. The IPCC practice is contrasted with that in the Guide to the expression of uncertainty in measurement (GUM) [6] and other related JCGM guides. In particular, the possible effects of correlation, often ignored in practice, on the estimate are considered.

## <span id="page-39-2"></span>**5.1.2 Greenhouse gas inventories**

**5.1.2.1** GHG inventories of the United Kingdom (UK) and other countries are compiled according to IPCC 2006 Guidelines [56]. Each year the inventory is updated to include the latest data available. Together with many other countries, the UK submits a report to the United Nations Framework Convention on Climate Change (UNFCCC) annually via a consolidated report, which contains all EU countries.

**5.1.2.2** The GHG inventory includes seven gases under the Kyoto Protocol, namely, carbon dioxide (CO $_2$ ), methane (CH $_4$ ), nitrous oxide (N $_2$ O), hydrofluorocarbons (HFCs), sulfur hexafluoride (SF<sub>6</sub>), perfluorocarbons (PFCs) and nitrogen trifluoride (NF<sub>3</sub>).

**5.1.2.3** A bottom-up calculation of emissions is based on contributions of the form

<span id="page-39-0"></span> $E$ missions = (Activity data)  $\times$  (Emissions' factor). (5.1)

In all there are 743 emission factors and more than 1700 sources of emissions. There are 17 key data sources together with involvement from industry, government and academics.

**5.1.2.4** In terms of emissions from the three largest GHG contributors  $\sim$  CO<sub>2</sub>, CH<sub>4</sub> and  $N_2O$  — the majority of the uncertainty in the GHG inventory comes from the agriculture, land use and waste sectors [[17](#page-127-0)], with these three sectors contributing nearly 90 % of the uncertainty in the total inventory emissions.

**5.1.2.5** European countries, especially the UK, are moving towards metrological assessment of current uncertainty quantification including accounting for the effect of correlated quantities. They also want to extend the current top-down validation activity [3] to cover the complete GHG inventory.

**5.1.2.6** Figure [5.1](#page-40-0) shows the breakdown of the inventory contributions across sectors. The fifth biggest contribution is agriculture. As reductions in other contributions are made, somewhat more readily, agriculture will in future make a larger proportional contribution.



<span id="page-40-0"></span>Figure 5.1: Inventory contributions across sectors

**5.1.2.7** An emission inventory can be expressed as a linear combination of terms of the form [\(5.1\)](#page-39-0):

<span id="page-40-1"></span>
$$
E = \sum_{i} F_i A_i, \tag{5.2}
$$

with

*E*: total GHG emission for a given sector, geographic area and time period,

*Fi* : emissions' factor for the emissions of a given pollutant from source category *i*,

*Ai* : activity for source category *i*.

**5.1.2.8** In practice, the model [\(5.2\)](#page-40-1) is often enlarged to incorporate scaling factors to allow total GHG emission to be given in terms of 'CO<sub>2</sub> equivalent' (CO<sub>2</sub>e). That is, each GHG in the table has a conversion or scaling factor associated with it, giving

<span id="page-40-2"></span>
$$
E = \sum_{i} s_i F_i A_i,\tag{5.3}
$$

An example incorporating such scaling factors is given in clause [5.4.](#page-42-0)

## **5.1.3 Combination and correlation**

**5.1.3.1** Inventory uncertainties are generally combined in quadrature (root-sum-squares), assuming no correlation between the input quantities involved, using the variant of the law of propagation of uncertainty (LPU) applicable to such cases [[6,](#page-126-1) clause 5.1.2, formula (10)].

**5.1.3.2** Figure [5.2](#page-41-0) shows the expanded uncertainty contributions across sectors with agriculture being one of the biggest contributors.



<span id="page-41-0"></span>Figure 5.2: Expanded uncertainty contributions (at the 95 % level) across sectors

**5.1.3.3** Emissions' inventories are critical to many environmental decision-making processes [[13](#page-127-1)]. Typical questions that decision makers may ask that motivate the need to deal with uncertainties in emissions' inventories relate to precision and bias in the estimates, whether the estimates are based upon measurement, modelling or expert judgment, the main sources of uncertainty in these estimates, and how uncertainty can be reduced. To answer such questions, a reasonable indication of the uncertainty associated with the estimates of the quantities  $F_i$ ,  $A_i$  and  $E$  in expression [\(5.3\)](#page-40-2) that comprise an emission inventory is needed.

## **5.1.4 IPCC and JCGM reporting guidelines**

**5.1.4.1** The IPCC reporting guidelines refer to the 'error propagation approach' and the 'Monte Carlo approach' for uncertainty propagation. An extract from [[57](#page-130-1)], which uses these approaches, is:

'Comparing the results of the error propagation approach, and the Monte Carlo estimation of uncertainty by simulation, is a useful quality control check on the behaviour of the Monte Carlo model.

'The reason that the error propagation approach is used as a reference is because the approach to the error propagation approach has been defined and checked by the IPCC, and is clearly set out in the IPCC 2000 Good Practice Guidance and the 2006 Guidelines.'

**5.1.4.2** In this example, the IPCC approach to uncertainty evaluation in the context of greenhouse gas emissions is considered. Statements are made concerning the extent of

### **5.2 Specification of the measurand**

### **5.3 Measurement model**

### <span id="page-42-0"></span>**5.4 Uncertainty propagation**

**5.4.3** This table incorporates the scaling mentioned in clause [5.1.2.](#page-39-2) The scaling factors to give 'CO<sub>2</sub> equivalent' equivalent emissions are  $s = 11/3 = 3.667$  for CO<sub>2</sub> as C,  $s = 25$ for methane  $\left( \text{CH}_4 \right)$  and  $s = 298$  for nitrous oxide  $\left( \text{N}_2 \text{O} \right)$ .

<span id="page-42-1"></span>Table 5.1: Reported activity data and emissions' factors with associated relative expanded uncertainties for 95 % confidence

the alignment with guidance provided by the JCGM. Wherever measurement uncertainty is used, it is important that it is evaluated in the same manner by different practitioners so that users' interpretation is consistent and reliable. Universally agreed methods for uncertainty evaluation are hence required.										
5.2	Specification of the measurand									
A meaningful interpretation of the measurand, used here, is 'total emissions for a given pollutant, geographic area and time period' (see clause 5.1).										
<b>Measurement model</b> 5.3										
The IPCC does not seem explicitly to use the concept of 'measurement model' in their documents. However, expression (5.3) can be regarded as a basic measurement model central to IPCC considerations.										
5.4 Uncertainty propagation										
<b>5.4.1</b> The use of the model specified as expression (5.3), in which we take the sector as mobile machinery in agriculture (IPCC category A4ci), the region as the UK and the time period as 2018, is considered. Uncertainties are propagated using the law of propa- gation of uncertainty in JCGM 100:2008 [6], comparable to the IPCC's 'error propagation' method. The quantities concerned are taken as independent. Then, perceived correlation is considered.										
5.4.2 Data for this sector from the UK National Atmospheric Emissions Inventory site (2018) are given in table 5.1 where the number of significant decimal digits is stated as reported. This total comprises contributions from the use of two fuels, 'Gas Oil' and Petrol, reported in terms of the 2018 Activity Data (the $A_i$ ) for those fuels. The table also specifies the Emissions' Factors (the $F_i$ ) for each of the key GHGs (CO <sub>2</sub> , CH <sub>4</sub> and N <sub>2</sub> O) for those fuels. The reported expanded uncertainties $U_{rel}$ (at the 95% confidence level) on the individual activity data and emission factors are given. 5.4.3 to give 'CO <sub>2</sub> equivalent' equivalent emissions are $s = 11/3 = 3.667$ for CO <sub>2</sub> as C, $s = 25$ for methane (CH <sub>4</sub> ) and $s = 298$ for nitrous oxide (N <sub>2</sub> O).	This table incorporates the scaling mentioned in clause 5.1.2. The scaling factors									
Table 5.1: Reported activity data and emissions' factors with associated relative expanded uncertainties for 95 % confidence										
Fuel Gas	$A_i/(TJ)$		$U_{\rm rel}(A_i)/\%$ $F_i/(\text{kt }T\text{J}^{-1})$	$U_{\rm rel}(F_i)/\%$						
Gas oil CO <sub>2</sub>	56317.0920	38.6	$2.0438 \times 10^{-2}$	2.7						
Gas oil CH <sub>4</sub>	56317.0920	1.6	$3.5368 \times 10^{-6}$	80.0						
Gas oil $N_2O$	56317.0920	1.6	$3.0984 \times 10^{-6}$	216.3						
Petrol CO <sub>2</sub>	67.1856	50.7	$1.9127 \times 10^{-2}$	2.0						
Petrol CH <sub>4</sub>	67.1856	1.6	$4.8654 \times 10^{-5}$	80.0						
Petrol $N_2O$	67.1856	1.6	$3.3578 \times 10^{-7}$	216.3						

$$
E = \sum_{i=1}^{6} E_i = \sum_{i=1}^{6} s_i A_i F_i,
$$
\n(5.4)

<span id="page-43-0"></span>Table 5.2: Emissions as  $CO<sub>2</sub>$  equivalent for each contribution and the total  $CO<sub>2</sub>$  equivalent

					<b>5.4.4</b> Since $A_1, \ldots, A_6$ denote the activity data values and $F_1, \ldots, F_6$ the corresponding emissions' factors in table 5.1 in the order given there, comparing with expression (5.3),	
	$E = \sum_{i=1}^{6} E_i = \sum_{i=1}^{6} s_i A_i F_i,$					(5.4)
					where the $s_i$ are the scaling factors, assumed to have no associated uncertainty.	
5.4.5			635 kt CO <sub>2</sub> e (relative expanded uncertainty of 38%).		Total emissions for this sector were reported to be $4282 \text{ kt } CO_2$ e. The reported xpanded uncertainty (at the 95% level of confidence) associated with this total was	
5.4.6				rom the UK National Atmospheric Emissions Inventory site (2018).	Table 5.2 summarizes the emissions' calculations for this sub-sector, namely, the emissions as $CO2$ equivalent for each contribution and the total $CO2$ equivalent, obtained Table 5.2: Emissions as $CO_2$ equivalent for each contribution and the total $CO_2$ equivalent	
	GHG	Fuel	Reported	Scaling	CO <sub>2</sub> e/	
			emissions/kt	factor	kt	
	$CO2$ as C	Gas oil	1150.9806	3.667	4220.2623	
	CH <sub>4</sub>	Gas oil	0.1992	25 <sub>1</sub>	4.9795	
	$N_2O$	Gas oil	0.1745	298	51.9987	
	$CO2$ as C	Petrol	1.2851	3.667	4.7120	
	CH <sub>4</sub>	Petrol	0.0033	25	0.0817	
	$N_2O$	Petrol	0.0000	298	0.0067	
	Total				4282.0409	
	tandard uncertainty $u_{\text{ind}}(E)$ :		$u_{\text{ind}}^2(E) = \sum_{i=1}^6 s_i^2 \left[ A_i^2 u^2(F_i) + F_i^2 u^2(A_i) \right].$		The application of the LPU [6, eqn. (10)] for independent quantities gives the associated The use of the data in tables 5.1 and 5.2 gives $E$ , as above, and by applying expression (5.2),	(5.5)
			$E = 4282.04 \text{ kt}, \qquad u_{\text{ind}}(E) = 835.16 \text{ kt}.$			(5.6)
			The uncertainty budget is given in table 5.3.			
.4.7					Perceived correlations associated with the various quantities, which seem not to be trongly considered in the IPCC documents, are now considered. It appears that correlation s not entertained with the 'error propagation' approach, although it is stated, without giving detail, that the Monte Carlo method may be used for this purpose [55, clause 6.3].	

<span id="page-43-1"></span>
$$
u_{\text{ind}}^2(E) = \sum_{i=1}^6 s_i^2 \left[ A_i^2 u^2(F_i) + F_i^2 u^2(A_i) \right].
$$
\n(5.5)

$$
E = 4282.04 \,\text{kt}, \qquad u_{\text{ind}}(E) = 835.16 \,\text{kt}. \tag{5.6}
$$

<span id="page-44-0"></span>

**5.4.8** Examining the activity contributions in table 5.1,  $A_2$  and  $A_3$  are identical numerically as are their associated standard uncertainties  $u(A_2)$  and  $u(A_3)$ . A similar statement applies to  $A_5$  and  $A_6$  and their associated standard uncertainties. This observation raises the possibility that  $A_2$  and  $A_3$  have a common origin and might even be the same quantity, and similarly for  $A_5$  and  $A_6$ . In discussion with GHG experts, this perception is considered to be more than reasonable.

 $\begin{tabular}{|c|c|} \hline & {10b8-5.3: Uncertainly budget}\\ \hline \hline G16G & Fe101 & CO$_2$ & std. Inc/Al\\ \hline G18G & Sc101 & 2.0729\\ \hline G19c & Sc101 & 2.1798\\ \hline G10 & 57.3866\\ \hline G10 & 57.3866\\ \hline G11 & 8401\\ \hline G14 &$ **5.4.9** Assuming this perception to be correct, an adjustment to the above uncertainty calculation is needed. When several input quantities are associated via a common origin or effect, the resulting correlation can typically be handled either by constructing and using an appropriate covariance matrix, or by re-parametrizing to isolate the common effect. In this case, whether the (numerically identical) activities  $A_2$  and  $A_3$ , and  $A_5$  and  $A_6$ , are either very highly correlated or are simply repeated use of the same information, it is simplest to re-parametrize by replacing the nominally separate instances by a single term. So, taking  $A_2$  and  $A_3$  as the same quantity, and similarly for  $A_5$  and  $A_6$ , instead of the model [\(5.3\)](#page-40-2), we consider the model

<span id="page-44-1"></span>
$$
E = s_1 A_1 F_1 + A_2 (s_2 F_2 + s_3 F_3) + s_4 A_4 F_4 + A_5 (s_5 F_5 + s_6 F_6).
$$

For the current data, *E* is identical to the previously calculated value, as expected and, applying LPU, since the *s<sup>i</sup>* are constants,

$$
u_{\text{cor}}^2(E) = \sum_{i=1,4} s_i^2 \left[ A_i^2 u^2 \left( F_i \right) + F_i^2 u^2 \left( A_i \right) \right]
$$
  
+ 
$$
\sum_{i=2,5} \left\{ A_i^2 \left[ s_i^2 u^2 \left( F_i \right) + s_{i+1}^2 u^2 \left( F_{i+1} \right) \right] + \left( s_i F_i + s_{i+1} F_{i+1} \right)^2 u^2 \left( A_i \right) \right\},
$$
 (5.7)

that is,

$$
u_{\rm cor}(E) = 835.16 \,\text{kt},\tag{5.8}
$$

the same value to the number of decimal digits displayed as obtained when not taking correlation into consideration. Algebraically, however, using expressions [\(5.5\)](#page-43-1) and [\(5.7\)](#page-44-1),

$$
u_{\text{cor}}^2(E) - u^2(E) = 2s_2 F_2 s_3 F_3 u^2(A_2) + 2s_5 F_5 s_6 F_6 u^2(A_5),\tag{5.9}
$$

which is positive and demonstrated by displaying  $u^2(E)$  and  $u_{\rm cor}^2(E)$  to a greater number of decimal places. This result suggests that considering correlation is not meaningful for such calculations. However, we present rationale why accounting for covariance in emissions' calculation is generally appropriate.

**5.4.10** The reason for the increase in the standard uncertainty associated with total emissions being so small is that one particular contribution to the budget,  $CO<sub>2</sub>$  as C for Gas oil, the first in table [5.2,](#page-43-0) is by far the most dominant. This effect has the consequence that the consideration of correlation is *unnecessary* in this example.

**5.4.11** The principle stands that such correlations should generally be taken into account since they might make a significant difference where one uncertainty contribution is less dominant. It would be appropriate to identify an instance where the contribution of the different sub-sectors is more balanced and the effect of correlation therefore more pronounced.

**5.4.12** It would be possible to handle the uncertainty propagation using [6, formula (13)] taking correlation into account for the quantities concerned. The use of variable substitution as above reparametrizes the problem in such a way that the need to take correlation directly into account is avoided. The assumption is made that the remaining quantities are independent, which should be verified in practice to the extent possible.

**5.4.13** Insight into the difference in the standard uncertainty in the measurand under the assumptions of independence and dependence (in the above respect) is given by taking an extreme case.

**5.4.14** Consider the following idealized scenario. For all *i*, take  $A_i = A$ ,  $F_i = F$ ,  $u(A_i) = u_A$ ,  $u(F_i) = u_F$  and  $s_i = s$ . Then, expression (5.5) becomes

<span id="page-45-0"></span>
$$
u_{\text{ind}}^2(E) = 6s^2 A^2 u_F^2 + 6F^2 s^2 u_A^2 = s^2 A^2 F^2 \left[ 6u_{\text{rel}}^2(F) + 6u_{\text{rel}}^2(A) \right]
$$
(5.10)

with relative standard uncertainties

$$
u_{\rm rel}(F) = \frac{u_F}{F}, \qquad u_{\rm rel}(A) = \frac{u_F}{A}.
$$

On the other hand, from expression (5.7),

<span id="page-45-2"></span><span id="page-45-1"></span>
$$
u_{\text{cor}}^2(E) = 6s^2 A^2 u_F^2 + 18F^2 s^2 u_A^2 = s^2 A^2 F^2 \left[ 6u_{\text{rel}}^2(F) + 18u_{\text{rel}}^2(A) \right].\tag{5.11}
$$

The quotient of the terms in square brackets in expressions [\(5.10\)](#page-45-0) and [\(5.11\)](#page-45-1), namely,

$$
\lambda = \frac{u_{\text{cor}}^2(E)}{u_{\text{ind}}^2(E)} = \frac{6u_{\text{rel}}^2(F) + 18u_{\text{rel}}^2(A)}{6u_{\text{rel}}^2(F) + 6u_{\text{rel}}^2(A)} = \frac{u_{\text{rel}}^2(F) + 3u_{\text{rel}}^2(A)}{u_{\text{rel}}^2(F) + u_{\text{rel}}^2(A)},
$$
(5.12)

**L4.10** The reas[o](#page-43-1)n for th[e](#page-126-1) interests in the standard uncertainty associated with total angle<br>is the standard interest in the standard uncertainty associated with to<br>the first in Table 5.2, is typ for the most downthant in is informative regarding the respective contributions from the emissions' factor and the activity data under the above correlation and independence assumptions. Examining expression (5.12), for an activity data relative uncertainty that is small compared with that for the emissions' factor,  $\lambda$  is close to unity and in that case the assumption of independence is reasonable. For the converse,  $\lambda$  is close to 3 with the consequence that, accounting for is reasonable. For the converse,  $\lambda$  is close to 3 with the consequence that, accounting for<br>correlation,  $u_{\rm cor}(E)$  is  $\sqrt{3} \approx 1.7$  times  $u_{\rm ind}(E)$ , a 70% increase. We emphasize that such an increase is optimistic. For such cases, where better balanced data in terms of the values of the quantities involved is available, the expected increase is bounded by 70 %.

# **5.5 Interpretation of results**

5.5 Interpretation of results<br>
5.5 In elative expanded uncertainties of 100% or more are not unc[om](#page-42-1)mon in emission<br>
inventories, which is why it is important these large uncertainty sources are expedially tidentified<br>
inve **5.5.1** Relative expanded uncertainties of 100 % or more are not uncommon in emissions' inventories, which is why it is important these large uncertainty sources are carefully taken into consideration. The UK introduction to GHG inventories [[96](#page-133-0)], shows a highly asymmetric interval for inventory  $N_2O$  figures, in that emissions' estimates for  $N_2O$  are far from the centre of their 95 % confidence intervals, in its figure 5.2. The implication is that the probability distributions for  $N<sub>2</sub>O$  are strongly asymmetric, but in the previous clause (5.1) of that document uncertainties are expressed as 95 % confidence intervals (for example, table 5.1), with no indication that they may be asymmetric.

**5.5.2** Some areas within emissions' inventories are so uncertain that it is not known with confidence whether they are sources or sinks. However, the referenced figure implies that  $N<sub>2</sub>O$  from agriculture is definitely a source so the probability distribution must be asymmetric. However, the report does not provide an explanation for the figure.

**5.5.3** After some investigation, these aspects seem to be based on methodology described in Milne et al. [[77](#page-132-0)]. Figure 2 in that paper shows the distribution given by Monte Carlo modelling and the asymmetry is attributed to skew in the emissions' factors. Log-normal probability distributions are used to represent knowledge of such quantities.

**5.5.4** However, reference [[77](#page-132-0)] is focused on agricultural emissions (that is, manure, enteric fermentation and field burning), but table 5.1 in this document is based on non-road vehicle emissions (that is, from tractors, combine harvesters, etc.), which burn gas-oil and diesel. No information has been found by the authors of this chapter on the probability distributions for the quantities under consideration here.

**5.5.5** The focus here is on an agricultural sub-sector where there was sufficient data publicly available to carry out this type of assessment. The correlation issue seems more significant within other agricultural sub-sectors but the raw data needed to test this was not available. It is understood that such data and more extensive information concerning the numbers appearing in the UK inventory, from which correlations can be quantified, will subsequently be made available.

# **5.6 Uncertainty guidance promoted by the IPCC and the JCGM**

**5.6.1** In terms of regarding expression [\(5.3\)](#page-40-2) as the measurement model, JCGM 100:2008 [6] refers to the  $F_i$  and  $A_i$  as input quantities and  $E$  as the output quantity or measurand. In the model (5.3), the expression on the right-hand side constitutes the *measurement function*, with sensitivity coefficients given by the first-order partial derivatives of that function. IPCC documents seem not to emphasize that propagation of uncertainty should be carried out in the context of a *measurement model*. It is only by using measurement models as a basis, that measurement uncertainty evaluation can properly be conducted.

**5.6.2** The IPCC 2000 Good Practice Guidance [[100](#page-133-1)] and the 2006 Guidelines [[56](#page-130-0)] set out the 'error propagation approach' referred to as the LPU in the GUM [[6](#page-126-1)]. For example, the LPU occurs specifically in [[55,](#page-130-2) clause 6.3, formula (6.3)].

**LEGA** The nuitable of the UGCM is that the Moone Carlo method (NCO), as described<br>given the momenta of the CoM H22-2D is hearing suited to handle situations,<br>where the momentaties are large and the most uncertained its a **5.6.3** The attitude of the JCGM is that the Monte Carlo method (MCM), as described by JCGM 101:2008 and JCGM 102:2011 [[12,](#page-126-3) [7](#page-126-4)], is better suited to handle situations where the uncertainties are large and the measurement model is non-linear. In contrast, the error propagation approach (the LPU in JCGM 100:2008 [[6](#page-126-1)]) can produce invalid uncertainty statements. Thus, MCM can be used as a reference against which other methods can be assessed such as the evaluation of LPU, which requires linearization of the model. This way of thinking is very different from that of the IPCC, which regards the error approach as a reference. Indeed, MCM is seen, at least by some Bayesian statisticians, as a 'gold standard' for uncertainty propagation: Huggins et al. [[53](#page-130-3)] state:

'Classical Monte Carlo methods . . . remain the gold standard for approximate Bayesian inference because they have a robust finite-sample theory and reliable convergence diagnostics.'

There are many other references (such as [33, 52, 88, 101]) to the Monte Carlo method as a 'gold standard' for uncertainty propagation. In addition, non-pathological examples are well documented to show disagreement of the results produced by the two approaches. One example in mass calibration, from JCGM 101:2008 [12], shows that the LPU can underestimate the standard uncertainty by some 40 % compared to the MCM.

**5.6.4** Within the context of its GHG emissions' inventory, the UK has implemented the error propagation approach as set out in its guidance. It is stated:

'Uncertainty estimation by simulation (MCM) cannot be prescriptive, and will depend on how the country constructs its model and the correlations considered. Therefore, 'there is a greater likelihood of errors being introduced in the model used to estimate uncertainty by Monte Carlo simulation'.

**5.6.5** JCGM 101:2008 provides a different view on the matter [12]:

'Whereas there are some limitations to the JCGM 100:2008 uncertainty framework (GUF), the propagation of distributions [implemented by Monte Carlo] will always provide a PDF for the output quantity that is consistent with the model of the measurement and the PDFs for the input quantities.'

A comparison of MCM and the GUF is possible if the means and standard deviations of the probability distributions used by MCM are taken as the estimates and associated standard uncertainties used by LPU. There is no choice in the model itself. Both approaches use exactly the same model but, as stated, LPU linearizes it. Rather than errors being introduced by MCM, they will always be introduced by LPU unless the model is linear. Importantly, the main model used in GHG inventory work is [\(5.3\)](#page-40-2), which is non-linear.



# **6 Simple linear measurement models**

## **6.1 Preamble**

**6.1.1** This example considers the evaluation of uncertainty based on the simple linear measurement model

<span id="page-48-0"></span>
$$
Y = X_1 + \dots + X_N. \tag{6.1}
$$

when  $N = 4$ . The propagation of measurement uncertainty and the determination of coverage intervals using the provisions of the JCGM 100:2008 [6] and the Monte Carlo method (MCM) are shown. MCM is used to validate the application of JCGM 100:2008. Standard uncertainties are reported to three significant decimal places to facilitate their comparison.

6 Simpl[e](#page-126-1) linear measurement models and<br>solid Preamble considers the evaluation of uncertainty based on the simple<br>funcarion: This example considers the evaluation of uncertainty based on the simple<br>funcarion model  $Y = X_1 +$ **6.1.2** A special case of the generic linear model is considered for three different sets of probability density functions (PDFs)  $g_{x_i}(x_i)$  assigned to the input quantities  $X_i$ , which are all regarded as independent. The input quantities *X<sup>i</sup>* and hence the output quantity *Y* are all dimensionless in this example. For the first set, each  $g_{x_i}(x_i)$  is a standard normal PDF (with  $X_i$  having expectation zero and standard deviation unity). For the second set, each  $g_{_{X_i}}\!(x_i)$  is a rectangular PDF, also with  $X_i$  having expectation zero and standard deviation unity. The third set is identical to the second except that the PDF for  $g_{_{X_4}}\!(x_4)$  has a standard deviation of ten.

NOTE 1 Further information concerning linear models, such as the model (6.1), where the PDFs are normal or rectangular or a combination of both, is available [38].

NOTE 2 In accordance with JCGM 100:2008, for each instance, a probabilistically symmetric 95 % coverage interval for *Y* is based on a normal distribution with mean taken as the value of *Y* evaluated at the estimates of the input quantities and standard deviation as the evaluated standard uncertainty.

# <span id="page-48-1"></span>**6.2 Normally distributed input quantities**

**6.2.1** A standard normal PDF is assigned to each  $X_i$ . Thus, the estimates of the  $X_i$ are  $x_i = 0$ ,  $i = 1, 2, 3, 4$ , with associated standard uncertainties  $u(x_i) = 1$ .

**6.2.2** The law of propagation of uncertainty (LPU) [[6,](#page-126-1) clause 5.1.2] gives the standard uncertainty  $u(y) = 2.00$  associated with the estimate  $y = 0.00$  of *Y*, using a numerical tolerance of two significant decimal digits for  $u(y)$  [[12,](#page-126-3) clause 7.9.2]. A probabilistically symmetric 95 % coverage interval for *Y* is [−3.92, 3.92] based on a coverage factor of  $k = 1.96$ . These results are given in the first (non-header) row of table [6.1.](#page-49-0)

**6.2.3** The application of MCM [[12,](#page-126-3) clause 7] with  $M = 10^5$  trials and a numerical tolerance of  $\delta = 0.05$  gives the estimate  $\gamma = 0.00$ , its standard uncertainty  $u(\gamma) = 2.00$  and the probabilistically symmetric 95 % coverage interval [−3.92, 3.92]. See table [6.1,](#page-49-0) row 2. Two further applications of the method, with  $M = 10^6$  trials, agree with these results to within the numerical tolerance used. See table [6.1,](#page-49-0) row 2. These two further applications (different random samplings being taken from the PDFs) were made to demonstrate variation in the results obtained.

**6.2.4** Two applications of the adaptive Monte Carlo procedure [[12,](#page-126-3) clause 7.9] with the use of a numerical tolerance of *δ/*5 as recommended in [[12,](#page-126-3) clause 8.2] were then made, taking  $1.23 \times 10^6$  and  $1.02 \times 10^6$  trials. See table [6.1,](#page-49-0) rows 3 and 4.

**6.2.5** The PDF for *Y* obtained analytically is the normal PDF with expectation zero and standard deviation two. See table [6.1,](#page-49-0) bottom row.

<span id="page-49-0"></span>Table 6.1: Uncertainty evaluation for the linear model [\(6.1\)](#page-48-0) with a standard normal PDF assigned to each input quantity *X<sup>i</sup>*

Method	М		u(y)	Prob. sym. 95 % cov. int.	$d_{\rm lo}$	$a_{hi}$	<b>GUF</b> validated $(\delta = 0.05)$ ?
<b>GUF</b>		0.00	2.00	$[-3.92, 3.92]$			
MCM	10 <sup>5</sup>	0.00	2.00	$[-3.92, 3.92]$			
MCM	$10^{6}$	0.00	2.00	$[-3.92, 3.92]$			
MCM	$10^{6}$	0.00	2.00	$[-3.92, 3.92]$			
Adaptive MCM	$1.23 \times 10^{6}$	0.00	2.00	$[-3.92, 3.93]$	0.00	0.01	Yes
Adaptive MCM	$1.02 \times 10^{6}$	0.00	2.00	$[-3.92, 3.92]$	$0.00^{\circ}$	0.00	Yes
Analytical		0.00	2.00	$[-3.92, 3.92]$			

<span id="page-49-2"></span>**6.2.6** Figure [6.1](#page-49-1) shows the (normal) PDF for *Y* resulting from the JCGM 100:2008 uncertainty framework (GUF). It also shows one of the approximations (density histogram of  $M = 10^6$  model values of *Y*) constituting a discrete representation of this PDF provided by MCM [[12,](#page-126-3) clause 7.5]. The endpoints of the probabilistically symmetric 95 % coverage interval provided by both methods are shown as vertical lines. The PDF and the approximation are visually indistinguishable, as are the respective coverage intervals. For this example, such agreement would be expected (for a sufficiently large numerical value of *M*), because all the conditions hold for the application of the GUF [12, clause 5.7].



<span id="page-49-1"></span>Figure 6.1: Approximations for the model [\(6.1\)](#page-48-0), with each  $X_i$  assigned a standard normal PDF, to the PDF for *Y* provided by (a) the GUF (black line), (b) MCM (density histogram, clause 6.2.6). The analytical solution here coincides with the GUF solution. "Unit" denotes any unit

<span id="page-50-0"></span>
$$
g_Y(\eta) = \int_{-\infty}^{\infty} g_{X_1}(\eta - x_2) g_{X_2}(x_2) dx_2.
$$
 (6.2)

#### <span id="page-50-3"></span>**6.3 Rectangularly distributed input quantities with the same width**

The probabilistically symmetric 95 % coverage interval is determined, because the general linear measurement model (6.1) with independent input quantities has the property that the PDF for Y is symmetric when the PDFs for the $X_i$ are symmetric, which is proved as follows. For random variables $X_1$ and $X_2$ having PDFs that are symmetric about zero, their sum $Y = X_1 + X_2$ is also symmetric about zero. To establish this result (a more complicated proof appears in [94]), the PDF for $gY$ is given by the convolution integral							
	$g_Y(\eta) = \int_{-\infty}^{\infty} g_{X_1}(\eta - x_2) g_{X_2}(x_2) dx_2.$						(6.2)
To prove that $g_Y(\eta) = g_Y(-\eta)$ for expression (6.2), using symmetry negate the arguments of $g_{X_1}$ and $g_{X_2}$ , and set $\tau = -x_2$ . Comparison with the right side of (6.2) proves the result in this case. Observing that the result also applies for symmetry about nonzero expectations, by induction the result holds for the more general model (6.1).							
6.2.7 The validation procedures of [12, Clauses 8.1, 8.2] were applied. Using the ter- minology of [12, clause 7.9.2], $n_{\text{dig}} = 2$ , since two significant decimal digits in $u(y)$ are sought. Hence, $u(y) = 2.0 = 20 \times 10^{-1}$ , and so $c = 20$ and $\ell = -1$ in [12, clause 7.9.2, item a)]. Thus, according to that clause, the numerical tolerance is $\delta = \frac{1}{2} \times 10^{-1} = 0.05$ . The magnitudes $d_{\text{lo}}$ and $d_{\text{hi}}$ of the endpoint differences [12, eqns. (19), (20)] are shown in table 6.1, columns 6 and 7, for the two applications of adaptive MCM. Whether the GUF has been validated for $\delta = 0.05$ is shown in column 8.							
6.3 Rectangularly distributed input quantities with the same width							
<b>6.3.1</b> Assign a rectangular PDF to each $X_i$ , so that $X_i$ has an expectation of zero and a standard deviation of unity. The estimates of the $X_i$ are $x_i = 0$ ( $i = 1, 2, 3, 4$ ), with associated standard uncertainties $u(x_i) = 1$ . By following the analogous steps to those in clause 6.2, the results in table 6.2 were obtained. Table 6.2: Uncertainty evaluation for the linear model (6.1), with each input quantity $X_i$ assigned identical rectangular PDFs.							
Method				95 % cov. int.			M y $u(y)$ Prob. sym. $d_{\text{lo}}$ $d_{\text{hi}}$ GUF validated $(\delta = 0.05)$ ?
GUF <b>MCM</b> <b>MCM</b> MCM	10 <sup>5</sup> $10^{6}$ $10^6$	0.00 0.00 0.00 0.00	2.00 2.01 2.00 2.00	$[-3.92, 3.92]$ $[-3.90, 3.89]$ $[-3.89, 3.88]$ $[-3.88, 3.88]$			
Adaptive MCM <b>Adaptive MCM</b> Analytical	$1.02 \times 10^{6}$ $0.86 \times 10^{6}$	0.00 0.00 0.00	2.00 2.00 2.00	$[-3.88, 3.89]$ $[-3.87, 3.87]$ $[-3.88, 3.88]$	0.04 0.05	0.03 0.05	Yes No
The analytic solution for the endpoints of the probabilistically symmetric 95% coverage							
				interval, namely, $\pm 2\sqrt{3}[2-(3/5)^{1/4}] \approx \pm 3.88$ , is obtained as described below.			
The PDF of a rectangular distribution R(a, b) takes the constant value $(b-a)^{-1}$ for $a \le x \le$ and is zero otherwise. The <i>n</i> -fold convolution of R(0, 1) is the B-spline $B_n(x)$ of order <i>n</i> (degree $n-1$ ) with knots $0, \ldots, n$ [102] whose explicit expression is [25]							

<span id="page-50-1"></span>Table 6.2: Uncertainty evaluation for the linear model [\(6.1\)](#page-48-0), with each input quantity  $X_i$ assigned identical rectangular PDFs.

<span id="page-50-2"></span>
$$
B_n(x) = \frac{1}{(n-1)!} \sum_{r=0}^{n} {}^{n}C_r(-1)^r (x-r)_{+}^{n-1},
$$
\n(6.3)

where  ${}^nC_r = n!/[r!(n-r)!]$  and  $z_+ = \max(z, 0)$ . A piecewise cubic-polynomial representation of *B*<sup>4</sup> (*x*) obtained from expression [\(6.3\)](#page-50-2) or by applying the B-spline recurrence relation in [[25](#page-128-1)] is

$$
B_4(x) = \frac{1}{6} \begin{cases} x^3, & 0 \le x < 1, \\ 4 - 12x + 12x^2 - 3x^3, & 1 \le x < 2, \\ -44 + 60x - 24x^2 + 3x^3, & 2 \le x < 3, \\ (4 - x)^3, & 3 \le x \le 4, \\ 0, & \text{otherwise.} \end{cases}
$$
(6.4)

The left-hand endpoint  $y_{lo}$  of the probabilistically symmetric 95% coverage interval lies between zero and one, because

$$
\int_0^1 B_4(x) dx = \left[\frac{1}{24}x^4\right]_0^1 = \frac{1}{24},
$$

and  $0.025 = \frac{1}{40} < \frac{1}{24}$ . The area under the PDF that lies to the left of  $y_{\text{lo}}$  is given by

$$
\int_0^{d_{\rm L}} B_4(x) \, \mathrm{d}x = \frac{1}{24} d_{\rm L}^4 = \frac{1}{40}.
$$

That is,  $d_{\rm L} = (3/5)^{1/4}$ . By symmetry, the right-hand endpoint is  $d_{\rm R} = 4 - (3/5)^{1/4}$ . Thus, the probabilistically symmetric 95 % coverage interval is

 $\left[ (3/5)^{1/4}, 4-(3/5)^{1/4} \right]$  or  $2 \pm (2-(3/5)^{1/4}).$ 

The corresponding coverage interval for the four-fold convolution of the rectangular PDF p The corresponding coverage interval for the four-fold convolution of the rectangular PDF<br>R(−√3,√3) (which has zero expectation and unit standard deviation) is given by shifting  $R(-\sqrt{3}, \sqrt{3})$  (which has zero expectation and unit s<br>this result by two units and scaling it by 2 $\sqrt{3}$  units.

<span id="page-51-0"></span>**6.3.2** Figure [6.2](#page-52-0) shows the counterpart of figure 6.1 in this case, again based on  $M = 10^6$ trials and a numerical tolerance of  $\delta = 0.05$ . By comparison with figure [6.1,](#page-49-1) some modest differences between the approximations to the PDFs can be seen. Since normality of the measurand is assumed, the GUF provides exactly the same PDF for *Y* when the PDFs for the  $X_i$  are normal or rectangular, because the expectations of these quantities are identical, as are the standard deviations, in the two cases. The PDF provided by MCM takes smaller values than that provided by the GUF in the neighbourhood of the expectation and to a lesser extent towards the tails. It is slightly greater in the flanks. The endpoints of the coverage intervals provided are again almost visually indistinguishable, but table [6.2](#page-50-1) shows small differences.

where  ${}^{\circ}C_{c} = m/(t/(t-s^{2}))$  and  $z_{+} = \text{max}(z,0)$ . A piecewise cubic polynomial representation of  $R_{A}(x)$  obtained from expression (6.3) or by applying the B-spline reckures<br>attaining fig. ( $\frac{z^{2}}{4}$ ,  $\frac{1}{4}$ ,  $\frac{z^{$ **6.3.3** The probabilistically symmetric 95 % coverage interval determined on the basis of the GUF is in this case slightly more conservative than that obtained analytically. As for normally distributed quantities, the validation procedure was applied (columns 6–8 of table 6.2). As before,  $n_{\text{dig}} = 2$ ,  $u(y) = 20 \times 10^{-1}$ ,  $c = 20$ ,  $\ell = -1$  and  $\delta = 0.05$ . The endpoint differences  $d_{\text{lo}}$  and  $d_{\text{hi}}$  are larger than for the case of normally distributed quantities (table 6.1). For the first of the two applications of the adaptive Monte Carlo method, the GUF is validated. For the second application, it is not validated, although  $d_{\text{lo}}$  and  $d_{\text{hi}}$  for this application are close to the numerical tolerance  $\delta = 0.05$  (seen if more decimal digits than in table 6.2 are considered). Different validation results such as these are an occasional consequence of the stochastic nature of the Monte Carlo method, especially in a case such as that here.



<span id="page-52-0"></span>Figure 6.2: Counterpart of figure 6.1 for quantities having the same expectations and standard deviations (clause [6.3.2\)](#page-51-0), but distributed as rectangular PDFs. The black curve corresponds to GUF, the scaled histogram represents the MCM result and the red curve (close to the histogram) is the analytic solution

## **6.4 Rectangularly distributed input quantities with different widths**

<span id="page-52-2"></span>**6.4.1** Consider the example of clause 6.3, except that  $X_4$  has a standard deviation of ten rather than unity. Table [6.3](#page-52-1) contains the results obtained.

**6.4.2** The numbers *M* of Monte Carlo trials taken by the adaptive procedure  $(0.03 \times 10^6)$ and  $0.08 \times 10^6$ ) are much smaller than they were for the two previous cases in this example. The main reason is that, in this case,  $\delta = 0.5$ , the numerical tolerance resulting from requesting, as before, two significant decimal digits in  $u(y)$ , is ten times the previous value. Were the previous value to be used, *M* would be of the order of 100 times greater.

Method	$\,M$	$\mathsf{v}$	u(y)	Prob. sym. 95 % cov. int.	$d_{\rm lo}$	$d_{\rm hi}$	GUF validated $(\delta = 0.5)$ ?
GUF		0.0	10.1	$[-19.9, 19.9]$			
<b>MCM</b>	10 <sup>5</sup>	0.0		$10.2$ [-17.0, 17.0]			
<b>MCM</b>	$10^6$	0.0	10.2	$[-17.0, 17.0]$			
<b>MCM</b>	$10^6$	0.0	10.1	$[-17.0, 17.0]$			
<b>Adaptive MCM</b>				$0.03 \times 10^6$ 0.1 10.2 [-17.1, 17.1]	2.8	2.8	N <sub>0</sub>
Adaptive MCM $0.08 \times 10^6$ 0.0 10.1 [-17.0, 17.0]					2.9	2.9	N <sub>0</sub>

<span id="page-52-1"></span>Table 6.3: As table 6.2, except that the fourth input quantity has a standard deviation of ten rather than unity, and no analytic solution is provided (clause [6.4.1\)](#page-52-2)

<span id="page-52-3"></span>**6.4.3** Figure 6.3 shows the (normal) PDF for *Y* (black curve) resulting from the GUF. It also shows a (scaled) histogram obtained from the MCM results. They differ appreciably owing to the dominance of the PDF for  $X_4$ . The PDF for *Y* resembles that for  $X_4$ , but there is an effect in the flanks resulting from the PDFs for the other *X<sup>i</sup>* . The PDF for *Y* , closely resembling the histogram, is also shown as the red curve. It is based on an analytical



solution for the distribution of the sum of *n* non-identically distributed uniform random variables given in [[79](#page-132-2)].

<span id="page-53-0"></span>Figure 6.3: Counterpart of figure [6.2,](#page-52-0) except that the fourth input quantity  $(X_4)$  has a standard deviation of ten rather than unity (clause 6.4.3). The black, bell-shaped curve corresponds to GUF, the (scaled) histogram represents the MCM result, and the red line is the analytical solution. Full and broken vertical lines denote the limits of the 95 % coverage intervals for the GUF and analytical solution, respectively

**6.4.4** The probabilistically symmetric 95 % coverage interval determined on the basis of the GUF in this case is more conservative than that obtained using MCM. Again, the validation procedure was applied (columns 6–8 of table 6.3). Now,  $n_{\text{dig}} = 2$ ,  $u(y) = 10 \times 10^0$ ,  $c = 10$ ,  $\ell = 0$  and  $\delta = 1/2 \times 10^0 = 0.5$ . For the two applications of the adaptive Monte Carlo procedure, the GUF is not validated. For a numerical tolerance of one significant decimal digit in  $u(y)$ , i.e.  $n_{\text{dig}} = 1$ , for which  $\delta = 5$ , the validation status would be positive in both cases, the 95 % coverage intervals all being  $[-2 \times 10^1, 2 \times 10^1]$ .

NOTE The conditions for the Central Limit Theorem (CLT) to apply are not well met in this circumstance [[6,](#page-126-1) clause G.6.5], because of the dominating effect of the rectangular PDF for  $X_4$ . However, because these conditions are often in practice assumed to hold, especially when using proprietary software for uncertainty evaluation, the characterization of *Y* by a normal PDF on the assumption of the applicability of this theorem is made in this clause for comparison purposes.



# **7 Calibration of weights: second-order effects in uncertainty evaluation**

#### **7.1 Preamble**

7 **Calibration of weights:** second-order effects in uncertainty<br>
evaluation<br>
11 Pranmble<br>
11 Pranmble<br>
11 Pranmble<br>
11 Pranmble<br>
11 Decamble<br>
1 This example illustrates the uncertainty evaluation of a mass calibration using the JCGM 100:2008 uncertainty framework (GUF) and the Monte Carlo method (MCM). The evaluation using the measurement model and a simplification are compared. When applying the law of propagation of uncertainty (LPU) as part of the GUF, two variants are used, that is, the commonly used one considering only first-order terms  $\left(\mathrm{GUF}_1\right)$  and the one considering both first- and second-order terms  $(GUF<sub>2</sub>)$ .

#### **7.2 Measurement model**

**7.2.1** Consider the calibration of a weight W, whose density is  $\rho_W$ , against a reference weight R, whose density is  $\rho_R$ . The calibration is performed using a balance operating in air whose density is  $\rho_{\rm a}$ . Both weights have the same nominal mass  $m_{\rm nom}$   $=$  100 g. Since  $\rho_{\rm W}$ and  $\rho_R$  are generally different, it is necessary to account for buoyancy effects [[84](#page-132-3)]. This example appeared originally in [[12,](#page-126-3) clause 9.3].

**7.2.2** Applying Archimedes' principle, the measurement model takes the form

<span id="page-54-0"></span>
$$
m_{\rm W}(1 - \rho_{\rm a}/\rho_{\rm W}) = (m_{\rm R} + \delta m_{\rm R})(1 - \rho_{\rm a}/\rho_{\rm R}),\tag{7.1}
$$

where  $\delta m_R$  is the mass of a small weight of density  $\rho_R$  added to R to balance it with W.

**7.2.3** The conventional mass  $m_{W,c}$  of W is the mass of a (hypothetical) weight of density  $\rho_0 = 8000 \,\text{kg}/\text{m}^3$  that balances W in air at density  $\rho_{a_0} = 1.2 \,\text{kg}/\text{m}^3$  (these assigned value have no associated uncertainty). Thus,

$$
m_{\rm W}(1-\rho_{\rm a_0}/\rho_{\rm W})=m_{\rm W,c}(1-\rho_{\rm a_0}/\rho_0).
$$

**7.2.4** In terms of conventional masses  $m_{W,c}$ ,  $m_{R,c}$  and  $\delta m_{R,c}$ , model [\(7.1\)](#page-54-0) becomes

<span id="page-54-1"></span>
$$
m_{\rm W,c} \frac{1 - \rho_{\rm a}/\rho_{\rm W}}{1 - \rho_{\rm a_0}/\rho_{\rm W}} = (m_{\rm R,c} + \delta m_{\rm R,c}) \frac{1 - \rho_{\rm a}/\rho_{\rm R}}{1 - \rho_{\rm a_0}/\rho_{\rm R}}
$$
(7.2)

or, to an approximation adequate for most practical purposes ( $\rho_W$ ,  $\rho_R \gg \rho_a$ ) [[84](#page-132-3)],

$$
m_{\text{W,c}} = (m_{\text{R,c}} + \delta m_{\text{R,c}}) \bigg[ 1 + (\rho_{\text{a}} - \rho_{\text{a}_0}) \bigg( \frac{1}{\rho_{\text{W}}} - \frac{1}{\rho_{\text{R}}} \bigg) \bigg],
$$

from which, introducing  $\delta m = m_{W,c} - m_{nom}$ , the measurement model used in this example is eventually obtained:

<span id="page-54-2"></span>
$$
\delta m = (m_{\text{R,c}} + \delta m_{\text{R,c}}) \left[ 1 + (\rho_{\text{a}} - \rho_{\text{a}_0}) \left( \frac{1}{\rho_{\text{W}}} - \frac{1}{\rho_{\text{R}}} \right) \right] - m_{\text{nom}}.
$$
 (7.3)

**7.2.5** Applying LPU to either the exact model [\(7.2\)](#page-54-1) or the approximate model [\(7.3\)](#page-54-2) is made difficult by the algebraic complexity of the partial derivatives. It is easier to apply MCM because only model values need be formed.

22.6 The only information concerning  $m_k$ , and  $5m_k$ , is an estimate and an specified<br>tianded uncertainty for each of these quantities. Accordingly, normal distributions are<br>sistend to them, with estimates and standard un **7.2.6** The only information concerning  $m_{\text{R,c}}$  and  $\delta m_{\text{R,c}}$  is an estimate and an associated standard uncertainty for each of these quantities. Accordingly, normal distributions are assigned to them, with estimates and standard uncertainties as expectations and standard deviations, respectively [[12,](#page-126-3) clause 6.4.7.1]. As for quantities  $\rho_{\rm a}$ ,  $\rho_{\rm W}$  and  $\rho_{\rm R}$ , only lower and upper limits are known for each of these quantities. Accordingly, rectangular distributions are assigned to these quantities, with endpoints equal to the known limits [12, clause 6.4.2.1]. Table [7.1](#page-55-0) summarizes the input quantities and their probability density functions (PDFs). In the table, a normal distribution  $N(\mu, \sigma^2)$  is described in terms of expectation  $\mu$ and standard deviation *σ*, and a rectangular distribution R(*a*, *b*) with endpoints *a* and *b* ( $a$  <  $b$ ) in terms of expectation  $(a + b)/2$  and half-width  $(b − a)/2$  or standard deviation  $(b-a)/\sqrt{12}$ .



<span id="page-55-0"></span>

#### **7.3 Uncertainty propagation**

**7.3.1** The GUF that considers only first-order terms  $(GUF_1)$ , applicable to uncorrelated input quantities, corresponds to the following evaluations:

<span id="page-55-3"></span>
$$
y_{\text{GUF1}} = f(\mathbf{x}),\tag{7.4}
$$

$$
u^2(y_{\text{GUF1}}) = \sum_{i=1}^N \left(\frac{\partial f}{\partial X_i}\right)^2 u^2(x_i),\tag{7.5}
$$

where  $\mathbf{x} = [x_1, \dots, x_N]^\top$  and  $\mathbf{y} = f(\mathbf{x})$  denote the vector of estimates of the input quantities  $X_i$ ,  $i = 1,...,N$ , and the corresponding estimate of the measurand *Y*, respectively. Partial derivatives are evaluated at the estimates of the input quantities.

**7.3.2** The GUF that considers first- and second-order terms  $(GUF_2)$ , applicable to uncorrelated and normally distributed input quantities, corresponds to the following evaluations:

<span id="page-55-2"></span><span id="page-55-1"></span>
$$
y_{GUF2} = y_{GUF1} + \frac{1}{2} \sum_{i=1}^{N} \frac{\partial^2 f}{\partial X_i^2} u^2(x_i),
$$
\n
$$
u^2(y_{GUF2}) = u^2(y_{GUF1}) + \sum_{i=1}^{N} \sum_{j=1}^{N} \left[ \frac{1}{2} \left( \frac{\partial^2 f}{\partial X_i \partial X_j} \right)^2 + \frac{\partial f}{\partial X_i} \frac{\partial^3 f}{\partial X_i \partial X_j^2} \right] u^2(x_i) u^2(x_j).
$$
\n(7.7)

Formulæ [\(7.6\)](#page-55-1) and [\(7.7\)](#page-55-2) form a self-consistent pair since, as stated in [72], they are both based on the third-order Taylor expansion of the measurement function  $f(X)$ . The use of the estimate  $(7.4)$  in conjunction with formula  $(7.7)$ , as in [[6,](#page-126-1) clause 5.1.2], represents inconsistent working, since it is biased, not taking higher-order terms into consideration.

**7.3.3** The GUF and the adaptive MCM [[12,](#page-126-3) clause 7.9] were each used to obtain an estimate δ*m* of δ*m*, the associated standard uncertainty *u*(δ*m*), and the shortest 95 % coverage interval for δ*m*. The results obtained are shown in table [7.2.](#page-56-0) There, the GUF1 and GUF2 estimates of δ*m* are equal since the partial derivative for the second-order effect is zero in this instance.

<span id="page-56-0"></span>

Method	$\delta m$ /mg	$u(\delta m)$ /mg	Shortest 95% coverage interval/mg	<b>GUF</b> validated $(\delta = 0.005)$ ?
$GUF_1$	1.2340	0.0539	[1.1285, 1.3395]	No.
<b>MCM</b>	1.2340	0.0755	[1.0843, 1.3838]	
GUF <sub>2</sub>	1.2340	0.0750	[1.0870, 1.3810]	Yes

Table 7.2: Results of the calculation stage for the mass calibration model [\(7.3\)](#page-54-2), δ*m*

**7.3.4** A total of  $0.72 \times 10^6$  trials were taken by the adaptive MCM [[12,](#page-126-3) subclause 7.9] with the use of a numerical tolerance of  $\delta/5$  [12, subclause 8.2] with  $\delta$  set for the case where one significant decimal digit in  $u(\delta m)$  is regarded as meaningful.

<span id="page-56-1"></span>**7.3.5** Figure [7.1](#page-57-0) shows the approximations to the PDF for  $\delta m$  obtained from GUF<sub>1</sub> and MCM. The continuous curves represent a normal PDF with parameters given by the GUF. The inner pair of (broken) vertical lines indicates the shortest 95 % coverage interval for δ*m* based on this PDF.

The histogram is the scaled frequency distribution obtained using MCM as an approximation to the normal distribution provided by the GUF. The outer pair of (continuous) vertical lines indicates the endpoints of the shortest 95 % coverage interval for δ*m* based on the discrete representation of the distribution function determined as in [[12,](#page-126-3) subclause 7.5].

**7.3.6** Although GUF<sub>1</sub> and MCM both give estimates of  $\delta m$  in good agreement, the associated standard uncertainties are noticeably different. The value of  $u(\delta m)$  returned by either MCM or GUF $_2$  is 40 % larger than that returned by GUF $_1.$  The latter method is therefore understating the uncertainty associated with δ*m*.

Formule (7.6) and [7.7) from a self-consisten pair sine, as stated in [72], they are hold in the self-consistent pair of the measurement function  $f(X)$ , the self-consistent pair sine estimate (7.4) in consistent pair of t **7.3.7** Table 7.2 also shows in the right-most two columns the results of applying the JCGM 101:2008 validation procedure [[12,](#page-126-3) clause 8] in the case where one significant decimal digit in *u*(δ*m*) is regarded as meaningful. Since a numerical tolerance of one significant decimal digit in  $u(\widehat{\delta m})$  is required,  $n_{\text{dig}} = 1$ . Hence,  $u(\widehat{\delta m}) = 0.08$  mg =  $8 \times 10^{-2}$  mg, and so the *c* in [12, subclause 7.9 item a)] equals 8 and  $\ell = -2$ . Thus  $\delta = 1/2 \times 10^{-2} = 0.005$ . Whether the results were validated to one significant decimal digit in  $u(\delta m)$  is indicated in the final column of the table. Thus, if only first-order terms are considered, the application of GUF provides inadequate results. If, however, higher-order terms are considered [[6,](#page-126-1) clause 5.1.2 note], the GUF becomes valid. Thus, the non-linearity of the model is such that accounting only for first-order terms is inadequate and higher-order terms cannot be neglected.



<span id="page-57-0"></span>Figure 7.1: Approximations to the PDF for the output quantity δ*m* obtained using GUF<sup>1</sup> (sharply peaked curve),  $GUF_2$  (more rounded curve) and MCM scaled histogram, subclause [7.3.5\)](#page-56-1)

#### **7.4 Uncertainty budget**

#### **7.4.1 GUF**

In many measurement contexts it is common practice to list the uncertainty components  $u_i(y) = |c_i| u(x_i)$ ,  $i = 1, ..., N$ , where  $c_i$  is the *i*th sensitivity coefficient and  $u(x_i)$  the standard uncertainty associated with the *i*th input estimate *x<sup>i</sup>* , contributing to the standard uncertainty  $u(y)$ . Usually these are presented in a table (such as table 7.3), the "uncertainty budget". This practice may be useful to identify the dominant terms contributing to  $u(y)$  associated with the estimate of the output quantity.

Table [7.3](#page-57-1) lists the partial derivatives of first order for the model [\(7.3\)](#page-54-2) with respect to the input quantities together with the sensitivity coefficients, namely, these derivatives evaluated at the estimates of the input quantities.

<span id="page-57-1"></span>

Since the uncertainty associated with the output quantity is the sum of squared individual uncertainty contributions,  $u_i(y)$ , the squared uncertainties of the two input quantities  $m_{\text{R,c}}$  and  $\delta m_{\text{R,c}}$  contribute to the squared uncertainty  $u^2(y)$  in the respective proportions  $0.050^2/(0.050^2+0.020^2) = 86\%$  and  $0.020^2/(0.050^2+0.020^2) = 14\%$ .

<span id="page-58-0"></span>Table 7.4: Uncertainty budget – GUF $_2$  (only terms having non-zero partial derivatives are shown)

	The uncertainty budget can also include higher-order uncertainty contributions in accor- dance with expression (7.7). Table 7.4 shows that the interaction of variables $\rho_a$ and $\rho_w$ has significant impact on the overall uncertainty estimate $u(y)$ .	
shown)	Table 7.4: Uncertainty budget $-GUF2$ (only terms having non-zero partial derivatives are	
Input quantity $X_i$	Partial derivative $c_i = \partial Y / \partial X_i$	Rel. contrib./% $c_i^2 u^2(x_i)/u^2(y)$
$m_{R,c}$ $\delta m_{\rm R,c}$	$1 + (\rho_a - \rho_{a_0})(1/\rho_W - 1/\rho_R) = 1$ $1 + (\rho_a - \rho_{a_0})(1/\rho_W - 1/\rho_R) = 1$	44 7
Input quantities $X_i, X_j$	Partial derivative $c_{ij} = \frac{\partial^2 Y}{\partial X_i \partial X_j}$	Rel. contrib./% $c_{i}^{2}u^{2}(x_{i})u^{2}(x_{j})/u^{2}(y)$
$\rho_{\rm a}, \rho_{\rm W}$ $\rho_{\rm a}, \rho_{\rm R}$	$-(m_{R,c} + \delta m_{R,c})/\rho_W^2 = -0.00156 \,\text{m}^4/\text{mg}$ $+(m_{R,c} + \delta m_{R,c})/\rho_R^2 = +0.00156 \text{ m}^4/\text{mg}$	49 0
	$u(y) = 0.0750$ mg	100
	fixed at their estimates, MCM can be used to provide the PDF for the output quantity for the model having just that input quantity as a variable [6]. The quotient of the standard deviation of the resulting model values and the standard uncertainty associated with the estimate of the relevant input quantity can be taken as a sensitivity coefficient. This quo- tient corresponds to that which would be obtained by taking all higher-order terms in the Taylor series expansion of the model into account. This approach may be viewed as a generalization of the approximate partial-derivative formula in [6, clause 5.1.3 note 1]. Both the sensitivity coefficients and the contributions for each input quantity to the uncer- tainty associated with the estimate of the output quantity will in general differ from those	
obtained with JCGM 100:2008 [6].		

#### **7.4.2 MCM**

# **8 Gauge block calibration**

## **8.1 Preamble**

**3 G[a](#page-126-1)uge block calibration**<br> **3.1 Prea[m](#page-130-5)ble**<br> **3.1 Preamble**<br> **3.1 Security** stems the uncertainty evaluation for the calibration of a gauge block units<br>
tractice standard by applying the JCGM 100-2008 uncertainty framewor This example shows the uncertainty evaluation for the calibration of a gauge block using a reference standard by applying the JCGM 100:2008 uncertainty framework (GUF) [6, clause 5] and the Monte Carlo method (MCM) [[12,](#page-126-3) clause 7]. The example starts with the transformation of a model for the difference between the length of the blocks, the quantity actually measured, into an explicit measurement model. Since some of the input quantities in that model are correlated, a re-parametrization is used to remove that correlation. Probability density functions are assigned to the transformed quantities based on the metrologist's knowledge. Using those PDFs, GUF and MCM are applied and the results discussed.

### **8.2 Measurement model**

**8.2.1** The length of a gauge block is determined by comparing it with a known reference standard having the same nominal length  $L_{\text{nom}} = 50 \text{ mm}$ . The direct output of the comparison of two gauge blocks is the difference *d* in their lengths given by

<span id="page-59-0"></span>
$$
d = L(1 + \alpha \theta) - L_s(1 + \alpha_s \theta_s),
$$
\n(8.1)

where

*L* denotes the length (at 20 °C temperature [58]) of the gauge block being calibrated,

 $L<sub>s</sub>$  denotes the length of the reference standard, also at 20  $^{\circ}$ C, as given in its calibration certificate,

 $\alpha$  and  $\alpha$ <sub>s</sub> are the coefficients of thermal expansion for the two blocks,

*<sup>θ</sup>* and *<sup>θ</sup>*<sup>s</sup> are the deviations in temperature of the two gauge blocks from the 20 ◦<sup>C</sup> reference temperature.

(JCGM 100:2008 refers to a gauge block as an end gauge and the symbol *L* for the length of a gauge block is used in this document in place of the symbol *ℓ* used in JCGM 100:2008 for that quantity.)

**8.2.2** From expression (8.1), the output quantity *L* is given by

<span id="page-59-2"></span><span id="page-59-1"></span>
$$
L = \frac{L_s(1 + \alpha_s \theta_s) + d}{1 + \alpha \theta}.
$$
\n(8.2)

By using  $1/(1+x) \approx 1-x$  when |*x*| is small compared with unity, model [\(8.2\)](#page-59-1) becomes

$$
L = L_{s} + d + L_{s}(\alpha_{s}\theta_{s} - \alpha\theta) - d\alpha\theta - L_{s}\alpha_{s}\theta_{s}\alpha\theta.
$$
\n(8.3)

The term  $L_1 \alpha_s \theta_s \alpha \theta$  is a second-order term and can be safely neglected. The term  $d\alpha\theta$ is the thermal-expansion correction for the difference *d* between the two gauge blocks. Even considering an extreme situation (say,  $d = 1 \,\mu$ m and  $\theta = 3 \degree C$ ), both the value and the associated standard uncertainty lie well below a negligible 0.1 nm. Model [\(8.3\)](#page-59-2) thus reduces to

<span id="page-59-3"></span>
$$
L = Ls + d + Ls(\alphas \thetas - \alpha \theta),
$$
\n(8.4)

an approximation adequate for most practical purposes.

<span id="page-60-2"></span>If the difference in temperature between the gauge block being calibrated and the reference standard is written as  $\delta\theta = \theta - \theta_s$  and the difference in their thermal expansion coefficients as  $\delta \alpha = \alpha - \alpha_s$ , model [\(8.4\)](#page-59-3) becomes

$$
L = L_{\rm s} + d - L_{\rm s} (\delta \alpha \theta + \delta \theta \alpha_{\rm s}).
$$

This re-parametrization [[10,](#page-126-5) clause 8.3.2] avoids the introduction of covariances between *θ* and  $\theta_S$  and between  $\alpha$  and  $\alpha_S$ .

**8.2.3** The difference *d* in the lengths of the gauge block being calibrated and the reference standard is determined as the average of a series of five indications, obtained independently, of the difference using a calibrated comparator. The difference *d* can be expressed as

<span id="page-60-0"></span>
$$
d = D + d_1 + d_2,\tag{8.6}
$$

where *D* denotes the observed length difference, and  $d_1$  and  $d_2$  are additive corrections modelling, respectively, the random and systematic effects associated with using the comparator [[10,](#page-126-5) clause 10.4.3].

**8.2.4** The quantity *θ*, representing deviation of the temperature from 20 ◦C of the gauge block being calibrated, can be expressed as

<span id="page-60-1"></span>
$$
\theta = \theta_0 + \Delta,\tag{8.7}
$$

where  $\theta_0$  represents the average temperature deviation of the gauge block from 20  $^{\circ}$ C and  $\varDelta$  describes a cyclic variation of the temperature deviation from  $\theta_0.$ 

<span id="page-60-3"></span>**8.2.5** Substituting expressions (8.6) and [\(8.7\)](#page-60-1) into model [\(8.5\)](#page-60-2) and introducing  $\delta L = L - L_{\text{nom}}$ gives

$$
\delta L = L_{\rm s} + D + d_1 + d_2 - L_{\rm s} [\delta \alpha (\theta_0 + \Delta) + \delta \theta \alpha_{\rm s}] - L_{\rm nom}
$$
\n(8.8)

as the measurement model employed here.

an approximation adequate for most practical purposes.<br>
Jf the difference in term<br>cataristic between the gauge block being calibrated and the reference<br>
is standard is written as  $6\theta = 0 - \theta$ , and the difference in their t **8.2.6** This measurement problem is treated in terms of model [\(8.8\)](#page-60-3) with output quantity δ*L* and input quantities  $L_s$ , *D*,  $d_1$ ,  $d_2$ ,  $\alpha_s$ ,  $\theta_0$ ,  $\Delta$ , δ $\alpha$  and δ $\theta$ . This treatment differs from JCGM 100:2008 example H.1 where the models [\(8.6\)](#page-60-0) and [\(8.7\)](#page-60-1) are regarded as sub-models to model [\(8.5\)](#page-60-2). The approach here avoids having to use the results obtained from MCM applied to the sub-models [\(8.6\)](#page-60-0) and [\(8.7\)](#page-60-1) to provide information about the distributions for the input quantities  $d$  and  $\theta$  in model [\(8.5\)](#page-60-2).

 $(8.5)$ 

## **8.3 GUF**

**8.3.1** The application of the GUF is based on the Taylor series approximation to the model [\(8.8\)](#page-60-3), the use of the Welch-Satterthwaite formula to evaluate an effective degrees of freedom (rounded towards zero) associated with the uncertainty obtained from the law of propagation of uncertainty, and assigning a scaled and shifted *t*-distribution with the above degrees of freedom to the output quantity. For a detailed account of the GUF approach, including the determination of the sensitivity coefficients, the assignment of estimates and associated uncertainties, and the consideration of second-order terms, see [6, Annex H.1].

## **8.4 MCM**

The application of MCM requires sampling from a rectangular distribution, normal distribution, *t*-distribution, U-shaped distribution, and rectangular distribution with inexactly prescribed limits, and implements adaptive MCM [12, clause 7.9] with a numerical tolerance ( $\delta$  = 0.5) set to deliver  $n_{\text{dig}}$  = 2 significant decimal digits in the standard uncertainty.

## **8.4.1 Assignment of PDFs**

**8.4.1** In the following clauses the available information about each input quantity in model [\(8.8\)](#page-60-3) is provided. For each item of information, the JCGM 100:2008 [[6](#page-126-1)] subclause from which the item is extracted is identified. Also provided is an interpretation of the information in terms of an assignment of a distribution to the quantity. Table [8.1](#page-63-0) summarizes the assignments made.

3.3 **CUF**<br>
3.3 **CUF**<br>
3.3 **CUF**<br>
3.3 **A** The application of the GUF is base[d](#page-126-1) on the Taylor series approximation in the<br>
model (8.8), the use of the Welch-Satrenthwaite formula to evaluate an effective degrees of<br>
recordin **8.4.2** The calibration certificate for the reference standard gives  $\hat{L}_s = 50.000623$  mm as its length at 20 °C [[6,](#page-126-1) Annex H.1.5]. It gives  $U(\hat{L}_s) = 75$  nm as the expanded uncertainty of the reference standard length and states that it was obtained using a coverage factor of *k* = 3 [[6,](#page-126-1) Annex H.1.3.1]. The certificate states that the effective degrees of freedom associated with the combined standard uncertainty, from which the quoted expanded uncertainty was obtained, is  $v_{\text{eff}}(u(L_s)) = 18$  [[6,](#page-126-1) Annex H.1.6]. Thus, a scaled and shifted *t*distribution  $t_{\nu}(\mu, \sigma^2)$  is assigned to  $\widehat{L}_s$  [[12,](#page-126-3) clause 6.4.9.7]. Here  $\mu$ ,  $\sigma$  and  $\nu$  are the location, scale and degrees of freedom, respectively, with values

$$
\mu = 50000623 \text{ nm}, \qquad \sigma = \frac{U}{k} = \frac{75}{3} \text{ nm} = 25 \text{ nm}, \qquad \nu = 18.
$$

**8.4.3** The average of the five indications of the length difference between the gauge block being calibrated and the reference standard is 215 nm [[6,](#page-126-1) Annex H.1.5]. The pooled experimental standard deviation characterizing the comparison of *L* and *L*<sup>s</sup> was determined from 25 indications, obtained independently, of the length difference of two standard gauge blocks, and equal 13 nm [[6,](#page-126-1) Annex H.1.3.2]. Thus, a scaled and shifted *t*distribution  $t_{\nu}(\mu, \sigma^2)$  is assigned to *D* [[12,](#page-126-3) clause 6.4.9.2], with

$$
\mu = 215 \text{ nm}, \qquad \sigma = \frac{13}{\sqrt{5}} \text{ nm} = 5.8 \text{ nm}, \qquad \nu = 24.
$$

**8.4.4** According to the calibration certificate of the comparator used to compare *L* with *L*<sup>s</sup> , the associated uncertainty due to random effects is 10 nm for a coverage probability of 95 % and is obtained from six indications, obtained independently [[6](#page-126-1)][Annex H.1.3.2]. Thus, a scaled and shifted *t*-distribution  $t_y(\mu, \sigma^2)$  is assigned to  $d_1$  [[12,](#page-126-3) clause 6.4.9.2], with

$$
\mu = 0 \text{ nm}, \qquad \sigma = \frac{U}{k} = \frac{10}{2.57} \text{ nm} = 3.9 \text{ nm}, \qquad \nu = 5.
$$

Here, *k* is obtained from table G.2 of JCGM 100:2008 with  $\nu = 5$  degrees of freedom and  $p = 0.975$ .

**8.4.5** The uncertainty of the comparator due to systematic effects is given in the certificate as 20 nm at the "three sigma level" [6, Annex H.1.3.2]. This uncertainty may be assumed to be reliable to 25%, and thus the degrees of freedom is  $v_{\text{eff}}(u(d_2)) = 8$ [[6,](#page-126-1) Annex H.1.6]. Thus, a scaled and shifted *t* distribution  $t_{\nu}(\mu, \sigma^2)$  is assigned to  $d_2$  [[12,](#page-126-3) clause 6.4.9.2], with

$$
\mu = 0 \text{ nm}, \qquad \sigma = \frac{U}{k} = \frac{20}{3} \text{ nm} = 6.7 \text{ nm}, \qquad \nu = 8.
$$

**8.4.6** The coefficient of thermal expansion of the length reference standard is given as  $\hat{\alpha}_s = 11.5 \times 10^{-6} \,^{\circ}\text{C}^{-1}$  with possible values of this quantity represented by a rectangular distribution with limits  $\pm 2 \times 10^{-6} \,^{\circ}\text{C}^{-1}$  [6] Annoy H 1.2.2.1. Thus, a rectangular distri distribution with limits  $\pm 2 \times 10^{-6}$  °C<sup>-1</sup> [6, Annex H.1.3.3]. Thus, a rectangular distribution R(*a*, *b*) is assigned to  $\alpha_s$  [12, clause 6.4.2] with limits  $a = 9.5 \times 10^{-6}$  °C<sup>-1</sup> and  $b = 13.5 \times 10^{-6} \,^{\circ} \text{C}^{-1}.$ 

**8.4.7** The temperature of the test bed is reported as  $(19.9 \pm 0.5)$  °C. The average temperature deviation  $\widehat{\theta}_0 = -0.1$  °C is reported as having an associated standard uncertainty due to the uncertainty associated with the average temperature of the test bed of  $u(\hat{\theta}_0) = 0.2 \degree C$  [[6,](#page-126-1) Annex H.1.3.4]. Thus, a normal distribution N( $\mu, \sigma^2$ ) is assigned to  $\theta_0$  [[12,](#page-126-3) clause 6.4.7], with  $\mu = -0.1$  °C and  $\sigma = 0.2$  °C. Since there is no information about the source of the evaluation of the uncertainty, a normal distribution is assigned.

**8.4.8** The temperature of the test bed is reported as  $(19.9 \pm 0.5)$  °C. The stated maximum offset of −0.5 ◦Cfor *∆* is said to represent the amplitude of an approximately cyclical variation of temperature under a thermostatic system. The cyclic variation of temperature results in a U-shaped (arc sine) distribution [[6,](#page-126-1) Annex H.1.3.4]. Thus, an arc sine distribution U(*a*, *b*) is assigned to  $\overline{\Delta}$  [12, clause 6.4.6], with limits  $a = -0.5$  °C and  $b = 0.5$  °C.

and is obtained from air indications, chuolerly diadependently [6]].Annex H.1.3.2]. Thus, accords diadependently  $\mu = 0$ nm,  $\sigma = \frac{U}{k} = \frac{10}{2.57}$ nm = 3.9 mm,  $V = 5$ . chuos 6.4.9.2 h,with excels of readom  $\mu = 0$ nm,  $\sigma$ **8.4.9** The estimated bounds on the variability of  $\delta \alpha$  are  $\pm 1.0 \times 10^{-6}$  °C<sup>-1</sup> with an equal probability of δ*α* having any value within those bounds [[6,](#page-126-1) Annex H.1.3.5]. These bounds are deemed to be reliable to 10%, giving  $\nu(u(\delta \alpha)) = 50$  [[6,](#page-126-1) Annex H.1.6]. Thus, a rectangular distribution with inexactly prescribed limits is assigned to δ*α* [[12,](#page-126-3) clause 6.4.3], with parameters  $a = -1.0 \times 10^{-6} \degree \text{C}^{-1}$ ,  $b = 1.0 \times 10^{-6} \degree \text{C}^{-1}$ , and  $d = 0.1 \times 10^{-6} \degree \text{C}^{-1}$ . The stated reliability of 10 % on the estimated bounds provides the basis for this value of the parameter *d*.

**8.4.10** The reference standard and the gauge block being calibrated are expected to be at the same temperature, but the temperature difference  $\delta\theta$  could lie with equal probability anywhere in the estimated interval <sup>−</sup>0.05 ◦C to 0.05 ◦<sup>C</sup> [[6,](#page-126-1) Annex H.1.3.6]. This difference is believed to be reliable only to 50%, giving  $v(u(\delta \tilde{\theta})) = 2 [6,$  $v(u(\delta \tilde{\theta})) = 2 [6,$  Annex H.1.6]. Thus, a rectangular distribution with inexactly prescribed limits is assigned to  $\delta\theta$  [[12,](#page-126-3) clause 6.4.3],

Quantity PDF d b $\sigma$ $\boldsymbol{\nu}$ a $\mu$ $t_{\nu}(\mu, \sigma^2)$ 50000623 $L_{\rm s}/\rm{nm}$ 25 18 $t_{\nu}(\mu, \sigma^2)$ D/nm 215 5.8 24 $t_{\nu}(\mu, \sigma^2)$ $d_1/nm$ 3.9 5 0 $d_2/nm$ $t_{\nu}(\mu, \sigma^2)$ 6.7 8 0 $\alpha_{\rm s}/(10^{-6\,\circ}{\rm C}^{-1})$ 13.5 R(a, b) 9.5 $N(\mu, \sigma^2)$ $\theta_0$ /°C 0.2 $-0.1$ $\Delta$ /°C U(a, b) 0.5 0.5 $\delta \alpha/(10^{-6}$ °C <sup>-1</sup> ) 1.0 CTrap(a, b, d) 1.0 0.1 $-0.050$ $\delta\theta$ /°C CTrap(a, b, d) 0.050 0.025 Results The estimate $\delta L$ is the sum of three terms of model (8.8), $L_s$ , D and $L_{\text{nom}}$ . The The distribution obtained from the GUF is a t distribution with $\nu = 16$ degrees
3.5 <b>3.5.1</b> Table 8.2 gives the results obtained using the information summarized in table 8.1. Figure 8.1 shows the PDFs for $\delta L$ obtained from the application of the GUF (red solid curve) ) ind MCM (blue scaled histogram). For the latter, $1.3\times10^6$ trials were taken by the adaptive Monte Carlo procedure. 3.5.2 emaining terms have estimates equal to zero, thus not contributing to the measurand estimate. Yet, the associated uncertainties are different from zero and do contribute to the issociated uncertainty. See also [10, clauses 10.3.4 and 10.4.3]. 3.5.3 of freedom. The endpoints of the shortest 99% coverage intervals for $\delta L$ obtained from he PDFs are indicated as (broken, red) vertical lines (obtained from the GUF) and (con-
inuous, blue) vertical lines (obtained from MCM).
Table 8.2: Results for model (8.8) using the information summarized in table 8.1
$\widehat{\delta L}$ $u(\widehat{\delta L})$ Method Shortest 99 % coverage
interval for $\delta L / nm$ /nm /nm
<b>GUF</b> 838 [745, 931] 32
<b>MCM</b> [745, 932] 838 36

<span id="page-63-0"></span>Table 8.1: PDFs assigned to input quantities for the gauge block model (8.8) on the basis of available information

## **8.5 Results**

<span id="page-63-1"></span>





<span id="page-64-0"></span>Figure 8.1: PDF for δ*L* obtained using the GUF (red curve) and scaled histogram from MCM for the approximate model (8.8) using the information summarized in table [8.1](#page-63-0)



# **9 Calibration of a torque measuring system**

## **9.1 Preamble**

**Calibration of a torque measuring system**<br> **1.1 Preamble**<br>
This example and a simple information of a torque-measuring somen assign restriction restrictions profile and different torque-measuring somen assign reference s This example addresses the straight-line calibration of a torque-measuring sensor against a reference system using measurements taken at different torque values. The goal is to determine a linear relationship that relates the results of the torque measuring sensor with those of the reference system. Two approaches are provided: (i) ordinary and weighted least-squares estimation (see, for example, ISO/TS 28037 [[61](#page-130-6)]), accompanied by an uncertainty evaluation based on the law of propagation of uncertainty (LPU) [[6](#page-126-1)], and (ii) a statistical approach applying Bayesian inference (compare, for example, [[48](#page-129-1)]). Explicit expressions are given for the Bayesian uncertainty analysis, which simplifies its application.

## **9.2 Experimental**

Measurements have been carried out at different values of torque by the sensor and the reference system. Table [9.1](#page-65-0) shows the data, where repeated measurement results of the sensor have been summarized through their means and standard deviations, respectively. For further details about the measurements the reader is referred to [112].

<span id="page-65-0"></span>Table 9.1: Summary statistics for part of the measurement data given in the guideline [[112](#page-134-1)][table B6, page 43]. The summary statistics include mean  $y_i$  and standard deviations  $S_i$  of the  $n_i$  measurement results  $y_i$  of the sensor at the *i*th torque level;  $x_i$  denotes the corresponding value given by the reference system. These summary statistics are available online in repository [[73](#page-131-1)]



# **9.3 Specification of the measurand**

Let *X* denote the applied torque, in what follows called stimulus, and *Y* the corresponding quantity measured by the considered sensor, denoted below as response. The linear relation

<span id="page-65-1"></span>
$$
Y = \beta X \tag{9.1}
$$

is assumed to model the relationship between the measured responses of the sensor and the applied stimulus. Model [\(9.1\)](#page-65-1) represents a straight line with zero offset. Such a straight line has been chosen for physical reasons, which is supported by the observed data. The measurand is the slope parameter  $β$  in the particular straight line model [\(9.1\)](#page-65-1). The input quantities are  $Y_1, \ldots, Y_p$  (with  $p = 8$  in this example), which correspond to the

measured values at the considered torques  $X_1,\ldots,X_p.$  The variability associated with the measured values  $x_1, \ldots, x_p$  from the reference device are considered sufficiently small to be neglected.

### **9.4 Measurement model**

#### **9.4.1 General**

measure[d](#page-126-1) values as the considered torque  $X_1, \ldots, X_p$ . The variability associated with the measured values  $x_1, \ldots, x_p$  from the reference device are considered antificiently ganding<br>be neglected.<br> **9.4 Measurement model 9.4.1.1** The uncertainty evaluations presented in this example are based on two different models. The uncertainty evaluation following the JCGM 100:2008 [6] in connection with ordinary and weighted least-squares estimation is based on a measurement model in which the measurand is represented as a function of the input quantities. An estimate of the measurand is then obtained by evaluating this measurement model using the estimates of the input quantities. The uncertainty associated with the resulting estimate for the measurand results from a propagation of the uncertainties associated with the estimates of the input quantities through this measurement model.

**9.4.1.2** The Bayesian inference is based on a statistical model for the observed data, and the measurand enters as one of the parameters of the statistical model. Bayesian inference can account for prior knowledge about the measurand, resulting in a probability distribution for the measurand, which can be viewed as the final complete result. Mean and standard deviation of that distribution can be taken as an estimate and standard uncertainty for the measurand. Bayesian uncertainty analysis can be viewed as being reached through Bayesian inference, rather than by a process of propagating input uncertainties through a measurement model in the sense of JCGM 100:2008 [6].

### **9.4.2 Ordinary and weighted least-squares**

**9.4.2.1** The application of weighted least squares (WLS) estimation determines an estimate  $\hat{\beta}$  for the measurand by minimizing

<span id="page-66-0"></span>
$$
Q = \sum_{i=1}^{p} \sum_{j=1}^{n_i} W_i (y_{ij} - \beta x_i)^2 = \sum_{i=1}^{p} W_i [n_i (y_i - \beta x_i)^2 + (n_i - 1) S_i^2]
$$
(9.2)

with respect to *β*. In expression (9.2),  $y_i$  and  $S_i$  represent mean and standard deviation of the repeated measurement results  $y_{ij}$ ,  $j = 1, \ldots, n_i$ , of the sensor at the *i*th stimulus  $x_i$ , and *W<sub>i</sub>* denote some weights,  $i = 1, ..., p$ . The solution  $\widehat{\beta}$  to this minimization problem is

<span id="page-66-1"></span>
$$
\widehat{\beta} = \left(\sum_{i=1}^{p} n_i W_i x_i^2\right)^{-1} \sum_{i=1}^{p} n_i W_i x_i y_i.
$$
\n(9.3)

The measurement model will be defined as

<span id="page-66-2"></span>
$$
\beta = \left(\sum_{i=1}^{p} n_i W_i x_i^2\right)^{-1} \sum_{i=1}^{p} n_i W_i x_i Y_i, \qquad (9.4)
$$

that is, by replacing the estimates  $y_i$  in expression [\(9.3\)](#page-66-1) with corresponding quantities  $Y_i$ . Note that a measurement model in the sense of JCGM 100:2008 [[6](#page-126-1)] is always a model between quantities. Since the estimates  $x_i$  of the stimulus are treated as being exact, the actual quantity  $X_i$  in [\(9.4\)](#page-66-2) has already been replaced with the known values  $x_i$  in this example. Ordinary least squares (OLS) estimation is obtained by choosing weights  $W_i = 1$ for all observations  $i = 1, \ldots, p$ .

#### **9.4.3 Statistical model**

A statistical model specifies the distribution from which the observed data is taken as a realization. The subsequent statistical model assumes that all individual observations  $y_{ij}$ , for  $j = 1, \ldots, n_i$ ,  $i = 1, \ldots, p$ , are realizations of independently and normally distributed random variables  $Y_{ij}$  with means equal to  $\beta x_i$  and variances  $\sigma_i^2$ , that is,

$$
y_{ij}|\beta, \sigma_i^2 \sim N(\beta x_i, \sigma_i^2).
$$

<span id="page-67-2"></span><span id="page-67-0"></span> $(9.5)$ 

In expression [\(9.5\)](#page-67-0), the  $x_i$  denote the known stimuli and  $\sigma_i^2$  are the variances to be inferred. The likelihood function is the probability for the observed data viewed as a function of the measurement model parameters. For the statistical measurement model [\(9.5\)](#page-67-0), the likelihood function is

$$
\ell(\beta, \sigma^2 | \text{data}) \propto \prod_{i=1}^p (\sigma_i^2)^{-n_i/2} \exp\left\{-\frac{1}{2\sigma_i^2} [(n_i-1)S_i^2 + n_i(y_i-\beta x_i)^2] \right\},\tag{9.6}
$$

where  $\sigma^2 = (\sigma_1^2, \ldots, \sigma_p^2)^\top$ ,  $y_i$  and  $S_i$  denote the mean and standard deviation of  $y_{ij}$ ,  $j = 1, \ldots, n_i$ , and "data" summarizes the sufficient statistics  $y_1, \ldots, y_p, S_1, \ldots, S_p$  of the data.

#### **9.5 Uncertainty evaluation**

#### **9.5.1 Law of propagation of uncertainty**

1.4.3 Statistical model<br>
statistical model specifies the distribution from which the observed dua is taken as a<br>
statistical model specifies the distribution from which the observed dua is taken as a<br>
verified model spec **9.5.1.1** The measurement model (9.4) contains input quantities  $Y_i$ ,  $i = 1, \ldots, p$ . For each of these input quantities, a series of repeated measurement results  $y_{ij}$ ,  $j = 1, \ldots, n_i$ , each of these input quantities, a series of repeated ineasurement results  $y_{ij}$ ,  $j = 1, \ldots, n_i$ , is available. In following JCGM 100:2008, the mean and scaled standard deviation  $S_i / \sqrt{n_i}$ are taken, respectively, as an estimate  $y_i$  of  $Y_i$  with associated standard uncertainty  $u(y_i)$ . The estimates and standard uncertainties for the input quantities are listed in table [9.2.](#page-67-1)

<span id="page-67-1"></span>Table 9.2: The estimates  $y_i$  and their associated standard uncertainties  $u(y_i)$  given the data in table [9.1](#page-65-0) and measurement model [\(9.4\)](#page-66-2)

Input $Y_1$ quantity	$Y_2$		$Y_{\epsilon}$	ΥΩ
$u(y_i)/N$ m 0.0022 0.0021 0.0017 0.0017 0.0040 0.0033 0.0049 0.0219	$y_i/N$ m 0.0950 0.1966 0.3016 0.4983 1.0083 3.0266 4.0466 5.0666			

**9.5.1.2** According to JCGM 100:2008 [[6](#page-126-1)], the estimate  $\hat{\beta}$  of  $\beta$  is taken as the value given by the measurement model [\(9.4\)](#page-66-2) when inserting the estimates from table [9.2](#page-67-1) for the input quantities, that is,

$$
\widehat{\beta} = \left(\sum_{i=1}^{p} n_i W_i x_i^2\right)^{-1} \sum_{i=1}^{p} n_i W_i x_i y_i.
$$
\n(9.7)

The associated variance (squared standard uncertainty) is

$$
u^{2}(\widehat{\beta}) = \sum_{i=1}^{p} \left( \frac{\partial \beta}{\partial Y_{i}} \bigg|_{Y_{i}=y_{i}} \right)^{2} u^{2}(y_{i}) = \frac{\sum_{i=1}^{p} (n_{i}W_{i}x_{i})^{2} u^{2}(y_{i})}{\left(\sum_{i=1}^{p} n_{i}W_{i}x_{i}^{2}\right)^{2}}.
$$
\n(9.8)

**9.5.1.3** The WLS estimate assigns weights as  $n_iW_i = 1/u^2(y_i)$ , whereas OLS estimate corresponds to setting  $W_i = 1$  for all *i*.

**9.5.1.4** Assuming a *t*-distribution with degrees of freedom  $\nu = n - 1$  for  $\beta$ , a 95% coverage interval is  $\beta \pm t_{0.975.7}u(\beta)$ .

#### **9.5.2 Bayesian uncertainty analysis**

9.5.1.3 The WiS esti[m](#page-67-2)ate assigns weights as  $a_1W_t = 1/x^2(y_1)$ , whereas OLS estimate<br>corresponds to setting  $W_t = 1/x^2(y_1)$ , whereas OLS estimate<br>orresponds to setting  $W_t = 1$  for all.<br> [C](#page-68-2)omparison in the commitment of the c **9.5.2.1** In a Bayesian inference, prior knowledge about the measurand (and other unknowns) is combined with the information contained in the data through application of Bayes' theorem. The result is the posterior distribution that summarizes the knowledge about the measurand (and other unknowns in the statistical model (9.5)), conditional on the observed data. In our case, the posterior is given through the following probability density function (PDF):

<span id="page-68-0"></span>
$$
p(\beta, \sigma^2 | \text{data}) \propto \pi(\beta, \sigma^2) \ell(\beta, \sigma^2 | \text{data}),
$$
\n(9.9)

where  $\ell(\beta,\sigma^2|{\rm data})$  is the likelihood function (9.6) for the assumed statistical model [\(9.5\)](#page-67-0), and  $\pi(\beta,\sigma^2)$  is the adopted prior for the model parameters  $\beta$  and  $\sigma^2$  where  $\sigma^2$  denotes a vector  $(\sigma_1^2, \ldots, \sigma_p^2)$ .

**9.5.2.2** Once the posterior distribution of all model parameters is obtained, expression [\(9.9\)](#page-68-0), the posterior  $p(\beta|\text{data})$  for the parameter of interest (the measurand) is obtained through marginalization:

<span id="page-68-1"></span>
$$
p(\beta|\text{data}) = \int_0^\infty \cdots \int_0^\infty p(\beta, \sigma^2|\text{data}) \, d\sigma_1^2 \dots d\sigma_p^2.
$$
 (9.10)

The marginal posterior [\(9.10\)](#page-68-1) is a PDF that constitutes the complete Bayesian uncertainty analysis for the measurand. Summary statistics of this PDF may be sufficient in many cases, and the posterior mean

<span id="page-68-3"></span>
$$
\widehat{\beta} = \int_{-\infty}^{\infty} p(\beta | \text{data}) \beta \, d\beta \tag{9.11}
$$

can be considered as the Bayesian estimate of the slope *β*, and the posterior standard deviation as its associated standard uncertainty,  $u(\beta)$ :

$$
u^{2}(\widehat{\beta}) = \int_{-\infty}^{\infty} p(\beta | \text{data})(\beta - \widehat{\beta})^{2} d\beta.
$$
 (9.12)

Finally, a 95 % credible interval [ $\beta_{\rm L}$ ,  $\beta_{\rm H}$ ] can be calculated from the posterior [\(9.10\)](#page-68-1), which satisfies

<span id="page-68-2"></span>
$$
\int_{\beta_{\rm L}}^{\beta_{\rm H}} p(\beta | \text{data}) \, \mathrm{d}\beta = 0.95. \tag{9.13}
$$

Since Eq. (9.13) does not uniquely determine a credible interval, further conditions need to be imposed, for example, that the credible interval is probabilistically symmetric, that is, symmetric around the Bayes estimate, or is the 95 % credible interval of shortest length. <span id="page-69-1"></span>**9.5.2.3 Informative prior** Assume that prior knowledge is available about the *σ* in a form of prior estimates  $\hat{\sigma}_i^2$  ( $i = 1, ..., p$ ). Assume also that the reliability of these prior estimates can be suppressed in terms of socialization continuition contribution of  $50\%$ estimates can be expressed in terms of coefficients of variation  $c_i$  , which are set to  $c_i =$  50 % for all values of *i*. This set of prior knowledge can be modelled using inverse Gamma distributions. Specifically, the parameters  $a_i$  and  $b_i$  of the inverse Gamma distributions are then given as follows:

$$
a_i = 2 + 1/c_i^2
$$
,  $b_i = (a_i - 1)\hat{\sigma}_i^2$ 

<span id="page-69-0"></span> $(9.14)$ 

i.e., the prior knowledge about each  $\sigma_i^2$  is modelled by an inverse Gamma distribution with mean  $\hat{\sigma}^2_i$  and variance  $c^2_i(\hat{\sigma}^2_i)$ .

**9.5.2.4** The resulting marginal posterior for the slope is then obtained as

$$
p(\beta|\text{data}) \propto \pi(\beta) \prod_{i=1}^{p} t_{n_i-1+2a_i} \left(\beta; \ y_i/x_i, \ \frac{(n_i-1)S_i^2 + 2b_i}{n_i x_i^2 (n_i-1+2a_i)}\right),\tag{9.15}
$$

where  $\pi(\beta)$  denotes the prior PDF for  $\beta$ , and  $t_{\nu}(x;m,s^2)$  stands for the PDF of a scaled and shifted *t*-distribution with *ν* degrees of freedom, that is,

$$
t_{\nu}(x; m, s^2) \propto \left(1 + \frac{1}{\nu} \frac{(x - m)^2}{s^2}\right)^{-(\nu + 1)/2}.
$$
 (9.16)

**9.5.2.5** The univariate PDF (9.15) can be evaluated and the summary statistics [\(9.11\)](#page-68-3) to [\(9.13\)](#page-68-2) obtained through standard procedures of numerical quadrature. For evaluating [\(9.15\)](#page-69-0), it is advantageous to calculate the logarithm of *p*(*β*|data) first, applying the exponential function afterwards.

**9.5.2.6** Markov Chain Monte Carlo method (MCMC) sampling can also be employed to draw samples from the posterior distribution arising from the above considerations. It is additionally assumed that prior knowledge is available about  $\beta$  in the form of a normal distribution with mean 1 and standard deviation 0.1. In Stan [104], the following measurement model can be written:

```
JCGM WG1 Committee draft
       data {
          int<lower=1> N
          vector[N] x;
          vector[N] y;
          vector[N] uy;
          vector[N] a_prior;
          vector[N] b_prior;
       }
       parameters {
          real beta;
          real<lower=0> sigma2[N];
       \sqrt{2}model {
          // priors for sigma^2
          sigma2 ~ inv_gamma(a_prior, b_prior);
          // prior for beta
```

```
beta \tilde{ } normal(1.0, 0.1);// likelihood
    for(i in 1:N) y[i] \tilde{ } normal(beta * x[i], sqrt(sigma[1]));
}
```
<span id="page-70-1"></span>**9.5.2.7 Weakly informative prior** The case of vague prior knowledge can be modelled by choosing a vague prior for the measurand *β*, such as a normal distribution with mean 1 and standard deviation 100. Likewise, values of  $c<sub>i</sub> = 1000\%$  can be assigned and Markov chain Monte Carlo sampling performed as before.

# **9.6 Reporting the result**

Table [9.3](#page-70-0) contains the estimate, its associated standard uncertainty, and the 95 % coverage interval obtained by application of the JCGM 100:2008 uncertainty framework (GUF) to OLS and WLS estimation, together with corresponding results for the Bayesian uncertainty analysis. The credible intervals determined by the Bayesian uncertainty analysis were taken as probabilistically symmetric intervals. Figure 9.1 shows the PDF for the measurand obtained by Bayesian uncertainty analyses in comparison with the results achieved by WLS estimation with uncertainty evaluated according to JCGM 100:2008 [6].

Table 9.3: Results obtained by ordinary least-squares (OLS) and weighted least-squares (WLS) with uncertainty evaluation according to JCGM 100:2008, as well as results from a Bayesian uncertainty evaluation

<span id="page-70-0"></span>

# **9.7 Discussion**

**9.7.1** The results obtained by application of the GUF to OLS and WLS estimation are different, which is due to the difference of the corresponding models [\(9.4\)](#page-66-2) used. Specifically, the WLS estimation with weights  $n_iW_i = 1/u^2(y_i)$  leads to a different estimate for the slope and a smaller uncertainty than OLS estimate. These weights are 'optimal' in the sense that they lead to a minimum uncertainty under all measurement models [\(9.4\)](#page-66-2).

be a "normal (1.0, 0.1);<br>
leta " normal (beta  $*$  x[1], nqrt (si.gma2(13));<br>  $f(x)$  for  $(x + \ln x)$  y[1] " normal (beta  $*$  x[1], nqrt (si.gma2(13));<br>  $\frac{1}{2}$  for  $(x + \ln x)$  y[1] " normal (beta;  $x + \ln x$ ) and the<br>[mi](#page-71-0)nimization **9.7.2** On the other hand, OLS does not apply 'optimal' weights and results in a larger uncertainty associated with its different estimate of the measurand. Since the corresponding measurement model is linear in the data, the squared standard uncertainty provides an unbiased estimate of the variance of the OLS estimator under hypothetical repeated sampling from the statistical model [\(9.5\)](#page-67-0). However, OLS estimate utilizes a measurement model that does not account for the fact that different observations have different variability. That is, measurements are assigned the same weight although their variability differs by almost an order of magnitude.



<span id="page-71-0"></span>Figure 9.1: Marginal posterior distribution *p*(*β*|data) for the regression slope from informative Bayes procedure and the estimate from classical OLS procedure

**9.7.3** The WLS estimate can be also be viewed as a solution to the statistical model [\(9.5\)](#page-67-0) if the variances  $\sigma_i^2$  were known. For unknown variances, however, the optimal weights  $n_i$ *W*<sub>*i*</sub> = 1/( $\sigma_i^2/n_i$ ) are also unknown. While the standard uncertainties *u*(*y*<sub>*i*</sub>) provide ap $n_i w_i = 1/(v_i/n_i)$  are also unknown. While the standard uncertainties  $u(y_i)$  provide approximate estimates of these optimal weights in the form of  $\sigma_i/\sqrt{n_i}$ , they can become unreliable when derived from a small number of repeated measurements.

**9.7.4** The Bayesian uncertainty analysis is based on the statistical model [\(9.5\)](#page-67-0) and does account for the different variability in the observations. At the same time, it does not use a single estimate of that variability to be used in a subsequent estimation of the measurand, but rather estimates the measurand and the variability in the observations simultaneously. Due to the straight-line model, all observations influence the estimation of all different variabilities in the observations, and observations with large variability will have less influence in the final result for the measurand. Furthermore, Bayesian inference allows prior knowledge about the measurand to be taken into account. For these reasons, the Bayesian uncertainty analysis is recommended for this example. It should be noted that methods from classical statistics can also be used to analyze the data on the basis of the statistical model [\(9.5\)](#page-67-0), which has not been considered in this example.

**9.7.5** Bayesian inference using the informative prior described in clause [9.5.2.3](#page-69-1) yields very similar results to those using the non-informative prior [\(9.5.2.7.](#page-70-1) The reason is that the data dominate the prior information taken for the measurand, and that the (hypothetical) prior knowledge about the variances has been taken only vaguely and in accordance with the observed variances. If either of these two latter conditions for the prior for the variances is removed, the results of an informative Bayesian inference might look appreciably different because each variance is modelled individually for each stimulus value and only a small number of repeated measurement results are available. In this case, the prior for the variance will be more informative. In other applications it can be reasonable to assume a common variance, which would reduce the sensitivity with respect to the prior for the variance significantly. Furthermore, the procedure provided for the Bayesian inference would then result in a single *t* distribution for the measurand in the non-informative
case, or the product of a single *t* distribution and an informative prior for the measurand otherwise.

case, or the product of a single 1 distribution and an informative prior for the measurement of the case and directly account for possible errors.<br>
9.7.6 The statistical model (9.5) used here does not directly account for **9.7.6** The statistical model [\(9.5\)](#page-67-0) used here does not directly account for possible errors in the measured values provided by the reference system. In fact, the example B2 of the guideline VDI/VDE 2600 part 2 [[112](#page-134-0)] reports non-vanishing uncertainties for them. The statistical model [\(9.5\)](#page-67-0) accounts for such an additional variability to a certain extent, as it includes unknown, individual variances for the dependent variable, that are simultaneously inferred together with the parameters of the straight line.

# **10 Conformity assessment of total suspended particulate matter in air**

## **10.1 Preamble**

The main goal of this example is to calculate, according to JCGM 106:2012 [8], risks of false decisions in conformity assessment when a normal distribution is not adequate to model prior information about the measurand. As a case study, mass concentration of total suspended particulate matter (TSPM) in ambient air is considered.

# **10.2 Objective and data**

**10.2.1** A total of 220 measurement results of the mass concentration of TSPM in ambient air, collected in 2009 in the proximity of a stone quarry located in Israel, was obtained according to Environmental Protection Agency (EPA) method IO-2.1 [41]. Such results were compared with the national (Israeli) regulation limit for air quality to study the occurrence of out-of-specification (OOS) results, as detailed in [69, 70].

10 Confor[mi](#page-131-1)try assessment of total suspended particulate-matrices<br>ter in air.<br>
101. Preamble<br>
12. The mathematic scanner when a normal distribution is both although<br>
the main goal of this example is to calculate, accordin **10.2.2** In this example, the focus is on the calculation of global and specific risks of false decision in the conformity assessment of the air quality, based on such kind of measurement results. The risk of underestimating the pollutant concentration is the consumer's (inhabitants') risk and that of overestimating it is the producer's (owner's) risk. Calculation of such risks is as important for the regulator protecting the inhabitants' quality of life in the area surrounding the quarry as for the manufacturers' association acting in the interests of the stone producers in the country.

## **10.3 Specification of the measurand**

For the characterization of TSPM, the EPA method IO-2.1 [[41](#page-129-0)] indicates the use of a highvolume sampler for collection of particles with aerodynamic diameters of  $100 \mu m$  or less. For each measurement, a large volume  $V$  of air between 1600  $\mathrm{m}^{3}$  and 2400  $\mathrm{m}^{3}$  was sampled at a constant flow rate and the mass *m* of the matter in the sampled air volume, collected on the sampler filter, was measured as the difference between the results of weighing the filter before and after sampling. The measurand is the average value of the mass concentration TSPM over the sampling period:  $\rho = m/V \text{ (mg m}^{-3)}$ .

## **10.4 Measured values and associated measurement uncertainty**

**10.4.1** The measured values of the TSPM mass concentration,  $\hat{\rho}_i$ , ( $i = 1, \ldots, 220$ ), are reported in magn-3 within the O14ste tut file (ovailable at [97]) and depicted in fig. reported in mg m<sup>-3</sup> within the Q1data.txt file (available at [[87](#page-132-0)]) and depicted in figure 10.1. (The running index *i* will be dropped hereafter for simplicity.) Note further that in JCGM 106:2012 a measured value would be denoted as  $\rho_{\rm m}$ , whereas in this document it is denoted by  $\hat{\rho}$ .

**10.4.2** A full uncertainty budget for the considered results is available in [[70](#page-131-1)], where it is shown that the major contribution to the uncertainty is that from the measurement of the sampled air volume. The relative standard uncertainty associated with any  $\hat{\rho}$  is 7.0 %.



<span id="page-74-0"></span>Figure 10.1: Histogram of the measured values of TSPM mass concentration and the corresponding lognormal probability density function (expression 10.1). The quantiles of this distribution are  $q_{0.025} = 0.042$  mg m<sup>-3</sup> and  $q_{0.975} = 0.229$  mg m<sup>-3</sup>

## **10.5 Tolerance limits**

The Israeli national regulations of ambient air quality prescribe an upper tolerance (regulation) limit  $T_U = 0.2$  mg m<sup>-3</sup> for TSPM mass concentration averaged over 24 h. This limit applies for any location, including those close to the quarry.

## **10.6 Decision rule and conformity assessment**

### **10.6.1 General**

**10.6.1.1** Regulations require direct comparison of measured values  $\hat{\rho}$  with  $T_U$ , as shown in figure [10.1.](#page-74-0) In the present example, acceptance limits  $A_U$  will be made varying in order to show their impact on the risk values of false decisions. When acceptance limits are taken to coincide with the tolerance limits (that is,  $A_U = T_U$ ), a 'shared risk' rule is considered as the decision rule for conformity assessment [[8,](#page-126-0) clause 8.2.1].

**10.6.1.2** In this instance, the inhabitants living in the area are considered the 'consumers', whereas the stone quarry owners are the 'producers'. So, the 'consumer's risk' relates to the inhabitants, whereas the 'producer's risk' relates to the quarry owners.

**EVALUATION**<br>
<sup>Trades</sup> <sup>2</sup> **Comparison**<br>
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<sup>2</sup> **Comparison**<br> **EVALUATION**<br> **EVALUATION**<br> **EVALUATION**<br> **EVALUATION**<br> **EVALUATION**<br> **EVALUATION**<br> **EVALUATION**<br> **EVALUATION**<br> **EVALUATION**<br> **10.6.1.3** The global and specific risks of false decisions in conformity assessment are defined in [8, clause 3.3] for both the consumer and the producer, and have different interpretations. While a specific risk is the risk of an incorrect decision based on a measurement result of a particular air sample, global risk refers to the probability of incorrect decisions based on future measurement results. Both kinds of risks rely on a Bayesian framework but require the calculation of different probability objects. Indeed, the posterior distribution (obtained through Bayes' theorem [[48](#page-129-1)]) is used for specific risks while the joint distribution is used for global risks. The posterior distribution is for the true value *ρ given* the measured value  $\hat{\rho}$ . The joint distribution is for the true value *and* the measured value. See clauses [10.6.3](#page-76-0) and [10.6.4.](#page-78-0)

#### **10.6.2 Bayesian framework**

**10.6.2.1** The evaluation of specific and global risks of false decisions on a characteristic of an item is described in JCGM 106:2012 [[8](#page-126-0)].

The underlying approach requires defining a prior probability density function (PDF)  $g_0(\rho)$ for TSPM mass concentration. Here, the prior PDF is based on experimental data (see JCGM 106:2012 [[8,](#page-126-0) clause 6.2.2]) and describes the distribution of the population of the TSPM mass concentration. An adequate fit for the experimental results, among possible candidate distributions, is the lognormal distribution:

<span id="page-75-0"></span>
$$
g_0(\rho) = \frac{1}{\rho \sigma \sqrt{2\pi}} \exp\left[-\frac{(\ln \rho - \mu)^2}{2\sigma^2}\right].
$$
\n(10.1)

The lognormal prior PDF [\(10.1\)](#page-75-0) is depicted as the curve approximating the histogram in figure [10.1.](#page-74-0) The location and scale parameter estimates of this distribution are  $\hat{\mu} = -2.325$ and  $\hat{\sigma}$  = 0.434, respectively. They are the sample mean and standard deviation, respectively, of the log-transformed data.

**10.6.2.2** The aforementioned estimates of  $\mu$  and  $\sigma$  ignore the fact that the TSPM data are observed with measurement uncertainty, which is now taken into account using Bayesian methods. A hierarchical Bayesian model can be specified as follows. The likelihood is given by

$$
\rho \sim \text{LN}(\mu, \sigma),
$$
  

$$
\hat{\rho}_i \sim \text{N}(\rho, 0.07 \hat{\rho}_i),
$$

and the priors are

$$
\mu \sim N(\mu_0, \sigma_0),
$$
  
\n
$$
\sigma \sim \text{Cauchy}(0, A),
$$
  
\n
$$
\rho_i \sim N(\hat{\rho}_i, 0.07\hat{\rho}_i).
$$

16.<br>
16.62 **Experiment Committee Committe** The prior distribution of  $\mu$  can be tuned, for example, such that its median corresponds to *T*<sub>U</sub>/2 ( $\mu$ <sub>0</sub> = −2.30) and its standard deviation is arbitrarily wide ( $\sigma$ <sub>0</sub> = 2.0). Similarly, in the absence of any strong prior information,  $A$  is set sufficiently large  $(A = 2.0)$  to yield a prior distribution for *ρ* whose 95 % quantiles span the observed data. The Cauchy distribution is chosen as it is a weakly informative prior and concentrates the density between 0 and *A* [[47](#page-129-2)]. The scale parameter *A* is set based on historical data.

In Stan, the posterior distributions of parameters  $\mu$  and  $\sigma$  can be obtained using the following model:

```
data {
    int<lower=1> N;
    vector<sup>[N]</sup> rho hat;
}
parameters {
    real mu;
    real<lower=0> sigma;
    vector[N] rho;
```

```
12.06<br>
\frac{1}{2} and 1<br>
\frac{1}{2}ergns priors<br>
\frac{1}{2} with \frac{1}{2} vergns priors<br>
\frac{1}{2} vergns priors<br>
\frac{1}{2} vergns \frac{1}{2} (2,2,2,0);<br>
\frac{1}{2} vergns \frac{1}{2} (2,2,2,2,0);<br>
\frac{1}{2} w
                  }
                  model {
                        // vague priors
                         mu ~ normal(-2.3, 2.0);
                         sigma \sim cauchy(0, 2.0);
                         rho \tilde{ } normal(rho_hat, 0.07 * rho_hat);
                         // likelihood
                         rho ~ lognormal(mu, sigma);
                         rho_hat \tilde{ } normal(rho, 0.07 * rho_hat);
                  }
```
**10.6.2.3** With Markov Chain Monte Carlo method (MCMC) sampling, facilitated by the R function stan() available from the package  $rstan$  [104], parameter estimates are obtained which, in this case, match closely with the simple method-of-moments estimates. Posterior estimates of  $\mu$  and  $\sigma$  are  $\hat{\mu} = -2.329$  and  $\hat{\sigma} = 0.434$ , which are nearly identical to the estimates provided by the method of moments.

**10.6.2.4** Each measurement result  $\hat{\rho}$  was modelled by a normal distribution with expectation equal to  $\hat{\rho}$  and standard deviation equal to the standard measurement uncertainty  $u = 0.07\hat{\rho}$  [[70](#page-131-1)]. The corresponding likelihood is hence

<span id="page-76-1"></span>
$$
h(\hat{\rho}|\rho) = \frac{1}{u\sqrt{2\pi}} \exp\left[-\frac{(\hat{\rho} - \rho)^2}{2u^2}\right].
$$
 (10.2)

**10.6.2.5** When both the prior PDF and the likelihood are normal distributions, the posterior PDF [[8,](#page-126-0) expression (1)] is normal [8, clause 7.2.1]. In this example, the prior PDF is log-normal. However, the method described in JCGM 106:2012 [8] to evaluate risks applies for any PDF.

## <span id="page-76-0"></span>**10.6.3 Global risks**

**10.6.3.1** For any considered upper acceptance limit  $A<sub>U</sub>$ , global risks for the consumer and the producer were calculated by combining the lognomal prior PDF [\(10.1\)](#page-75-0) and the normal likelihood (10.2), according to [[8,](#page-126-0) expressions (19) and (20)]. The joint PDF can be obtained in a numerically stable way by taking the logarithms of both the prior and the likelihood, summing the two transformed distributions and taking the exponential of the result, as follows:

```
P = function(xy) {
     p1 = dlnorm(xy[1], meanlog = -2.325, sdlog = 0.434, log = TRUE)
    p2 = \text{dnorm}(xy[2], \text{ mean } = xy[1], \text{ sd } = 0.07 * xy[1], \text{ log } = TRUE)exp(p1 + p2)}
```
**10.6.3.2** The integration of this joint PDF can be carried out using the R function cubintegrate defined in the package cubature [[81](#page-132-1)]. In the considered case, since all the involved PDFs were defined on the positive axis only, the lower integration limits were both taken as zero:

 $T_U = 0.2$  $A_U = 0.2$ require(cubature) Int = cubintegrate # rename the function cubintegrate  $C = 1/Int(P, lower = c(0,0), upper = c(Inf, Inf))$ \$int risk\_c =  $C * Int(P, lower = c(T_U, 0), upper = c(Inf, A_U))$ \$int risk\_p = C \* Int(P, lower = c(0, A\_U), upper = c(T\_U, Inf))\$int

**10.6.3.3** The obtained consumer's (red upward sloping line) and producer's (blue downward sloping line) global risks are displayed in figure 10.2, for  $A_U$  values varying in the interval  $[T_U - 0.05 \text{ mg m}^{-3}, T_U + 0.05 \text{ mg m}^{-3}]$ .



<span id="page-77-0"></span>Figure 10.2: Global consumer's risk (red upward sloping line) and global producer's risk (blue downward sloping line) as functions of the acceptance limit

Considering, for example, the special case in which  $A_{\text{U}} = T_{\text{U}}$ , consumer's and producer's global risks were 0.6 % and 0.7 %, respectively. These risk estimates can also be obtained without the need for numerical integration of the posterior PDF. Instead, a Monte Carlo method as shown in the R code below can be used:

```
N = 1e6y_true = nlnorm(N, meanlog = -2.325, sdlog = 0.434)y_meas = rnorm(N, mean = y_ttrue, sd = 0.07 * y_ttrue)
T_U = 0.20A_U = 0.20risk_c = sum(y_t = x) T_U & y_meas < A_U)/N
risk_p = sum(y_t true < T_U & y_mmeas > A_U)/N
```
**10.6.3.4** Consider determining the maximum acceptable value of  $A_U$  corresponding to a small consumer's risk, such as 0.01 %. The result is that such an acceptance limit should not exceed 0.17 mg m<sup>-3</sup>. However, in this instance, the global producer's risk would increase from 0.7% to 5%. The other way round,  $A_U$  should be at least equal to 0.23 mg m<sup>-3</sup> in

order to assure a producer's risk smaller than 0.01 %. In this case, the global consumer's risk would increase from 0.6 % to 2 %.

### <span id="page-78-0"></span>**10.6.4 Specific risks**

**10.6.4.1** Whereas global risks are concerned with future outcomes from this quarry, specific risks are concerned with a particular measured value at hand, as discussed in [8, clause 9.3.2].

**10.6.4.2** For a specific value  $\hat{\rho} < A_{\text{U}}$  (that is, the measured TSPM mass concentration is assessed as conforming to the regulatory limit), the consumer's specific risk is the integral of the posterior PDF  $g(\rho|\hat{\rho})$  over the region  $[T_U,\infty]$  — the interval of true values that would be deemed as nonconforming. For a specific value  $\hat{\rho} > A_{\text{U}}$  (that is, a value not conforming to the regulatory limit), the producer's specific risk is the integral of the posterior PDF over the interval  $[0, T_U]$  — the region of actually conforming true values.

**10.6.4.3** In both cases, the posterior PDF  $g(\rho|\hat{\rho})$  [8, expression (A.11)] is needed, which is not available in closed form. The integral of the posterior PDF is calculated by means of the R function integrate:

```
order to assure a producer's risk smaller than 0.01%. In this case, the global consigner's<br>
risk would coress from 0.6% to 2%.<br>
10.6.4.1 Whereas global risks are concerned with future outcomes from this quarry, specifies<br>
      P = function(rho, rho_hat)p1 = dlnorm(rho, meanlog = -2.325, sdlog = 0.434, log=TRUE)
        p2 = \text{dnorm}(\text{rho\_hat}, \text{mean} = \text{rho}, \text{sd} = 0.07 * \text{rho\_hat}, \text{log=TRUE})exp(p1 + p2)}
      T_{U} = 0.2# specific consumer's risk for rho_hat = 0.190
      rcl = integrate(P, T_U, Inf, rho_hat=0.190)$value
      rc2 = integrate(P, 0, Inf, rho_hat=0.190)$value
      risk_c = rc1/rc2# specific producer's risk for rho_hat = 0.225
      rpl = integrate(P, 0, T_U, rho_hat=0.225) $value
      rp2 = interrate(P, 0, Inf, rho_hat=0.225)$value
      risk_p = rp1/rp2
```
The obtained consumer's and producer's specific risks for various values of  $\hat{\rho}$  are displayed in figure [10.3](#page-79-0) (for  $A_{\text{U}} = T_{\text{U}}$ ).

Owing to the comparatively large width of the prior distribution, and thus to the meagre prior information, the posterior is essentially determined by the likelihood. The implication is that the specific risks as calculated in a purely frequentist context would be very close to those given here.

## **10.7 Interpretation of results**

**10.7.1** Studies on global risks, such as that conducted in clause [10.6.3,](#page-76-0) allow the involved parties (consumers and producers) to agree on an acceptance limit, thus balancing the safeguarding of the inhabitants' health and the economical interests of the quarry owners, in the considered example.



<span id="page-79-0"></span>Figure 10.3: Consumer's and producer's specific risks depending on the measurement results for the TSPM concentration in the quarry

**10.7.2** The approach in clause [10.6.4](#page-78-0) provides risks of false decision for a specific result and for a particular acceptance limit ( $A_U = T_U$ , in the considered case). From a practical point of view, no action will be undertaken when a measurement result falls below the acceptance limit, that is, when it is conforming with requirements. However, when a result exceeds the limit, it will be declared as non-conforming and some corrective action will be required. In this case, producers have at hand a tool for assessing the extent of their responsibility for such failure and possibly elaborate an appropriate reaction. As an example, for a non-conforming measured value  $\hat{\rho} = 0.225$  mg m<sup>-3</sup>, the specific producer's risk<br>is about 11.%, meaning that there is a non-negligible 11.% probability that such a value is about 11 %, meaning that there is a non-negligible 11 % probability that such a value corresponds to an actually conforming true value of the TSPM levels.

Consumer's specific risk<br>  $\frac{3}{5}$ <br>  $\frac{3}{5}$ 

trade-off is often sought.

# **puted quantity 11.1 Preamble**

**11 Effect of considering a 2D image as a set of pixels on a com-**

**11.1.1** This example is concerned with properties of a physical object extracted from a raster image. With a greater number of pixels to reproduce the image, a better description

result to other instrumental effects. In some cases, an increase in resolution does not lead to appreciable improvement of the global result because of other uncertainty sources. A

**11.1.2** In radiation therapy anatomical segmentation, image pixel size impacts on the area of a 2D section of a tumour (or an organ) given by the delineation of a 2D image. Such a delineation is made by an operator or an algorithm and constitutes one step in the estimation of the volume of a 3D region (the ultimate measurand) corresponding to a reconstructed 3D image of a region of interest (ROI) such as from a set of parallel planar sections.

# **11.2 Estimation of organ or tumour mass**

**11.2.1** The estimation of tumour mass necessitates some form of outlining of the tumour on an image. Under assumptions of homogeneity, tumour mass is directly proportional to its volume and so it can be obtained using suitable imaging. Various outlining or segmenting techniques are used for this purpose [51].

11 Effect of considering a 2D image as a set of pixels on a sómmetric<br>
11 Effect of considering a 2D image as a set of pixels on a sómmetric<br>
11.1 Protections, with a general with properties of a physical object extracted **11.2.2** The area of a tumour section is generally obtained from an ROI (region of interest) outlined on anatomical imaging data in the form of an array of pixels [[46](#page-129-3)]. Knowing that outline, the corresponding area of the ROI can be estimated. Approaches are considered for evaluating a measure of the outlining uncertainty, and the use of that measure to evaluate the uncertainty associated with estimated area. The method used will depend largely on the information and resources available at the time of outlining and the method employed by the operator or algorithm to define the ROI. When determining areas manually, an outline is typically drawn by an operator across all images that comprise the data set of concern.

**11.2.3** In practice, a set of planar sections of the ROI, usually at a constant spacing, is considered, from each of which the area is deduced and hence, by considering all such sections, the volume is estimated. Here, attention is paid to a single section as one step in the overall process.

**11.2.4** In general, segmentation methods can be divided into the following categories: manual methods, threshold-based methods, boundary-based methods and stochastic- and machine learning-based methods. More details of these categories are given in [[54](#page-130-1)] and the references therein. Here, just boundary-based methods are considered, such as gradientbased edge detection, which require an initial region of interest (ROI) to be defined inside which an operator or algorithm estimates the object boundary [[46](#page-129-3)].

## <span id="page-81-2"></span>**11.3 Measurement model**

**11.3.1** The measurand is the area of a plane section of the tumour.

**11.3.2** Figure [11.1](#page-81-0) shows an outline of a tumour obtained from a pixelated image, at two image resolutions, produced by scanning equipment, which will be used to provide numerical results and examine the effect of pixel size on the reconstructed shape. It is assumed that the outline of the tumour is contained within the darker-shaded pixels chosen by an operator or algorithmically. Each pixel is taken as of unit size (dimension  $1 \times 1$ ).



<span id="page-81-0"></span>Figure 11.1: Tumour profiles for fine and coarse pixel sizes: 85 and (right) 585 pixels defining the outline and interior

**11.3.3** The measurement model is specified by the algorithm used to compute the area of the tumour section. The area depends on knowledge of the boundary of the section, specifically a mathematical description of a smooth curve representing the boundary, for which a parametric representation is used:

<span id="page-81-1"></span>
$$
x = F(t), \qquad y = G(t), \tag{11.1}
$$

where the parameter *t* is (ideally) taken as arc length from some starting point.

**11.3.4** The tumour section area could also be computed, more approximately, as some average (such as the arithmetic mean) of two extreme areas. Referring to either panel of figure 11.1, one area would be that of the union of the grey pixels and the white pixels they surround. The other area would be that of the union of those white pixels. Such a method may not be effective for a small number of pixels. In clause [11.4,](#page-83-0) an argument is made in favour of representing the boundary by a smooth curve [\(11.1\)](#page-81-1). Clause [11.5.1](#page-85-0) makes a simple comparison of those approaches in terms of the outlines in figure [11.1.](#page-81-0)

**11.3.5** Various empirical functions such as periodic splines and parametric Fourier series can be used for *F* and *G* to form the profile. Parametric Fourier series are selected based on considerable experience [[26,](#page-128-0) [28,](#page-128-1) [35,](#page-128-2) [68](#page-131-2)] with the use of Fourier series for data analysis. For this purpose, some guidance in (53) concerned with the determination and/ise of the committee of the committee of the [dr](#page-132-2)aft of the state of the continents of the game of polynomial contents of the committee of the con For this purpose, some guidance in [[63](#page-130-2)] concerned with the determination and use of polynomial calibration functions is adapted. That standard gives advice on determining the degree of polynomial most appropriate for the data by first constructing polynomials of all degrees *n* from one up to some value  $n_{\text{max}}$  and then selecting a particular degree from that set. Here, parametric Fourier series involving an increasing number of harmonics are determined, choosing the number regarded as suitable according to criteria to be discussed. By analogy with polynomials, the number of harmonics is called *the degree* [85].

**11.3.6** A parametric Fourier series [\(11.1\)](#page-81-1) of degree *n* has  $2n+1$  parameters defining  $F(t)$ and an identical number defining *G*(*t*). The *x*- and *y*-components can be expressed as

<span id="page-82-1"></span><span id="page-82-0"></span>
$$
F_n(t) = b_1 + \sum_{r=1}^n [b_{2r} \cos(2\pi rt) + b_{2r+1} \sin(2\pi rt)],
$$
  
\n
$$
G_n(t) = c_1 + \sum_{r=1}^n [c_{2r} \cos(2\pi rt) + c_{2r+1} \sin(2\pi rt)].
$$
\n(11.3)

**11.3.7** Here, the outline is approximated by such a parametric Fourier series. The area *A* enclosed by the curve can be expressed in terms of the coefficients in the parametric Fourier series [\(11.2\)](#page-82-0) and [\(11.3\)](#page-82-1) by

<span id="page-82-2"></span>
$$
A = \pi \sum_{r=1}^{n} r(b_{2r+1}c_{2r} - b_{2r}c_{2r+1}),
$$
\n(11.4)

an expression derived in [[111,](#page-134-2) annex E5.1.A].

<span id="page-82-3"></span>**11.3.8** The data comprises a set of  $m = 38$  ordered values  $x_i$  and  $y_i$ ,  $i = 1, ..., m$ , taken as the midpoints of the grey pixels in figure 11.1 (right panel) through which the tumour boundary passes. Since the true values of of each coordinate pair lie somewhere in the containing pixel, it is assumed that these coordinates can be modelled by rectangular distributions R(−1*/*2, 1*/*2) relative to that midpoint. Accordingly, the standard uncertainties  $u(x_i)$  and  $u(y_i)$ , denoted by  $u_0$ , associated with the  $x_i$  and  $y_i$  are taken as the standard deviations of these distributions, namely,  $1/(2\sqrt{3}) \approx 0.29$ .

The points (*x<sup>i</sup>* , *y<sup>i</sup>* ) are regarded as values of the dependent variables *x* and *y* corresponding to an increasing set of values  $t_i$ ,  $i = 1, \ldots, m$  of the independent variable  $t$ . The values of  $t_i$  used are the cumulative chordal distance from  $(x_1, y_1)$  to  $(x_i, y_i)$ :

$$
t_1 = 0
$$
,  $t_i = t_{i-1} + [(x_i - x_{i-1})^2 + (y_i - y_{i-1})^2]^{1/2}$ ,  $i = 2,...,m$ .

These *t<sup>i</sup>* -values are normalized so that the total cumulative chord length becomes unity:

$$
t_i := \frac{t_i}{t_m}, \qquad i = 1, \ldots, m.
$$

Since, for sufficiently closely-spaced data, cumulative chord length approximates well arc length and is often used for that purpose, normalized cumulative chord length is employed as realizations of the parameter *t* [[45](#page-129-4)].

Parametric Fourier series of degrees  $n = 1, \ldots, n_{\text{max}}$ , are fitted to the data and a degree that represents a compromise between fidelity and smoothness is established. The *x*data and *y*-data can be treated separately since they are modelled by the functions [\(11.2\)](#page-82-0) and  $(11.3)$ , which have no parameter in common.

A goodness of fit mestate is formed for each degree. A common mestar, used here, is equiped<br>packing the fitted same of same of squares of the deviations of the first signere of<br>the fitted same of same of squares of the de A goodness-of-fit measure is formed for each degree. A common measure, used here, is the chi-squared statistic  $\chi^2_{\rm obs}$ , the sum of squares of the deviations of the fitted series of degree *n* from the data ( $x_i$ ,  $y_i$ ), weighted inversely by the squared standard uncertainties associated with those values ([[62](#page-130-3)] provides an analogous description for polynomial regression). When covariances associated with these data are present, a modified measure is used, but that is not the case here.

In terms of the computed Fourier series for degree *n*, the observed chi-squared value is

$$
\chi_{\text{obs}}^2(n) = \sum_{i=1}^m \frac{[x_i - F_n(t_i)]^2 + [y_i - G_n(t_i)]^2}{u_0^2} = \frac{1}{12} \sum_{i=1}^m \{[x_i - F_n(t_i)]^2 + [y_i - G_n(t_i)]^2\}
$$

with degrees of freedom  $\nu = 2m-4n-2$  (the number of data less the number of adjustable parameters), from which is obtained the root-mean-square residual RMSR(*n*) for each *n*:

<span id="page-83-1"></span>RMSR(n) = 
$$
\left[\frac{\chi_{\text{obs}}^2(n)}{2m - 4n - 2}\right]^{1/2}.
$$
 (11.5)

**11.3.9** In addition to using these RMSR-values to choose a value of *n*, various modelselection criteria such as Akaike's Information Criterion (AIC), corrected Akaike's Information Criterion (AICc) and the Bayesian Information Criterion (BIC) [16] are used. Generically, for *m* data points and a model with  $n + 1$  parameters, these criteria are

$$
AIC(n) = \chi_{obs}^{2}(n) + 2(n + 1),
$$
  
 
$$
AICc(n) = AIC(n) + 2(n + 1)(n + 2)/(m - n - 2),
$$
  
 
$$
BIC(n) = \chi_{obs}^{2}(n) + (n + 1)ln(m).
$$

In this application, there are 2*m* data, the *x*- and the *y*-values, and  $4n + 2$  parameters in the functions [\(11.2\)](#page-82-0) and [\(11.3\)](#page-82-1). So, the generic *m* and *n* in the above information criteria and expression [\(11.5\)](#page-83-1) for RMSR and *S* are replaced by 2*m* and 4*n* + 2, respectively.

**11.3.10** All three information criteria are designed to balance reproducing the data and parsimony of model [80]. Given a number of candidate models, here parametric Fourier series of degrees  $n = 1, ..., n_{max}$ , the model having the smallest value of AIC (or AICc or BIC) would usually be selected. According to [[63](#page-130-2)], experience with polynomial calibration problems indicates that the same degree of polynomial is often selected by all three criteria, although there may be some exceptions such as when the data set is small.

## <span id="page-83-0"></span>**11.4 Model selection**

**11.4.1** Consider the profile of the tumour section shown in the left panel (85 pixels) of figure 11.1. The midpoints of the boundary pixels, each of unit size, are given by the  $m = 38$ coordinate pairs with the bottom-left corner of the grid corresponding to a location  $(1, 1)$ and, consequently, the top-left corner pixel of the tumour outline corresponding to location (3, 12).

**11.4.2** Consider approximations to this set of discrete data points by continuous profiles in the form of curves represented by the parametric Fourier functions [\(11.2\)](#page-82-0) and [\(11.3\)](#page-82-1) with the values of *t* taken as cumulative chord length as in clause [11.3.](#page-81-2)

**11.4.3** Parametric Fourier series [\(11.4\)](#page-82-2) with  $n = 1, 2, \ldots, 15$  harmonics were fitted to the data by least squares Ordinary least squares was used because of the common uncertainty. The corresponding values of RMSR, AIC, AICc and BIC, and the value  $\widehat{A}$  of area  $A$ , were calculated and given in table [11.1.](#page-84-0) Figure [11.2](#page-85-1) displays the parametric Fourier series for these degrees.

<span id="page-84-0"></span>Table 11.1: Root-mean-square residuals and information criteria for various numbers, *n*, of harmonics and tumour section area estimates  $\widehat{A}$  in the example of the coarse-pixelated tumour image in figure [11.1](#page-81-0) (left)

n	RMSR	AIC	AICc	ВIС	$\widehat{A}$
1	0.96	38	39	46	65.1
2	0.63	23	25	34	63.9
3	0.60	25	30	40	64.2
4	0.43	23	31	41	65.7
5	0.43	27	39	48	65.7
6	0.39	30	48	55	65.2
7	0.40	33	59	62	65.2
8	0.39	37	73	68	65.4
9	0.40	41	90	76	65.4
10	0.41	45	112	83	65.4
11	0.43	49	141	90	65.3
12	0.44	52	180	97	65.3
13	0.44	56	236	104	65.3
14	0.40	59	325	110	65.2
15	0.37	63	485	117	65.4

11.43 Numeric Couries exists (11.4) with  $n = 1, 2, \ldots, 15$  have notices we used because of the common ungertainty<br>from by less equares Ordinary best squares was used because of the common ungertainty<br>calculated and given **11.4.4** The intention is to select an appropriate number of harmonics that offers a balance between accounting for the data uncertainties and yielding an acceptably smooth representation of the profile. For the coarsely pixelated image, the sequence of RMSR values (column 2 of table 11.1) initially shows a tendency to decrease to a minimum value of approximately 0.4 when  $n = 6$  (six harmonics). Then, the sequence saturates at about that level for a few degrees, subsequently tending to behave erratically for larger *n* due in part to the errors in the data being followed too closely. The Fourier series of degree 6 was selected, for which the RMSR value is 0.39. This value is framed in figure [11.2](#page-85-1) as is the corresponding area estimate of 65.2. Curiously, the three information criteria (AIC, AICc and BIC) all exhibit minimum values for  $n = 2$ , shown underlined in table [11.1,](#page-84-0) for which the RMSR value of 0.63 is some 50% higher than that, 0.39, for  $n = 6$ .

<span id="page-84-1"></span>**11.4.5** This observation about the information criteria is not consistent with general experience of selecting an appropriate degree in polynomial calibration studies. Further, this choice of degree  $(n = 2)$  does not seem particularly consistent with the data on which it was based as seen from the poor fit when this number of harmonics are used to fit the tumour outline (figure [11.2\)](#page-85-1). For instance, the re-entrant feature in the 'north-east' region of the outline is not reproduced. This and other features, however, are better reproduced by Fourier series of degree  $n \geq 5$ . By the time  $n = 15$  is reached, spurious effects (not shown), due to following the 'noise' in the data, appear in the curves that are not apparent in the perceived outline of the tumour.



<span id="page-85-1"></span>Figure 11.2: Fourier representations of a section of a tumour using 1 to 9 Fourier harmonics

**11.4.6** It would be unwise to select a low degree such as  $n \leq 3$  since the corresponding Fourier series over-smooth the tumour outline. At the other extreme, the choice of a large number of harmonics such as 13 or greater means that the Fourier representations follow too closely the sequence of pixel midpoints specified by the data. An argument could be made purely on visual grounds that any number of harmonics between 5 and 9 might be suitable. Degree six was selected for the reasons above.

**11.4.7** Recently [[80](#page-132-3)], modified AIC and BIC criteria for model selection have come available based on incorporating the Shannon entropy in the criteria. Although their use for synthetic data sets containing 50 to 550 points has been demonstrated to provide better discrimination in the choice of model, to the knowledge of the JCGM they have not yet been applied in metrology.

## **11.5 Uncertainty evaluation**

<span id="page-85-0"></span>**11.5.1** The corresponding area contained within the curve modelled by a parametric Fourier series is 65.2 (clause [11.4.4\)](#page-84-1). This value should be compared with the minimum and maximum values given by the inner and outer profile of the boundary pixels, namely, 47 and 85, whose mean is 66 — close in terms of this simplistic approach to the value from the Fourier representation.

**11.5.2** The RMSR value of 0.39 (clause [11.4.4\)](#page-84-1) does not relate to the uncertainty associated with the estimate of the area of the section of the tumour. Because of the complexity of propagating uncertainties associated with pixel midpoint locations through the measurement model [\(11.4\)](#page-82-2), a Monte Carlo method (MCM) was applied. The *t*-values have uncertainties associated with the uncertain *x*- and *y*-values, on which they based. Their effect is taken into account by applying MCM, which is relatively straightforward.

**11.5.3** As in clause [11.3.8,](#page-82-3) the *x*- and *y*-coordinates of the midpoints of each boundary pixel (square and of unit area) are modelled by rectangular distributions R(−1*/*2, 1*/*2) relative to that midpoint . By sampling from these distributions, a new set of equally plausible *x*, *y* coordinates is created, a Fourier series with  $n = 6$  harmonics computed, and the corresponding area within the outline calculated. Repetition a large number of times  $(10<sup>5</sup>)$  yields a distribution of values for the area whose average is the Monte Carlo estimate for the area and whose standard deviation is the associated standard uncertainty.

**11.5.4** This calculation gave the estimated PDF for tumour section area shown in figure [11.3.](#page-86-0) From the Monte Carlo results, the area estimate was 65.0 and the associated standard uncertainty was 1.6. This estimate is practically indistinguishable from the aforementioned point-estimate of 65.2 square pixels (table 11.1), indicating little bias in the process.



<span id="page-86-0"></span>Figure 11.3: Scaled histogram approximation to the PDF for the cross-sectional area of the tumour given by Monte Carlo sampling using parametric Fourier series with six harmonics

## **11.6 Reporting the result**

**11.6.1** For the example presented in clause [11.4,](#page-83-0) two estimates of the area *a* of the considered section of the tumour and the associated standard uncertainties were provided. One result (clause [11.5.1\)](#page-85-0) was

 $\hat{a} = 67, \quad u(\hat{a}) = 11$ 

using the simple approach of taking the mean of the areas defined by the inner and outer perimeters of the shaded pixels in figure [11.1](#page-81-0) (centre) and taking these two extremes as endpoints of a rectangular distribution of possible values. The second result was

$$
\hat{a} = 65.0,
$$
  $u(\hat{a}) = 1.6,$ 

which was obtained by modelling the midpoints of the boundary pixels by a parametric Fourier series of degree six.

**11.6.2** To demonstrate stability of the results with respect to the model used, MCM results were also obtained for neighbouring degrees. For  $n = 5$ , the estimate 65.3 and standard uncertainty 1.5 were obtained. For  $n = 7$ , the corresponding values were 65.0 and 1.5. These estimates are consistent.

## **11.7 Interpretation of results**

which was channed by modelling the midpoints of the boundary pixels by a parametric<br>courier series of degrees aix.<br>  $1.62$  To demonstrates and the small in the results with respect to the model used, of  $\mathcal{M}$  legalits<br> **11.7.1** The volume or mass of an organ or tumour is generally obtained from a region of interest (ROI) outlined on anatomical imaging data [46]. It is therefore possible to estimate the outlining accuracy from consideration of factors that affect delineation. The method used will depend largely on the information and resources available at the time of outlining and the method employed by the operator or algorithm to define the ROI. The process adopted here uses regression techniques to fit parametric Fourier functions to the midpoints of the pixels approximating the outline while taking into account that the midpoints are a crude guide to points on the boundary. A Fourier series of degree six appears to capture the given outline quite satisfactorily bearing in mind that only data points representing the midpoint of the boundary pixels are used.

**11.7.2** The importance of the principles used to select a measurement model of an appropriate complexity cannot be overemphasized. Identification of the point where the RMSR values start to saturate with respect to model complexity is generally reliable [[24,](#page-128-3) [26,](#page-128-0) [27,](#page-128-4) [31,](#page-128-5) [32,](#page-128-6) [34,](#page-128-7) [76](#page-131-3)]. That approach is better than the use of traditional information criteria such as AIC for the example considered here. The reason is believed to be the relative paucity of profile boundary data, whereas the information criteria generally function better for large quantities of data. However, it is accepted that identification of where the RMSR values start to saturate is difficult to automate and there is much value in having the analysts involved in such decision making.

**11.7.3** To illustrate the numerical accuracy of the parametric Fourier representation, the cross-sectional area of the fine pixelated image shown in figure [11.1](#page-81-0) (right), taking each pixel as unit size, is also analyzed. This contour consists of *m* = 113 darkly shaded pixels, each of dimension  $0.35 \times 0.35$ . Table [11.2](#page-88-0) shows the results of fitting parametric Fourier series to these data along with RMSR and the other model selection criteria.

The RMSR residuals again saturated essentially at degree 6. The information criteria were consistent in selecting degree 4, again different from that according to RMSR saturation but better than for the coarser-pixel data, presumably because of the larger quantity of data involved. The corresponding area estimates shown in table [11.2](#page-88-0) are corrected by the factor  $(0.35)^2$  to take account of the relative pixel size between the two images.

**11.7.4** The sequence of boundary coordinates might have associated correlation because there is pattern in the arrangement of the corresponding pixels as a consequence of underlying smoothness in the tumour boundary. This correlation has no tangible effect on the results presented here and is discussed in greater detail in [[111,](#page-134-2) example E5.1].

65.6 1.28 91 92 101 $\mathbf{1}$ 64.8 $\overline{2}$ 0.76 39 40 53 3 0.67 35 39 65.2 54 66.1 4 0.42 31 26 49 56 66.1 5 0.43 30 37 66.1 0.40 33 63 6 43 0.39 66.1 36 50 71 7 66.0 0.39 8 40 58 78 66.0 9 0.40 44 68 86 66.0 10 0.41 78 48 94 66.1 11 0.41 102 51 90 66.0 12 0.41 $1\overline{1}0$ 55 103 66.0 13 0.42 59 119 118 65.9 14 0.40 62 137 125 15 0.40 66.0 66 158 133 The Fourier-series approach used is expected to be suitable for representing smootl tumour outlines. For cases in which the outline has 'sharp corners', very fine pixelation would be required to support such features. Moreover, Fourier series would not be able to represent such features efficiently. Piecewise representations such as periodic splines [37] would be beneficial in such circumstances although their use would not be so straightfor-	RMSR	AIC	AICc	<b>BIC</b>	$\widehat{A}$

<span id="page-88-0"></span>Table 11.2: Root-mean-square residuals, information criteria for various numbers, *n*, of harmonics and tumour section area estimates  $\widehat{A}$ , in the case of a finer-pixelated image

# **12 Between-bottle homogeneity of reference materials**

## **12.1 Preamble**

## **12.2 Specification of the measurand**

12		Between-bottle homogeneity of reference materials					
	12.1 Preamble						
		12.1.1 An essential element in the production of certified reference materials (CRMs) and proficiency test materials in batch form is the evaluation of the between-bottle ho- nogeneity [107]. This form of (in)homogeneity accounts for the (small) differences in he property of interest between the bottles (or more generally, items [66]) and including t in the uncertainty budget of the property value ensures that the value and associated incertainty are valid for each bottle in the batch, rather than for the batch as a whole 108, 107]. This example shows how a Bayesian hierarchical model can be used to de- ermine the between-bottle standard deviation of the amount fraction nitrogen in a set of ynthetic natural gas mixtures. The model is demonstrated for a data set involving nitrogen where the classical analysis does not provide a solution.					
		12.2 Specification of the measurand					
		12.2.1 The measurand is the amount fraction nitrogen in a set of 10 synthetic natural as mixtures. This example focuses on determining an uncertainty component that char- cterizes the between-bottle homogeneity in a batch of a reference material [66]. The data rom the homogeneity study are shown in table $12.1$ [105]. For each of the 10 mixtures, i values for the amount fraction nitrogen were obtained (called "replicates"). Table 12.1: Amount fraction of nitrogen in 10 gas mixtures expressed in cmol mol $^{-1}$					
	Mixture				Replicate 1 Replicate 2 Replicate 3 Replicate 4 Replicate 5		
	D520472 H95396 VSL190663 D520467 D520834 D520361 D520270 D520446 VSL190485 VSL190977	0.424 577 0.425 572 0.424 152 0.426320 0.424855 0.425 104 0.425 750 0.425 547 0.426326 0.425 968	0.425 167 0.425411 0.425 517 0.424672 0.425079 0.424773 0.424917 0.426483 0.424646 0.424069	0.425 379 0.423 638 0.425 638 0.425 211 0.425 413 0.426 424 0.424779 0.424631 0.425 205 0.425 988	0.424522 0.425 301 0.424 207 0.425 533 0.424729 0.424266 0.425 086 0.425 968 0.426302 0.425 489	0.424805 0.424 527 0.425 135 0.425 864 0.424725 0.424632 0.425 318 0.424620 0.425 020 0.423936	
	lways possible [105].	12.2.2 Traditionally, classical analysis of variance (ANOVA) is used [66, 59] for this pur- bose, which is more fully described in [107]. The parameter of interest is the between- roup standard deviation, which in this specific case is called the between-bottle standard leviation [107, 105]. Whereas classical ANOVA works well if the between-bottle homo- eneity effect is of similar magnitude as the measurement repeatability or greater, diffi- ulties arise when the between-bottle homogeneity effect is (substantially) smaller than he repeatability effect [105]. Such situations should be avoided [66, 107] but that is not 12.2.3 Using the traditional ANOVA, the following values of sum-of-squares (SS), degrees					
		of freedom $(v)$ , and mean squares (MS) are obtained (see table 12.2).					

<span id="page-89-0"></span>Table 12.1: Amount fraction of nitrogen in 10 gas mixtures expressed in cmol mol<sup>-1</sup>



<span id="page-90-0"></span>

**12.2.4** The standard deviation corresponding to the between group effect is, in this case,

$$
s_{\text{between}} = \sqrt{\frac{\text{MS}_{\text{between}} - \text{MS}_{\text{within}}}{5}},
$$

which cannot be computed since  $MS<sub>between</sub>$  is smaller than  $MS<sub>within</sub>$ .

Table 12.2: Analysis of variance resul[t](#page-131-4)s for the homogeneity data for the amount ranging<br>nitrogen<br>
There is a second to the complete draft of the complete draft of the same is a set of the commutation of the same<br>
Theorem **12.2.5** The data set suffers from poor repeatability, but that does not rule out that the between-bottle standard deviation is not negligible when compared to the uncertainty associated with the batch average for the amount fraction nitrogen [66]. The use of a Bayesian hierarchical model which, apart from the use of prior probability density distributions for the parameters, is similar to the traditional one-way ANOVA model widely used in the evaluation of homogeneity studies [[107](#page-134-3)], but does not suffer from the shortcomings of ANOVA [[48,](#page-129-1) [105](#page-134-5)].

#### **12.3 Measurement model**

**12.3.1** The statistical model relating the observed amount fractions  $y_{ij}$  for mixture *i*  $(i = 1...10)$  and replicate *j*  $(j = 1...5)$  to the mean amount fraction  $\mu$ , the error in the amount fraction in mixture *i*,  $B_i$  and the random measurement error  $\varepsilon_{ij}$  takes the form [[107](#page-134-3)]

$$
y_{ij} = \mu + B_i + \varepsilon_{ij}.\tag{12.1}
$$

The objective of the evaluation is to determine  $\tau^2 = \text{Var}(B_i)$  and  $\sigma^2 = \text{Var}(\varepsilon_{ij})$ . If no pooling is used, then  $\sigma_i^2 = \text{Var}(\varepsilon_{ij})$ , that is, a variance is computed for each mixture. In a between-bottle homogeneity study, the assumption is typically made that the withingroup standard deviations are independent and identically distributed. In that case, *σ* is the pooled within-group standard deviation. The Bayesian approach presented here also makes such an assumption[105].

**12.3.2** In this example, a Bayesian model is used, which implies that a joint prior probability density function (PDF) should be chosen for the model parameters. In the case that the parameters are assumed to be mutually independent a priori, then this joint prior PDF can be factored as three PDFs, one for each of the parameters. These probability density functions are specified as follows:

<span id="page-90-1"></span>
$$
\mu \sim N(\mu_0, \sigma_{\text{target}}^2),\tag{12.2}
$$

$$
\tau \sim \text{Cauchy}(0, \tau_0),\tag{12.3}
$$

<span id="page-90-2"></span>
$$
\sigma \sim \text{Cauchy}(0, \sigma_0). \tag{12.4}
$$

The prior probability density function (hereafter *prior*) for  $\mu$  is a normal distribution with mean  $\mu_0$  (elicited from the specification of the composition of the gas mixtures) and a

standard deviation that reflects how close the amount fraction for the component of interest of the batch is expected to be to the specified value ( $\mu_0 = 0.4250 \text{ cmol mol}^{-1}$ ). The manufacturer specifies that the actual amount fraction will not differ more than 5 % from the specified amount fraction. This specification is interpreted as a 95 % coverage interval, and hence a relative standard deviation of 2.5 % is used. This standard deviation is sufficiently large to ensure that the posterior probability density function (hereafter *posterior*) will be dominated by the data [[105,](#page-134-5) [106](#page-134-6)].

**12.3.3** The prior for the between-bottle standard deviation  $\tau$  is chosen to be the half-Cauchy distribution with location parameter 0 and scale parameter  $\tau_0$ . The latter is obtained from the specification for the production of the batch gas mixtures, which is usually larger than the value expected for *τ*. The half-Cauchy distribution concentrates half of the density between 0 and the scale parameter, taking only non-negative values [[47](#page-129-2)].

**12.3.4** A similar approach is used for the prior of  $\sigma$ . The scale parameter  $\sigma_0$  is set to be equal to the repeatability standard deviations of the amount fractions in this type of mixtures, as observed in previous measurements.

**12.3.5** The likelihood is, conditionally on the parameters, a normal distribution [[106](#page-134-6)]:

$$
y_{ij}|\mu,\tau,\sigma \sim N(\mu,\tau^2+\sigma^2). \tag{12.5}
$$

## **12.4 Implementation**

**12.4.1** The model as described in the previous sections is used with Bayes' rule. From the weakly informative priors for  $\mu$ ,  $\tau$ , and  $\sigma$  (expressions (12.2)-(12.4)), using the data and the likelihood, a joint posterior for the model parameters is obtained. From this posterior, the value for  $\tau$ , the between-bottle standard deviation, is calculated.

**12.4.2** In Stan [[104](#page-134-1)] code, the statistical model of the between-bottle homogeneity study with pooling of the within-group standard deviations reads as

```
64<br>
1. CCM GUN 5:2026-12.06<br>
mandand devention that reflexes how close the amount fraction for the component of the<br>
memorializary excites that the neutral memorialization for the component of the<br>
memorializary specifie
            data {
                    int<lower=1> N;
                    int<lower=1> K;
                    matrix[N,K] y;
            }
            parameters {
                    real<lower=0> mu;
                    real<lower=0> tau;
                    real<lower=0> sigma;
                    vector[N] eta;
            }
            transformed parameters {
                   vector[N] theta = mu + tau * eta;
             }
            model {
                  // priors
                  mu \degree normal(0.4250, 0.0025 * 0.4250);
```

```
tau ~ cauchy(0, 0.0015 * 0.4250);
    sigma \sim cauchy(0, 0.0020 * 0.4250);
    eta \tilde{ } normal(0, 1);
    // likelihood
    for (i in 1:N) \{ y[i, ] \in \text{normal}(theta[i], signa) \}}
```
**12.4.3** The statistical model consists of two blocks, which are briefly described below:

- The data block declares all data along with the explicit dimensions and variable types. Here, the number of gas mixtures ("bottles") N ( $N = 10$ ) and the number of replicates K ( $K = 5$ ) are declared, following by the  $N \times K$  matrix of observed amount fractions y.
- In the parameters block, the model parameters mu, tau and sigma are declared, as well as an additional variable eta, which is used for an efficient implementation of the hierarchical model [48]. Indeed, sampling eta is more efficient than directly trying to obtain the group means. Instead, group means are calculated in the transformed parameters block. Lastly, the model block specifies the statistical model which includes the specifications of prior probability distributions for all parameters as well as the likelihood.

**12.4.4** Experience from previous between-bottle homogeneity studies for nitrogen in natural gas has indicated that the relative repeatability standard deviation for the nitrogen amount fraction is 0.20 % and the specification for the between-bottle homogeneity is 0.30%. The latter is interpreted as an expanded uncertainty with coverage factor  $k = 2$ , thus  $\tau/\mu$  = 0.15% and  $\sigma/\mu$  = 0.20%.

<span id="page-92-0"></span>**12.4.5** Running the Markov Chain Monte Carlo method (MCMC) sampling with 50 000 iterations and a warm-up of 25 000 iterations, using 4 chains [[106,](#page-134-6) [105](#page-134-5)] yields the posterior PDF for  $\mu$ ,  $\tau$ , and  $\sigma$  shown in figure [12.1](#page-93-0) and the following output:

```
tau "conchy(0, 0.0016 + 0.4256);<br>
etax "conchy(0, 0.0004 + 0.4256);<br>
etax "socialvy(0, 0.0004 + 0.4256);<br>
etax "socialvy(0, 0.0004 + 0.4256);<br>
// 1xkelihood .);<br>
/or (4 in 1:0) f y[i, 1] "sormal (cheta[il], signal);<br>
\sqrt{# Inference for Stan model: anon_model.
 # 4 chains, each with iter=5000; warmup=2500; thin=1;
 # post-warmup draws per chain=2500, total post-warmup draws=10000.
 #
 # mean se_mean sd 2.5% 97.5% n_eff Rhat
 # mu 0.42514 0 0.00011 0.42492 0.42536 9269 1.00007
 # tau 0.00014 0 0.00011 0.00000 0.00039 6334 0.99994
 # sig 0.00069 0 0.00007 0.00057 0.00085 12789 0.99985
  #
 # Samples were drawn using NUTS(diag_e) at Sun Sep 18 17:19:33 2022.
 # For each parameter, n_eff is a crude measure of effective sample
  # size, and Rhat is the potential scale reduction factor on split
  # chains (at convergence, Rhat=1).
```
The uncertainty due to the MCMC is given in the column se\_mean and is smaller than the last digit stated in the columns mean and sd. Repeated runs of the MCMC are expected reproduce these values with the number of decimals displayed.



<span id="page-93-0"></span>

**12.4.6** A sensitivity analysis, aimed at assessing how influential the choices of the priors are, revealed that the posterior of *τ* is quite sensitive with respect to the choice of the scale parameter of the Cauchy distribution [105]. Yet, in this example, the scale parameter was chosen based on the specification of the performance of the preparation process. As this process had been used several times before, these specifications were deemed to be adequate for eliciting this prior.

**12.4.7** Where traditional ANOVA fails at quantifying the between-bottle homogeneity effect, the Bayesian approach provides a probability density function for  $\tau$ , from which the between-bottle standard deviation can be derived. As the uncertainty evaluation in accordance with the GUM [6] and ISO Guide 35 [[66](#page-131-4)] requires only a point estimate for the uncertainty component  $\tau$ , there are several ways to obtain such an estimate from the posterior probability density function – the mean, the mode, or the median. Given the skewness of the posterior of  $\tau/\mu$  (figure [12.1\)](#page-93-0), these three estimates are not equivalent: the mode of the posterior of  $\tau$  is 0.3 µmol mol<sup>-1</sup>, the median is 1.1 µmol mol<sup>-1</sup>, and the mean is 1.4 µmol mol<sup>-1</sup>.

**12.4.8** From the MCMC draws, the shortest coverage intervals can also be obtained. The following code, requiring the R packages HDInterval and rstan performs the calculation:

```
require(HDInterval)
require(rstan)
fit. mcmc = extract(fit)tau.hdi = hdi(fit.mcmc$tau/fit.mcmc$mu, credMass = 0.95)
sigma.hdi = hdi(fit.mcmc$sigma/fit.mcmc$mu, credMass = 0.95)
```
**12.4.9** The third line converts the output from Stan into the appropriate form and then the function hdi is used to compute the highest posterior density intervals [[75](#page-131-5)]. The lower (L) and upper (H) limits of the 95 % highest posterior density intervals are shown in table [12.3.](#page-94-0)

<span id="page-94-0"></span>Table 12.3: Highest posterior density intervals for the uncertainty components *τ/µ* and *σ/µ*.

Parameter	$L_{95\%}$	$H_{95\%}$	Unit
$\tau/\mu$		$0.0000 - 0.0008$	
$\sigma/\mu$		0.0013 0.0020	

#### **12.5 Reporting the result**

**12.5.1** The primary result is the value for the between-bottle standard deviation *τ*. In a previous paper [[105](#page-134-5)], the mean was chosen as estimate for *τ*, which is the most cautious option (it leads to the largest value for this uncertainty contribution). For data sets where the between-bottle variability is larger, the differences between the three options become smaller.

**12.5.2** Alternatives to using the mean include the use of the median or mode of the posterior probability density function for *τ*. Even  $τ = 0$  could be justified, as the lower end of the 95 % highest posterior density interval is practically zero.

**12.5.3** There are different choices for the credible interval as well. The example provides two intervals, the shortest in table 12.3 and the probabilistically symmetric in [12.4.5.](#page-92-0) Shortest intervals are generally more appropriate for skewed posteriors, as in the case of the posterior for *τ*.

lower (1) and upper (1) limits of the 95% highest posterior density intervals are shown<br>in the 12.2,<br>this has let 12.2,<br>The primary established by the summer  $\frac{1}{\sqrt{\mu}} \frac{W_{\mu\nu} - W_{\nu}}{0.0013 - 0.0029 - 1}$ <br>The primary res

# **13 Measurement of Celsius temperature using a resistance thermometer**

## **13.1 Preamble**

13 Measurement of Celsius temperature using a resistance-thermometer<br>
2D. The example is concerned with the measurement of Gelsius temperature by com-<br>
10.11 President and the including latituding the residence of the mea **13.1.1** This example is concerned with the measurement of Celsius temperature by comparing the resistance of an industrial platinum resistance thermometer with that of a standard resistor using a resistance bridge. The measurement of a single Celsius temperature (see [13.2\)](#page-95-0) is described by an implicit univariate measurement model  $[12,$  $[12,$  CLause 4.1, note 2] and that of several temperatures (see [13.3\)](#page-97-0) by an implicit multivariate measurement model [[7,](#page-126-3) clause 6.3.1]. The example illustrates the treatment of these (univariate and multivariate) measurement models using the law of propagation of uncertainty (LPU) generalized to implicit measurement models and models with more than one measurand [[7,](#page-126-3) [6](#page-126-1)].

**13.1.2** The application of LPU in this example depends on the linearization of a nonlinear (univariate or multivariate) measurement model. Ideally, either a statement should be made that any replacement of a non-linear measurement model by a linear model has been checked to be acceptable, or that such a check has not been made.

## <span id="page-95-0"></span>**13.2 Measurement of a single Celsius temperature**

<span id="page-95-2"></span>**13.2.1** The Celsius temperature  $\theta$  is measured by comparing the resistance  $R(\theta)$  of a resistance thermometer with the resistance  $R<sub>S</sub>$  of a standard resistor using a resistance bridge. In the temperature interval from  $0^{\circ}$ C to  $30^{\circ}$ C, the resistance of the thermometer is approximated by a parabolic function of its Celsius temperature *θ*:

<span id="page-95-5"></span>
$$
R(\theta) = (1 + A\theta + B\theta^2)R_0,\tag{13.1}
$$

where  $R_0$ ,  $A$  and  $B$  are determined from a calibration of the thermometer. Estimates of  $R_0$ , *A* and *B* and the associated standard uncertainties are given in table [13.1,](#page-95-1) and the non-zero correlation coefficients associated with pairs of the estimates in table [13.2.](#page-96-0)

Table 13.1: Estimates of the input quantities  $X \equiv (R_0, A, B, R_S, r)^\top$  for the measurement of a single Celsius temperature and associated standard uncertainties [\(13.2.1,](#page-95-2) [13.2.2](#page-95-3) and [13.2.3\)](#page-95-4)

<span id="page-95-1"></span>

	$\widehat{A}/^{\circ}C^{-1}$ $\widehat{R}_0/\Omega$	$\widehat{B}/^{\circ}C^{-2}$	$\widehat{R}_S/\Omega$	$\widehat{r}/1$
Std une	Estimate $99.99610$ $0.0039096$ $-6.0 \times 10^{-7}$ $0.00050 - 0.0000027$	$1.1 \times 10^{-7}$	99.99947 0.00010	1.078 005 7 0.0000050

<span id="page-95-3"></span>**13.2.2** The estimate  $\widehat{R}_{\rm S}$  and the associated standard uncertainty, determined by calibra-tion, are given in table [13.1.](#page-95-1)  $R_{\rm S}$  is independent of the parameters  $R_0$ , A and B.

<span id="page-95-4"></span>**13.2.3** The quantity measured with the resistance bridge is the resistance ratio

<span id="page-95-6"></span>
$$
r = \frac{R(\theta)}{R_S}.
$$
\n(13.2)

The measured value  $\hat{r}$  and the associated standard uncertainty are given in table [13.1.](#page-95-1) The resistance ratio is independent of the parameters  $R_{0}$ ,  $A$  and  $B$  of the resistance thermometer

and the resistance  $R_S$  of the standard resistor; there are thus no further non-zero correlation coefficients beyond those given in table [13.2.](#page-96-0)

<span id="page-96-0"></span>Table 13.2: Non-zero correlation coefficients associated with pairs of estimates of the input quantities for the measurement of a single Celsius temperature (13.2.1, 13.2.2 and [13.2.3\)](#page-95-4)

<span id="page-96-1"></span>

**13.2.4** By combining expressions (13.1) and (13.2), the following measurement model for the Celsius temperature  $\theta$  is obtained:

$$
h(Y, X) = (1 + A\theta + B\theta^2)R_0 - rR_S = 0.
$$
\n(13.3)

The measurement model is an implicit univariate model with  $N = 5$  input quantities  $X \equiv (R_0, A, B, R_S, r)^\top$  and a single output quantity (measurand)  $Y = \theta$ .

<span id="page-96-3"></span>**13.2.5** Although the above measurement model can be solved for *θ* by considering the roots of a quadratic equation, the numerical evaluation of the transformed model can suffer from subtractive cancellation, and the evaluation of sensitivity coefficients is made more difficult by the transformation. Thus, measurement model 13.3 will be treated as implicit.

**13.2.6** The estimate  $y \equiv \hat{\theta}$  of the Celsius temperature generating the measured resistance ratio  $\hat{r}$  is found by inserting the estimates given in table 13.1 into equation [13.3](#page-96-1) and solving this equation for *θ*. The solution obtained is  $\theta = 20.0232$  °C.

**13.2.7** The standard uncertainty  $u_y \equiv u(\hat{\theta})$  associated with the estimate *y* is evaluated using [[7,](#page-126-3) expression (8)], namely

<span id="page-96-2"></span>
$$
C_y V_y C_y^\top = C_x V_x C_x^\top
$$
\n(13.4)

with the matrices of first-order partial derivatives given as follows:

$$
C_y = \frac{\partial h}{\partial Y} = \frac{\partial h}{\partial \theta} = (A + 2B\theta)R_0,
$$

and

$$
\mathcal{C}_x = \frac{\partial h}{\partial x} = \left[ \frac{\partial h}{\partial R_0}, \frac{\partial h}{\partial A}, \frac{\partial h}{\partial B}, \frac{\partial h}{\partial R_S}, \frac{\partial h}{\partial r} \right] = \left[ 1 + A\theta + B\theta^2, R_0\theta, R_0\theta^2, -r, -R_S \right].
$$

These expressions are evaluated at the estimates of the input quantities given in table [13.1](#page-95-1) and the corresponding estimate of the output quantity, which gives the following sensitivity matrices:

and the resistance *R*<sub>S</sub> of the standard resistor; there are thus no further non-zero correlation coefficients beyond those given in table 13.2.  
\nTable 13.2: Non-zero correlation coefficients associated with pairs of estimates of the input quantities for the measurement of a single Celsius temperature (13.2.1, 13.2.2 and 13.2.3)  
\n
$$
\frac{\widehat{R}_0}{\widehat{R}_0} \frac{\widehat{A}}{1} \frac{\widehat{B}}{1} -0.155 \frac{0.092}{-0.959}
$$
\n
$$
\frac{\widehat{B}}{1} \frac{1}{1} -0.959
$$
\n13.2.4 By combining expressions (13.1) and (13.2), the following measurement model for the Celsius temperature θ is obtained:  
\n
$$
h(Y, X) = (1 + A\theta + B\theta^2)R_0 - rR_5 = 0.
$$
\n(13.3)  
\nThe measurement model is an implicit univariate model with  $N = 5$  input quantities  $X = (R_0, A, B, R_5, r)^T$  and a single output quantity (measurable with  $N = \theta$ ).  
\n13.2.5 Although the above measurement model can be solved for θ by considering the roots of a quadratic equation, the numerical evaluation of the transformed model can suffer from subtractive cancellation, and the evaluating normal distribution of the transformation. Thus, measurement model 13.3 will be treated as implicit.  
\n13.2.6 The estimate  $y = \widehat{\theta}$  of the Celsius temperature generating the measured resistivity toefficient by the transformation. Thus, measurement model 13.3 will be treated as implicit.  
\n13.2.7 The standard uncertainty  $u_{xy} = u(\widehat{\theta})$  associated with the estimate  $y$  is evaluated using [7, expression (8)], namely  
\n
$$
C_y V_y C_y^\top = C_x V_x C_x^\top
$$
\n
$$
V_x = \frac{\partial h}{\partial V} = \frac{\partial h}{\partial \theta} = (A + 2B\theta)R_0,
$$
\nand  
\n
$$
C_x = \frac{\partial h}{\partial X} = \frac{\partial h}{\partial R_0}, \frac{\partial h}{\partial A}, \frac{\partial h}{\partial B}, \frac{\partial h}{\partial R_0}, \frac{\partial h}{\partial T_0} = [1 + A\theta + B\theta^2, R_0\theta, R_0\theta^2, -r, -R_5].
$$
\nThese expressions are evaluated at the estimates of the input quantities given in table 13.1 and the corresponding estimate of the output quantity, which gives the following sensitivity matrices:  
\n
$$
C
$$

#### <span id="page-97-2"></span><span id="page-97-0"></span>**13.3 Measurement of several Celsius temperatures**

<span id="page-97-1"></span>Table 13.3: Estimates of the resistance ratios for the measurement of several Celsius temperatures and associated standard uncertainties (13.3.2)

	s $u(\hat{\theta}) = 0.0045$ °C.	The elements of the covariance matrix $V_x$ associated with the estimates of the input quan- ities are calculated from the standard uncertainties in table 13.1 and the correlation coefficients in table 13.2. The result obtained after matrix algebra of expression (13.4)					
		13.3 Measurement of several Celsius temperatures					
13.3.1	beratures $\theta_1, \ldots, \theta_{10}$ .	The resistance thermometer, standard resistor and resistance bridge described n 13.2 are used to measure the resistance ratios $r_1, \ldots, r_{10}$ generated by ten Celsius tem-					
		13.3.2 The estimates $\widehat{R}_0$ , $\widehat{A}$ , $\widehat{B}$ and $\widehat{R}_S$ , and the associated standard uncertainties remain as before (table 13.1) and the estimates $\widehat{r}_1,\ldots,\widehat{r}_{10}$ and their associated standard uncertainties are shown in table 13.3. The only non-zero correlation coefficients associated with pairs of the estimates remain those given in table 13.2. The resistance ratios are assumed to be independent, an assumption that is valid if the magnitudes of the random errors in the neasured resistance ratios dominate. Table 13.3: Estimates of the resistance ratios for the measurement of several Celsius tem- eratures and associated standard uncertainties (13.3.2)					
		$\widehat{r}_1$	$\widehat{r}_2$	$\widehat{r}_3$	$\widehat{r}_4$	$\widehat{r}_5$	
	Estimate Std unc	1.000 0053 0.0000050	1.015 0054 0.0000050	1.0300055 0.0000050	1.045 005 6 0.0000050	1.0600056 0.0000050	
		$\widehat{r}_6$	$\widehat{r}_7$	$\widehat{r}_8$	$\widehat{r}_9$	$\widehat{r}_{10}$	
	Estimate Std unc	1.078 0057 0.0000050	1.090 0058 0.0000050	1.1050059 0.0000050	1.1200060 0.0000050	1.0780057 0.0000050	
	neasurement model:	13.3.3 Each resistance ratio $r_i$ is related to the corresponding Celsius temperature $\theta_i$ by in expression of the form (13.3): $(1 + A\theta_j + B\theta_j^2)R_0 - r_jR_S = 0, \quad j = 1, , 10.$ There are $N=14$ input quantities $\pmb{X}\equiv(R_0,A,B,R_{\rm S},r_1,\ldots,r_{10})^\top$ and $m=10$ output quan- ities $Y \equiv (\theta_1, , \theta_{10})^\top$ , all of which are related by the following implicit multivariate					(13.5)
		$\begin{bmatrix} h_1(Y,X) \\ \vdots \\ h_{10}(Y,X) \end{bmatrix} \equiv \begin{bmatrix} R_0 \left( 1 + A\theta_1 + B\theta_1^2 \right) - r_1 R_S \\ \vdots \\ R_0 \left( 1 + A\theta_{10} + B\theta_{10}^2 \right) - r_{10} R_S \end{bmatrix} = 0.$					
		As noted in clause 13.2.5, transformation of the implicit measurement model (13.5) into n explicit form is possible but inadvisable.					
		1 <b>3.3.4</b> The estimates $\mathbf{y} = (\widehat{\theta}_1, \dots, \widehat{\theta}_{10})$ of the Celsius temperatures $Y$ are given by insert- ng the estimates given in columns 1 to 4 of table 13.1 and in table 13.3 into equations 13.5) and solving these equations. These estimates are given in table 13.4.					

<span id="page-97-3"></span>
$$
(1 + A\theta_j + B\theta_j^2)R_0 - r_jR_S = 0, \qquad j = 1, ..., 10.
$$
 (13.5)

<span id="page-97-4"></span>
$$
\begin{bmatrix} h_1(Y,X) \\ \vdots \\ h_{10}(Y,X) \end{bmatrix} \equiv \begin{bmatrix} R_0 \left( 1 + A\theta_1 + B\theta_1^2 \right) - r_1 R_S \\ \vdots \\ R_0 \left( 1 + A\theta_{10} + B\theta_{10}^2 \right) - r_{10} R_S \end{bmatrix} = 0.
$$

<span id="page-98-0"></span>Table 13.4: Estimates of the output quantities *Y* and associated standard uncertainties for the measurement of several Celsius temperatures [\(13.3.4](#page-97-4) and [13.3.5\)](#page-98-1)

∼			$\theta_{\epsilon}$	$\theta_{\rm s}$	$\theta_{10}$ $\theta_{\rm o}$
					Estimate/°C 0.0100 3.8491 7.6928 11.5410 15.3938 20.0232 23.1131 26.9797 30.8509 20.0232
					Std. unc./°C 0.0018 0.0027 0.0040 0.0046 0.0047 0.0045 0.0046 0.0060 0.0089 0.0045

<span id="page-98-2"></span>Table 13.5: Correlation coefficients associated with pairs of estimates of the output quantities *Y* for the measurement of several Celsius temperatures [\(13.3.5\)](#page-98-1)



<span id="page-98-1"></span>**13.3.5** The covariance matrix  $C_v$  associated with y is evaluated using expression [\(13.4\)](#page-96-2) where  $\bm{C}_{\bm{y}}$  and  $\bm{C}_{\bm{x}}$  are sensitivity matrices given by evaluating  $\bm{C}_{\bm{Y}}$  and  $\bm{C}_{\bm{X}}$ , respectively, at the estimates of the input and output quantities.  $\boldsymbol{c}_{\boldsymbol{\gamma}}$  is a diagonal matrix of dimension  $10\times$ 10 with diagonal entries  $R_0(A+2B\theta_1), \ldots, R_0(A+2B\theta_{10})$ .  $C_X$  is a matrix of dimension 10  $\times$ 14 given by

$$
C_X = \left[ \begin{matrix} C_X^{(1)} & C_X^{(2)} \end{matrix} \right],
$$

where

$$
\mathbf{C}_{X}^{(1)} = \left[ \begin{array}{cccc} 1 + A\theta_1 + B\theta_1^2 & R_0\theta_1 & R_0\theta_1^2 & -r_1 \\ \vdots & \vdots & \vdots & \vdots \\ 1 + A\theta_{10} + B\theta_{10}^2 & R_0\theta_{10} & R_0\theta_{10}^2 & -r_{10} \end{array} \right]
$$

is a matrix of dimension  $10 \times 4$ , and  $C_X^{(2)}$  is a diagonal matrix of dimension  $10 \times 10$  with diagonal elements all equal to  $-R_s$ . The covariance matrix  $V_x$  is calculated from the standard uncertainties given in columns 1 to 4 of table [13.1](#page-95-1) and in table [13.2.](#page-96-0) The standard uncertainties associated with the estimates of the Celsius temperatures and the correlation coefficients associated with pairs of estimates, derived from the matrix *C <sup>y</sup>* , are given in table 13.4 and table [13.5,](#page-98-2) respectively. A numerically stable method to form the covariance matrix *C <sup>y</sup>* is given in [[7,](#page-126-3) annex B].

<span id="page-98-3"></span>**13.3.6** The results given in table [13.4](#page-98-0) and figure [13.1](#page-99-0) show how the standard uncertainty  $u(\theta_j)$  varies with the estimate  $\theta_j$  of Celsius temperature  $\theta_j$ . The uncertainty is smallest around  $0^{\circ}$ C and increases rapidly for temperatures greater than 25 °C. This effect is due to the fact that the resistance thermometer was calibrated at the temperatures 0  $^{\circ}$ C, 15  $^{\circ}$ C, 20  $\degree$ C and 25  $\degree$ C, and that the Celsius temperature 0  $\degree$ C during calibration was generated using an ice bath with a standard uncertainty three times smaller than those associated with the other three temperatures, which were generated using an oil bath.

NOTE In figures [13.1](#page-99-0) and [13.2](#page-100-0) the straight-line segments joining the plotted points are included for purpose of visualization.

<span id="page-99-1"></span>**13.3.7** Using the results given in the final column of table [13.5,](#page-98-2) figure 13.2 shows how the correlation coefficient associated with the pair of estimates  $\theta_j$  and  $\theta_{10}$  of Celsius temperatures  $\theta_j$  and  $\theta_{10} = 20$  °C varies with  $\widehat{\theta}_j$ ,  $j = 1, \ldots, 9$ . The correlation coefficient has a maximum for  $\widehat{\theta}_j = \widehat{\theta}_6$  and approaches zero as the absolute difference  $|\widehat{\theta}_j - \widehat{\theta}_{10}|$  becomes large. The example demonstrates that quantities measured with the same instrument can be highly correlated.



<span id="page-99-0"></span>Figure 13.1: Standard uncertainty  $u(\theta_j)$  associated with the estimate  $\theta_j$  of Celsius temperature *θ<sup>j</sup>* [\(13.3.6\)](#page-98-3)



<span id="page-100-0"></span>Figure 13.2: Correlation coefficient associated with the pair of estimates  $\hat{\theta}_j$  and  $\hat{\theta}_{10}$  of Celsius temperatures  $\theta_j$  and  $\theta_{10} = 20$  °C (13.3.7)



# **14 Activity of a radioactive source corrected for decay**

#### **14.1 Preamble**

**14.1.1** A <sup>207</sup>Bi (bismuth-207) standard is used at a time  $t_1$  to calibrate a detector. The calibration certificate of this radioactive source reports an activity  $\widehat{A}_0$  at the reference date  $t_0$ , with associated standard uncertainty  $u(A_0)$ .

**14.1.2** This example provides an estimate  $\widehat{A}_1$  of the activity  $A_1$  at time  $t_1$  and evaluates the associated standard uncertainty  $u(A_1)$ , illustrating the application of the law of propagation of uncertainty (LPU). Because the model is non-linear in time and half-life  $T_{1/2}$ , the example also verifies whether the LPU is adequate for this task.

#### **14.2 Measurement model**

The measurement model for the activity  $A_1$  at time  $t_1$  is

<span id="page-101-0"></span>
$$
A_1 = f(A_0, T_{1/2}, \Delta t) = A_0 \exp\left(-\frac{\Delta t \ln 2}{T_{1/2}}\right),\tag{14.1}
$$

with  $\Delta t = t_1 - t_0$ . Both quantities  $t_0$  and  $t_1$ , and thus  $\Delta t$ , are considered to have negligible uncertainty.

Introducing the *decay correction E* =  $\exp\left(-\frac{\Delta t \ln 2}{T_{1/2}}\right)$ *T*1*/*<sup>2</sup> ,

$$
A_1 = A_0 E. \tag{14.2}
$$

#### **14.3 Uncertainty evaluation**

**14.3.1** Estimates  $\widehat{A}_0$  and  $\widehat{T}_{1/2}$  of  $A_0$  and  $T_{1/2}$  and their associated standard uncertainties are taken from a calibration certificate and the literature [18], respectively:

$$
\widehat{A}_0 = 11.450 \text{ kBq},
$$
  $u(\widehat{A}_0) = 0.050 \text{ kBq},$   
\n $\widehat{T}_{1/2} = 32.9 \text{ a},$   $u(\widehat{T}_{1/2}) = 1.4 \text{ a}$ 

and *∆t* is taken as *∆t* = 2.73 a. The symbol a is used for the year in accordance with ISO 80000-3:2020 [[65](#page-131-6)]. No further information is taken into consideration.

**14.3.2** Substitution of the above values into measurement model [\(14.1\)](#page-101-0) gives

$$
\widehat{A}_1 = 11.450 \text{ kBq} \times \exp\left(\frac{-2.73 \text{ a} \times \text{ln}2}{32.9 \text{ a}}\right) = 0.94411 \widehat{A}_0 = 10.810 \text{ kBq}.
$$

Applying the LPU [6, Clauses 5.1.2, 5.1.3] to measurement model [\(14.1\)](#page-101-0) gives

$$
u^{2}(\widehat{A}_{1})=u_{\widehat{A}_{0}}^{2}(\widehat{A}_{1})+u_{\widehat{T}_{1/2}}^{2}(\widehat{A}_{1}) \text{ with } u_{\widehat{A}_{0}}(\widehat{A}_{1})=|c_{\widehat{A}_{0}}|u(\widehat{A}_{0}), u_{\widehat{T}_{1/2}}(\widehat{A}_{1})=|c_{\widehat{T}_{1/2}}|u(\widehat{T}_{1/2}).
$$

 $c_{\widehat{A}_0}$  and  $c_{\widehat{T}_1/2}$  are sensitivity coefficients given by differentiating the measurement model [\(14.1\)](#page-101-0), namely,

**14. Activity of a radioactive source corrected for decay**  
\n**14.1 Preamble**  
\n**14.1.1** A <sup>207</sup>Bi (bismuth-207) standard is used at a time 
$$
t_1
$$
 to callibrate a detector. The  
\nallibration certificate of this radioactive source reports an activity  $\hat{A}_0$  at the reference date  
\n<sub>0</sub>, with associated standard uncertainty  $u(\hat{A}_0)$ .  
\n**14.1.2** This example provides an estimate  $\hat{A}_1$  of the activity  $A_1$  at time  $t_1$  and evaluates  
\nthe associated standard uncertainty  $u(\hat{A}_0)$ . Illustrating the application of the law of propa-  
\nization of uncertainty (LPU). Because the model is non-linear in time and half-life  $T_{1/2}$ , the  
\nexample also verifies whether the LPU is adequate for this task.  
\n**14.2 Measurement model**  
\n
$$
A_1 = f(A_0, T_{1/2}, \Delta t) = A_0 \exp\left(-\frac{\Delta t \ln 2}{T_{1/2}}\right),
$$
\n
$$
A_1 = f(A_0, T_{1/2}, \Delta t) = A_0 \exp\left(-\frac{\Delta t \ln 2}{T_{1/2}}\right).
$$
\n
$$
A_1 = A_0 E.
$$
\n**14.3 Uncertainty evaluation**  
\n**14.3.1** Estimates  $\hat{A}_0$  and  $\hat{T}_{1/2}$  of  $A_0$  and  $T_{1/2}$  and their associated standard uncertainties  
\nrate taken from a calibration certificate and the literature [18], respectively:  
\n $\hat{A}_0 = 11.450 \text{ kBq}, \quad u(\hat{A}_0) = 0.050 \text{ kBq},$   
\n $\hat{T}_{1/2} = 32.9 \text{ a}, \qquad u(\hat{T}_{1/2}) = 1.4 \text{ a}$   
\n**14.3.2** Substitution of the above values into measurement model (14.1) gives  
\n $\hat{A}_1 = 11.450 \text{ kBq} \times \exp\left(-\frac{2.73 \text{ a} \times \ln 2}{32.9 \text{ a}}\right) = 0.94411 \hat{A}_0 = 10.810 \text{ kBq}.$ 

where the partial derivatives are evaluated at the estimates  $A_0$  and  $T_{1/2}$ . Thus,

$$
u_{\widehat{A}_0}(\widehat{A}_1) = 0.94411 \times 0.050 \text{ kBq} = 0.0472 \text{ kBq},
$$
  

$$
u_{T_{1/2}}(\widehat{A}_1) = 0.018898 \text{ kBq} \text{ a}^{-1} \times 1.4 \text{ a} = 0.0265 \text{ kBq}
$$

and hence

$$
u(\widehat{A}_1) = \left[ (0.0472)^2 + (0.0265)^2 \right]^{1/2} \text{kBq} = 0.054 \text{ kBq}.
$$

**14.3.3** To investigate whether the LPU is adequate in this case, the behaviour of the non-linear model [\(14.1\)](#page-101-0) is compared with the linearized model in the neighbourhood of  $\widehat{T}_{1/2}$  = 32.9 a. The linearized model is

<span id="page-102-0"></span>
$$
A'_{1} = f_{\text{lin}}(A_{0}, T_{1/2}) = \widehat{A}_{1} + c_{\widehat{A}_{0}}(A_{0} - \widehat{A}_{0}) + c_{\widehat{T}_{1/2}}(T_{1/2} - \widehat{T}_{1/2}). \tag{14.3}
$$

where the partial derivatives are evaluated at the estimates  $\vec{A}_6$  and  $\vec{T}_{ijx}$ . Thus,  $a_{\vec{X}_1}(\vec{A}_1) = 0.94441 \times 0.050 \text{ kJg} = 0.0472 \text{ kJt}_0$ <br>  $a_{\vec{X}_{12}}(\vec{A}_1) = 0.018898 \text{ kJg} = 4 \times 1.4a = 0.0265 \text{ kJg}$ <br>
and hen The criterion used here for adequacy of LPU is that the values of the non-linear model [\(14.1\)](#page-101-0) and the linearized model [\(14.3\)](#page-102-0) differ negligibly in the neighbourhood of  $T_{1/2}$ . This neighbourhood is taken as comprising values of the half-life  $T_{1/2}$  up to several multiples of  $u(T_{1/2})$  away from its estimate. Accordingly, the linearized and the exact model values are tabulated at several positive and negative multiples of  $u(T_{1/2}) = 1.4$  a away from  $\hat{T}_{1/2} = 32.9$  a. See table [14.1.](#page-102-1) The approach used is comparable to that in [[10,](#page-126-4) annex F.2]; see also [[6,](#page-126-1) clause 5.1.3, note 2].

Table 14.1: Estimates  $\widehat{A}'_1$  and  $\widehat{A}_1$  from linearized and exact models for <sup>207</sup>Bi decay, respectively

<span id="page-102-1"></span>

$\widehat{T}_{1/2}/a$	$\widehat{A}'_1$ /kBq	$\widehat{A}_1$ /kBq	Difference/kBq
28.7	10.731	10.719	0.011
30.1	10.757	10.752	0.005
31.5	10.784	10.782	0.001
32.9	10.810	10.810	0.000
34.3	10.836	10.835	0.001
35.7	10.863	10.859	0.004
37.1	10.889	10.881	0.009

**14.3.4** As seen from table [14.1,](#page-102-1) the values of the linearized model differ from the nonlinear model by at most 0.011 kBq for values of the half-life between  $T_{1/2} - 3u(T_{1/2})$ and  $T_{1/2} + 3u(T_{1/2})$ . This difference is considerably smaller than  $u(A_0) = 0.050$  kBq and  $u(A_1) = 0.054$  kBq and therefore it can be concluded that the validity of LPU is confirmed in this case.

### **14.4 Discussion**

**14.4.1** The time interval *∆t* = 2.73 a is small compared with the radionuclide half life  $\widehat{T}_{1/2}$  = 32.9 a. As a consequence, the value 0.944 11 of *E* in expression [\(14.1\)](#page-101-0) is reasonably close to unity, and so  $u(A_1)$  differs little from  $u(A_0)$ .

JCA 2. When the time difference As becomes close to  $\hat{v}_{1/2}$ ,  $|\mathbf{U}|$  may not popied set to  $\hat{v}_{2/2}$  and the interference of the interference of  $\hat{v}_{1/2}$  and the range of the interference of  $\hat{v}_{2/2}$  and the **14.4.2** When the time difference  $\Delta t$  becomes closer to  $T_{1/2}$ , LPU may not provide acceptable results for <sup>207</sup>Bi as judged by such a check. This happens because the relative uncertainty of its half-life of 4% is large and the range of values from  $\widehat{T}_{1/2} - 3u(\widehat{T}_{1/2})$  to  $T_{1/2}$  + 3*u*( $T_{1/2}$ ) is more likely to highlight non-linearities. However, in most cases encountered in radionuclide metrology and its applications, half-lives are known more precisely (relative uncertainty of 0.5 % or better) and LPU can be used safely unless *<sup>∆</sup>t/T*b1*/*<sup>2</sup> is very large like in  $^{14}$ C dating of very old samples.

When  $\Delta t / T_{1/2}$  and the relative uncertainty of the half-life *u*(*T*<sub>1/2</sub>)/*T*<sub>1/2</sub> are both sufficiently small, LPU can be used safely. In practice, the relative uncertainty of the decay correction

$$
u(E)/E = \ln 2 \left[ u(\hat{T}_{1/2})/\hat{T}_{1/2} \right] \Delta t/\hat{T}_{1/2},
$$

is a good measure of the non-linearity of the model. It can be shown that in general, when the relative uncertainty of the decay correction  $u(E)/E$  reaches 1%, it is important to check the validity of LPU as shown in the present example and compare the difference between the linearized and exact models with the uncertainty of the activity in order to decide whether the non-linear effect on LPU is significant or not. In the affirmative, JCGM 101:2008 [[12](#page-126-2)], concerned with the propagation of distributions using a Monte Carlo method, should be applied.

# **15 Breaking force of steel wire rope**

### **15.1 Preamble**

With thousands of workers transported up and down mine shafts each day (figure 15.1), measuring the safety of hoisting ropes forms an integral part of mine management. Strict laws cover the maintenance and testing of steel wire rope, the latter focussing mainly on the destructive testing of these ropes. The source of this information and the example is the Rope Testing Laboratory of the Council for Scientific and Industrial Research, South Africa (CSIR), from whom permission was obtained to use the example.



Figure 15.1: Steel wire rope used for hoisting in vertical mine shafts

## <span id="page-104-0"></span>**15.2 Measurement model**

**15.2.1** The measurand is the breaking force of a sample constituting the length of a piece of rope of a particular type. An observation of the breaking force cannot be repeated since the test is destructive on the sample. Because it is impossible to make a series of observations of the measurand under repeatability conditions, one observation is instead made of each of a number, *n*, of samples given by cutting pieces from the same rope. The pragmatic approach is therefore taken of testing six pieces from the same coil of rope and accepting that the variability across the six specimens will be embedded in the variability of the test results.

**15.2.2** The breaking forces that are recorded are influenced by two effects: the variability of the samples due to the inhomogeneity of the rope, and intrinsic measurement effects, essentially measurement corrections, such as resolution of the force measuring system.

Following the provisions of JCGM GUM-6:2020 [[10](#page-126-4)],  $n = 6$  breaking force observation equations are

<span id="page-104-1"></span>
$$
F_i = F_0 + B + B_{\text{ref}} + e_{\text{res}} + e_T + E_i, \qquad i = 1, ..., n,
$$
\n(15.1)

where

the  $F_i$  denotes quantities for which the breaking force observations are independent realizations,

 $F_0$  is the measurand, breaking force,

*B* is the error of the test machine load cell,

 $B_{ref}$  is the error of the reference load cell,

 $e_{res}$  is the resolution error of the force measuring system,

 $e_T$  is the temperature error,

 $E_i$  are the measurement errors associated with the  $F_i$ .

### **15.3 Assignment of PDFs**

#### <span id="page-105-1"></span>**15.3.1 Breaking force observations**

The available data comprise  $n = 6$  observations of the  $F_{i_2}$ 

 $\widehat{\textbf{F}}^{\top} = [\widehat{F}_1, \ldots, \widehat{F}_6] = (10\,006\ 10\,007\ 10\,005\ 10\,008\ 10\,003\ 10\,005)$ kN,

with associated errors  $E_i$  assumed to be independent, normal random variables with mean 0 kN and unknown standard deviation *σ*.

The data  $F_i$  have mean 10 005.7 kN and standard deviation 1.8 kN. This mean has associated standard uncertainty 1.8 kN*/* 6 = 0.7 kN. Figure 15.2 shows a normal distribution with mean 10 005.7 kN and standard deviation 0.7 kN as the full blue curve. An alternative choice of distribution, the Weibull, displayed as the broken black curve, is considered in clause [15.6.](#page-107-0)



<span id="page-105-0"></span>Figure 15.2: Normal distribution (full blue curve) and Weibull distribution (broken black curve) fitted to the breaking force data

Using the provisions of  $[12,$  $[12,$  clause 6.4.7]), the mean 10005.7 kN and standard deviation 0.7 kN with  $\nu = n - 1 = 5$  degrees of freedom are used as the parameters of the resulting *t* distribution for breaking force *F*.

## **15.3.2 Correction effects**

All other input quantities, representing measurement corrections, are modelled on the basis of available information with rectangular distributions centred at zero and having specified semi-widths: see table [15.1.](#page-106-0)

<span id="page-106-0"></span>

<b>Ouantity</b>		Semi-width/kN Standard uncertainty/kN	
В	6.8	3.9	
$B_{\text{ref}}$	8.2	4.7	
$e_{\text{res}}$	8.7	5.0	
$e_T$	10.2.	5.9	

Table 15.1: Rectangular PDFs assigned to correction effects for the rope breaking-force model on the basis of available information

### <span id="page-106-1"></span>**15.4 JCGM 100:2008 uncertainty framework (GUF)**

**15.4.1** The measurand  $F_0$  in expression (15.1) does not appear as the subject (compare [[91](#page-133-0)]). Making it so as in [[6](#page-126-1)] yields the measurement model

$$
F_0 = F - B - B_{\text{ref}} - e_{\text{res}} - e_T - E,\tag{15.2}
$$

using which the estimate of  $F_0$  is 10005.6 kN and applying the law of propagation of uncertainty (LPU) gives, using the standard uncertainty associated with the *E<sup>i</sup>* as identical to that of the  $F_i$  and hatted symbols as denoting the estimates of the various quantities,

$$
u(\widehat{F}_0) = \left[ u^2(\widehat{F}) + u^2(\widehat{B}) + u^2(\widehat{B}_{\text{ref}}) + u^2(\widehat{e}_{\text{res}}) + u^2(\widehat{e}_T) + + u^2(\widehat{E}) \right]^{1/2} = 9.9 \,\text{kN}.
$$

Since all but one of the input quantities have infinite degrees of freedom supporting their standard uncertainty estimates, the Welch-Satterthwaite formula [6, expression (G.2a)] in JCGM 100:2008 yields

$$
v_{\text{eff}} = (n-1) \left[ \frac{u(\widehat{F}_0)}{u(\widehat{F})} \right]^4 > 10^5.
$$

As a consequence, the measurement result is a breaking force of  $\hat{F} = 10005.6 \text{ kN}$  with an expanded uncertainty  $U(\widehat{F}) = 19.4$  kN corresponding to a coverage factor *k* of 1.96 at a level of confidence of 95 %.

**15.4.2** Following the provisions of JCGM 100:2008, a *t*-distribution with  $v_{\text{eff}}$  degrees of freedom [[6,](#page-126-1) annex G.3] is assigned as the PDF for  $F_0.$  Since  $\,_{\rm eff}$  is so large, the t distribution could be taken as normal for practical purposes, as demonstrated by the coverage factor being equal to three significant decimal digits to that for a normal distribution for the same level of confidence.

## **15.5 Monte Carlo method (MCM)**

Table 15.1: Rectangular PDFs ansigned to correction effects for the rope bracklage<br>Sigmoid of the also is available information  $\frac{8.9}{\epsilon_0}$ <br>  $\frac{8.2}{\epsilon_1}$ <br>  $\frac{8.2}{\epsilon_2}$ <br>  $\frac{8.2}{\epsilon_3}$ <br>  $\frac{8.2}{\epsilon_4}$ <br>  $\frac{8.2}{\epsilon_5}$ **15.5.1** To validate the results of clause [15.4,](#page-106-1) MCM was applied following the provisions of [12]. Its application in this example consists of taking  $10^6$  draws from the probability distributions considered in clause [15.4](#page-106-1) and using the measurement model for each set of draws to provide a corresponding value for the measurand  $F_0.$  For this simple measurement model, the Monte Carlo calculation to provide this sample of the probability distribution for *F*0 took 0.2 s using MATLAB version 'R2023b Update 3 (23.2.0.2409890) 64-bit (win64) October 4, 2023'.

**15.5.2** The mean and standard deviation of this MCM sample of  $F_0$  are 10 005.7 kN and 10.1 kN, respectively, and the equiprobabilistic 95 % coverage interval whose endpoints are the 2.5 % and 97.5 % quantiles of the sample ranged from 9986.0 kN to 10 025.3 kN. Since this interval is almost exactly symmetric about the mean, half the length of this interval is taken as an expanded uncertainty, namely,  $U(\widehat{F}_0) \approx 19.6 \, \text{kN}$ .

**15.5.3** These results were confirmed to the number of significant decimal digits quoted by re-running the Monte Carlo calculation with  $10^7$  trials, which took  $1.6\,\mathrm{s}$ .

## <span id="page-107-0"></span>**15.6 Discussion**

**15.6.1** The results provided by the GUF and MCM, having respective estimates 10 005.6 kN and 10 005.7 kN and associated standard uncertainties 9.9 kN and 10.1 kN compare well, although the GUF probability density function (PDF) is slightly more peaked than that provided by MCM: figure [15.3.](#page-107-1) Since the Monte Carlo method can be regarded as a reference [[33](#page-128-8)], judging by the considerable extent of agreement with the results from applying JCGM 100:2008 it can be concluded that the latter for the above set of data is valid.



<span id="page-107-1"></span>Figure 15.3: PDF for the breaking force obtained using the GUF (blue curve) and scaled histogram from MCM using the information summarized in table [15.1](#page-106-0)

**15.5.2.** The mean and standard desiration of this MCM sample of  $P_0$  are 10.006.74N sample in the second state in the second state in the second state in the second of this MCM sample ranged from 9986.61W in 10.025.3 Ne **15.6.2** A normal distribution was assigned to the breaking force *F* in clause [15.3.1](#page-105-1) since it is common practice in testing laboratories when preparing uncertainty budgets in the area. However, for destructive testing, as here, a [Weibull distribution](https://en.wikipedia.org/wiki/Weibull_distribution) is often used especially in research studies in the area [[74,](#page-131-7) [83](#page-132-4)]. Accordingly, a Weibull distribution was alternatively assigned to the quantity *F* for this purpose. Typically, such a distribution is conventionally fitted to the data for *F* and the mean and standard deviation extracted, which yields in this case exactly the same mean and standard uncertainty as for the normal distribution. There are readily discernible differences between the distributions, shown in figure 15.2: for instance, whilst the normal distribution is symmetric about its mean, the Weibull distribution displays some asymmetry.

**15.6.3** The question may be raised as to what are the conditions in which to use the Weibull distribution instead of the normal. For the data analyzed here, the uncertainty
contribution from the six observations was virtually negligible compared with those from the correction effects. That contribution, 0.7 kN is to be compared with a the standard uncertainty, about 10 kN, associated with the estimates of the measurand. Bearing in mind that the standard uncertainties associated with the input estimates are combined in quadrature, that contribution influences the combined standard uncertainty by only 0.2 %.

<span id="page-108-0"></span>contribution from th[e](#page-126-0) six doserocions was witnedly negligible compared with the phosic from the correction effects. That contribution, 6, 2 NH is no her operated with the drained of the mindston that the standard uncertain **15.6.4** In cases where this contribution is more meaningful, it is recommended to use the Weibull distribution, which seems state of the art in the area [[83](#page-132-0)]. It is to be noted that since the means of these fitted distributions are identical, as are their standard deviations, the GUF would yield identical results for the two distributions. Conversely, MCM would deliver different results since the PDF for the first input quantity would be different. More trust could be placed in the MCM results since all conditions [12, clause 5.7] required for the GUF to apply would not hold.

## **16 Comparison loss in microwave power meter calibration**

### **16.1 Preamble**

**16.1.1** This example illustrates the evaluation of measurement uncertainty for a power ratio that can occur during a microwave power meter calibration. Four methods for evaluating the uncertainty are considered:

- 1. JCGM 100:2008 uncertainty framework (GUF) with first-order terms (GUF<sub>1</sub>);
- 2. GUF with higher-order terms  $(GUF<sub>2</sub>)$ ;
- 3. Monte Carlo method (MCM);
- 4. analytic calculation.

<span id="page-109-2"></span>16 **Comparison loos in microwave power meter calibration**<br>
16.1 **Preamble**<br>
16.1 **Preamble**<br>
16.1 This example illustrates the evaluation of measurement uncertainty for a power<br>
16.1 This example illustrates the evaluatio The example includes scenarios where the Central Limit Theorem (CLT) [6, Annex G.2] holds from acceptably to poorly, or does not hold at all. Whilst the use of GUF yields correct estimates and standard uncertainties – provided that higher-order terms are taken into account – it does not in general provide reliable coverage intervals. MCM instead provides realistic uncertainties and coverage intervals for all scenarios. Clause 16.5 contains some information that is relevant to the analytic calculations presented in this example that may be of interest to advanced practitioners of uncertainty calculations. This information can be disregarded by other (more general) practitioners, who can therefore ignore references in the example to that clause.

**16.1.2** Microwave power meter calibration involves the measurement of several voltage reflection coefficients. Voltage reflection coefficients are complex-valued quantities. Manufacturers of microwave power meters usually try to produce power meters where these voltage reflection coefficients are as close to  $0+$ j0, where j<sup>2</sup> = −1, as possible. In practice, the absolute values for both the real and imaginary components of these voltage reflection coefficients are expected to be less than 0.050. Hence, this example considers absolute values for these components between zero and 0.050.

## **16.2 Formulation**

**16.2.1** During the calibration of a microwave power meter, the power meter and a standard power meter are connected in turn to a stable signal generator. The power absorbed by each meter will in general be different because their complex input voltage reflection coefficients are not identical. The ratio  $Y$  of the power  $P_M$  absorbed by the meter being calibrated and that,  $P_{\rm S}$ , by the standard meter is  $[98]$  $[98]$  $[98]$ 

<span id="page-109-0"></span>
$$
Y = \frac{P_{\rm M}}{P_{\rm S}} = \frac{1 - | \Gamma_{\rm M} |^2}{1 - | \Gamma_{\rm S} |^2} \times \frac{|1 - \Gamma_{\rm S} \Gamma_{\rm G}|^2}{|1 - \Gamma_{\rm M} \Gamma_{\rm G}|^2},\tag{16.1}
$$

where  $\Gamma_G$  is the voltage reflection coefficient of the signal generator,  $\Gamma_M$  that of the meter being calibrated and  $\varGamma_{\text{S}}$  that of the standard meter. This power ratio is an instance of "comparison loss" [5, 67].

**16.2.2** Consider the case where the standard and the signal generator are reflectionless, that is,  $\Gamma_{\rm S} = \Gamma_{\rm G} = 0 + i0$ , and measured values are obtained of the real and imaginary parts *X*<sub>1</sub> and *X*<sub>2</sub> of *Γ*<sub>M</sub> = *X*<sub>1</sub> + j*X*<sub>2</sub>. Since  $|I_M|^2 = X_1^2 + X_2^2$ , formula [\(16.1\)](#page-109-0) becomes

<span id="page-109-1"></span>
$$
Y = 1 - X_1^2 - X_2^2. \tag{16.2}
$$

**16.2.3** Given respectively are estimates  $x_1$  and  $x_2$  of the quantities  $X_1$  and  $X_2$  from measurement and the associated standard uncertainties  $u(x_1)$  and  $u(x_2)$ .  $X_1$  and  $X_2$  are often not independent. Denote by  $u(x_1, x_2)$  the covariance associated with  $x_1$  and  $x_2$ . Equivalently,  $u(x_1, x_2) = r(x_1, x_2)u(x_1)u(x_2)$ , where  $r = r(x_1, x_2)$  denotes the associated corre-lation coefficient [[6,](#page-126-1) clause 5.2.2].

NOTE The practitioner may sometimes have difficulty in quantifying the covariance. In such cases, the uncertainty evaluation can be repeated with various trial numerical values for the correlation coefficient in order to study its effect. This example carries out calculations using a correlation coefficient of zero and of 0.9 (compare [16.2.7\)](#page-110-0).

**16.2.4** On the basis of [[12,](#page-126-0) clause 6.4.8],  $X = (X_1, X_2)^\top$  is assigned a bivariate normal probability density function (PDF) in  $X_1$  and  $X_2$ , with expectation and covariance matrix

<span id="page-110-1"></span>
$$
\left[\begin{array}{c} x_1 \\ x_2 \end{array}\right], \quad \left[\begin{array}{cc} u^2(x_1) & ru(x_1)u(x_2) \\ ru(x_2)u(x_1) & u^2(x_2) \end{array}\right].
$$
 (16.3)

**16.2.5** Because the magnitudes of  $X_1$  and  $X_2$  in expression (16.2) are in practice small compared with unity, the resulting *Y* is close to unity. Results are accordingly expressed in terms of the quantity

<span id="page-110-2"></span>
$$
\delta Y = 1 - Y = X_1^2 + X_2^2,\tag{16.4}
$$

taken as the measurement model. For physical reasons,  $0 \le Y \le 1$ , and hence  $0 \le \delta Y \le 1$ .

<span id="page-110-0"></span>**16.2.6** The determination of an estimate  $\delta y$  of  $\delta Y$ , the associated standard uncertainty *u*(δ*y*) and a coverage interval for δ*Y* will be considered for choices of  $x_1, x_2, u(x_1), u(x_2)$ and  $r(x_1, x_2)$ . All quantities have dimension 1.

<span id="page-110-4"></span>16.2.3 Given respectively are estimates  $x_i$  cond  $x_j$  cof the quentities  $\int_{x_i}^{x_i} f(x) dx \leq \int_{x_i}^{x_i} f(x) dx \leq \int_{x_i}^{x_i}$ **16.2.7** Six cases are considered, in which  $x_2$  is taken as zero and  $u(x_1) = u(x_2) = 0.005$ . The first three cases correspond to  $x_1 = 0$ , 0.010 and 0.050, each with  $r = 0$ . The other three cases correspond to taking the same  $x_1$ , but with  $r = 0.9$ . The various numerical values of *x*<sup>1</sup> (comparable to those occurring in practice) are used to investigate the extent to which the results obtained using the considered approaches differ. In these six cases,  $x_2$ is set to zero. It would be equally applicable to set  $x_1$  to zero and use differing values for  $x_2$ in these examples. In most practical situations, it is expected that  $-0.05 \le x_1 \le 0.05$  and  $-0.05 \le x_2 \le 0.05.$ 

**16.2.8** For the cases in which  $r = 0$ , the covariance matrix given in formulæ [\(16.3\)](#page-110-1) reduces to a diagonal matrix with entries  $\left( u^2(x_1), u^2(x_2) \right)$  and the corresponding joint distribution for  $X_1$  and  $X_2$  to the product of two univariate normal distributions for  $X_i$ , for  $i = 1, 2$ , with expectation  $x_i$  and standard deviation  $u(x_i)$ .

### <span id="page-110-3"></span>**16.3 Propagation and summarizing:**  $r = 0$

#### **16.3.1 General**

**16.3.1.1** The evaluation of uncertainty is treated

1. analytically (for purposes of comparison),

- 2. using the GUF and
- 3. using the MCM.

NOTE These approaches do not constrain δ*Y* to be less than or equal to unity. However, for sufficiently small uncertainties  $u(x_1)$  and  $u(x_2)$ , as here, the PDF for  $\delta Y$  may adequately be approximated by a simpler PDF defined over all non-negative values of δ*Y* . A rigorous treatment, using Bayesian inference [[113](#page-134-0)], which applies regardless of the magnitudes of  $u(x_1)$  and  $u(x_2)$ , is possible, but beyond the scope here. Also see [[12,](#page-126-0) Scope, NOTE 2].

<span id="page-111-2"></span>**16.3.1.2**  $\delta y$  and  $u(\delta y)$  can generally be formed analytically as the expectation and standard deviation of δ*Y* , as characterized by the PDF for δ*Y* . See 16.5.2. The PDF for δ*Y* can be formed analytically when  $x_1 = 0$  and, in particular, used to determine the endpoints of the shortest 95 % coverage interval in that case. See 16.5.3.

<span id="page-111-3"></span>**16.3.1.3** GUF with first-order terms  $(GUF_1)$  and with higher-order terms  $(GUF_2)$  is applied for each of the three estimates  $x_1$  in the uncorrelated case. See 16.5.4.

**16.3.1.4** GUF<sub>1</sub> (see clause [7.3.1\)](#page-55-0) yields

$$
\delta y_{GUF1} = x_1^2 + x_2^2,
$$
\n
$$
u^2(\delta y_{GUF1}) = u^2(x_1^2) + u^2(x_2^2) = 4x_1^2 u^2(x_1) + 4x_2^2 u^2(x_2).
$$
\n(16.6)

**16.3.1.5** 
$$
GUF_2
$$
 (see clause 7.3.2) yields

$$
\delta y_{GUE2} = x_1^2 + x_2^2 + u^2(x_1) + u^2(x_2),\tag{16.7}
$$

$$
u^{2}(y_{GUF2}) = 4x_{1}^{2}u^{2}(x_{1}) + 4x_{2}^{2}u^{2}(x_{2}) + 2u^{4}(x_{1}) + 2u^{4}(x_{2}).
$$
\n(16.8)

Considering higher-order terms in the Taylor expansion of the measurement model impacts not only on the standard uncertainty, but on the estimate [[72](#page-131-1)]. This important fact is not taken into account in JCGM 100:2008 (see [[6,](#page-126-1) clause 5.1.2]).

**16.3.1.6** MCM is applied in each case with  $M = 10^6$  trials.

### <span id="page-111-1"></span>**16.3.2** Input estimate  $x_1 = 0$

<span id="page-111-4"></span>2. using the [G](#page-112-0)UI and<br>3. using Ne GUI and<br>3. using NeWAM.<br>
This phe MoM.<br>
S[t](#page-120-0)ript[e](#page-117-1) photon constants is to be less than or equal to unity they<br>exactly the differenty and therefore the most committed by a simple the<br>different **16.3.2.1** For the input estimate  $x_1 = 0$ , GUF<sub>2</sub> and MCM only can be used, because the partial derivatives of  $\delta Y$  with respect to  $X_1$  and  $X_2$ , evaluated at  $X_1 = x_1$  and  $X_2 = x_2$ , are identically zero when  $x_1 = x_2 = 0$ . Accordingly, first-order terms in the Taylor expansion would be identically zero. Thus, if GUF<sub>1</sub> were applied, both the estimate and the associated standard uncertainty would incorrectly be computed as zero!

NOTE A similar difficulty would arise for  $x_1$  close to zero.

<span id="page-111-0"></span>**16.3.2.2** Figure 16.1 shows the results for δ*Y* determined

- 1. analytically (the exponentially decreasing curve for  $\delta Y \ge 0$  and zero elsewhere),
- 2. using  $GUF_2$  (bell-shaped curve) and
- 3. using MCM (scaled frequency distribution shown as a histogram).

NOTE In all figures, normal PDFs are depicted for the GUF results. These PDFs are those that would hold for the measurand should the CLT be applicable. As such, they are to be understood as a mere graphical tool useful to assess visually the adequacy of the CLT approximation.

<span id="page-112-4"></span>

<span id="page-112-0"></span>Figure 16.1: Results for the model of comparison loss in the case  $x_1 = x_2 = 0$ , with  $u(x_1) = u(x_2) = 0.005$  and  $r = 0$  (16.3.2.2, 16.3.2.5, 16.3.2.8)

**16.3.2.3** Since the partial derivatives of the model function (16.4) of order higher than two are all identically zero, GUF $_2$  takes all Taylor-series terms, that is, the full non-linearity of the problem, into account. Thus, the GUF<sub>2</sub> solution is the best that is possible within the GUF.

**16.3.2.4** The analytic solution is a particular chi-squared distribution—the sum of squares of two standard normal variables (see [16.5.3\)](#page-118-0). It is seen in the figure that this PDF is very different from the normal PDF assumed when using GUF<sub>2</sub>. It can therefore be concluded that the CLT approximation is inadequate and, as a consequence, it can be expected that, although the estimate and associated standard uncertainty provided by GUF<sub>2</sub> are correct, coverage intervals are not.

<span id="page-112-1"></span>**16.3.2.5** It is also seen in figure 16.1 that the PDF provided by MCM is consistent with the analytic solution.

<span id="page-112-3"></span>**16.3.2.6** The estimates  $\delta y$  determined as the expectation of  $\delta Y$  obtained analytically (A), using GUF<sub>1</sub> (G<sub>1</sub>), GUF<sub>2</sub> (G<sub>2</sub>) and MCM (M) are given in columns 2 to 5 of the row corresponding to  $x_1 = 0.000$  in table [16.1.](#page-113-0) Columns 6 to 9 contain the corresponding  $u(\delta y)$ .

**16.3.2.7** Use of GUF<sub>1</sub> with model [\(16.4\)](#page-110-2) is not possible, as explained in subclause [16.3.2.1.](#page-111-1) Both GUF<sub>2</sub> and MCM correctly provide estimate and associated standard uncertainty in agreement with the analytic solution.

<span id="page-112-2"></span>**16.3.2.8** Figure 16.1 also shows the shortest 95 % coverage intervals for the corresponding approximations to the distribution function for δ*Y* . The 95 % coverage interval, indicated by dotted vertical lines, as provided by GUF $_2$  is infeasible: it erroneously implies





there is a non-zero probability that δ*Y* is negative. However, it provides a reasonably accurate upper endpoint. The continuous vertical lines are the endpoints of the shortest 95 % coverage interval derived from the analytic solution, as described in [16.5.3.](#page-118-0) The endpoints of the shortest 95 % coverage interval determined using MCM are indistinguishable to graphical accuracy from those for the analytic solution.

**16.3.2.9** The endpoints of the shortest coverage intervals for  $x_1 = 0.000$  in table [16.1](#page-113-0) are given in columns 10 to 13 of that table.

### **16.3.3** Input estimate  $x_1 = 0.010$

<span id="page-113-2"></span>**16.3.3.1** For the input estimate  $x_1 = 0.010$ , with correlation coefficient  $r = 0$ , figure [16.2](#page-113-1) shows the results obtained using GUF $_1$  (higher-peaked curve), GUF $_2$  (lower-peaked curve) and MCM (histogram).

<span id="page-113-3"></span><span id="page-113-0"></span>

<span id="page-113-1"></span>Figure 16.2: As figure [16.1](#page-112-0) except that  $x_1 = 0.010$ , and the results from GUF<sub>1</sub> (higherpeaked curve) and GUF $_2$  (lower-peaked curve) [\(16.3.3.1,](#page-113-2) [16.3.3.3,](#page-114-0) [16.3.4.1,](#page-114-1) [16.4.3\)](#page-116-0)

**16.3.3.2** The PDF provided by MCM exhibits a modest left-hand flank, although it is truncated at zero, the smallest feasible value of δ*Y* . Further, compared with the results for  $x_1 = 0$ , it is closer in form to a normal PDF. This outcome is due to the fact that firstorder terms in the Taylor expansion are no longer identically zero as for  $x_1 = 0$ , so that the CLT approximation starts becoming acceptable. The expectations for δ*Y* provided by GUF<sup>1</sup> and GUF<sub>2</sub> are 100 × 10<sup>-6</sup> and 150 × 10<sup>-6</sup>, and the standard uncertainties  $100 \times 10^{-6}$ and 112 × 10−<sup>6</sup> , respectively. This closeness of the standard uncertainties indicates that contributions from second-order terms are smaller than those from first order.

<span id="page-114-0"></span>**16.3.3.3** Figure [16.2](#page-113-1) also shows the endpoints of the shortest 95 % coverage intervals obtained by the three approaches. The continuous vertical lines denote the endpoints of the interval provided by MCM, the broken vertical lines those resulting from  $GUF<sub>1</sub>$  and the dotted vertical lines from GUF<sub>2</sub>. The interval provided by GUF<sub>1</sub> is shifted to the left compared with that from GUF<sub>2</sub> and with the shortest 95 % coverage interval for MCM. Both intervals provided by GUF again include infeasible values of  $\delta Y$ . Again, GUF<sub>2</sub> provides an accurate upper limit, graphically indistinguishable from that provided by MCM. The interval provided by MCM has its left-hand endpoint at zero, the smallest feasible value.

**16.3.3.4** The corresponding results are given in the penultimate row of table [16.1.](#page-113-0)

### **16.3.4 Input estimate**  $x_1 = 0.050$

<span id="page-114-1"></span>**16.3.4.1** Figure [16.3](#page-114-2) is similar to figure 16.2, but for  $x_1 = 0.050$ . Now, the PDF provided by MCM is much closer to a normal PDF, a clear indication that the CLT applies in this case. The reasons are that the first-order term in the Taylor expansion is dominant and the input quantities are assigned normal distributions. For the same reasons, the PDFs assumed with both variants of GUF are virtually indistinguishable from each other, – apart from a slight difference in the expectations – and from the PDF provided by MCM. The latter exhibits a slight skewness, as evidenced in the tail regions. The coverage intervals provided by the two variants of GUF are now feasible, that provided by  $GUF<sub>2</sub>$  being slightly closer to the MCM solution, as expected.

<span id="page-114-3"></span>

<span id="page-114-2"></span>Figure 16.3: As figure [16.2](#page-113-1) except that  $x_1 = 0.050$  [\(16.3.4.1,](#page-114-1) [16.4.3\)](#page-116-0)

**16.3.4.2** The corresponding results are given in the final row of table [16.1.](#page-113-0)

## **16.3.5 Discussion**

As  $x_1$  becomes increasingly removed from zero, the results given by GUF, with first-order and with higher-order terms, and those for MCM become closer to each other.

The numerical values  $x_1 = x_2 = 0$  lie in the centre of the region of interest to the practitioner, corresponding to the so-called "matched" condition for the power meter being calibrated, and thus in no sense constitute an extreme case.

Because of the symmetry of the model in  $X_1$  and  $X_2$ , exactly the same effect would occur were  $x_2$  used in place of  $x_1$ .

<span id="page-115-1"></span>**16.3.5 Discussion**<br> **16.3.5 Discussion**<br> **37.5** Sylecomes interceasingly termoved from zero, the results given by [G](#page-126-0)UP, with findebagin<br>
and wish higher-order terms, and those for MCM become close tro each other<br>
the nume One reason why  $GUF_1$  might be used in practice is that software for its implementation is readily available: results obtained from it might sometimes be accepted without question. For the case where  $x_1 = x_2 = 0$  (figure 16.1), the danger would be apparent because the standard uncertainty *u*(δ*y*) was computed as zero, and consequently any coverage interval for  $\delta Y$  would be of zero length for any coverage probability. For  $x_1 \neq 0$  (or  $x_2 \neq 0$ ),  $u(\delta y)$ and the length of the coverage interval for δ*Y* are both non-zero, so no such warning would be available without prior knowledge of likely values for *u*(δ*y*) and this length. Thus, a danger in implementing software based on GUF for these calculations is that checks of the software for  $x_1$  or  $x_2$  sufficiently far from zero would not indicate such problems, although, when used subsequently in practice for small values of  $x_1$  or  $x_2$ , the results would be invalid, but conceivably unwittingly accepted.

## **16.4 Propagation and summarizing:**  $r \neq 0$

### **16.4.1 General**

<span id="page-115-0"></span>**16.4.1.1** The three approaches used in the cases where the  $X_i$  are uncorrelated (see [16.3\)](#page-110-3) are now applied for the three cases in which they are correlated, using, as an example,  $r = 0.9$ . GUF<sub>1</sub> only is used. Unlike the cases where the  $X_i$  are uncorrelated, GUF<sub>2</sub> is not applied, no counterpart being provided in JCGM 100:2008 for the formula containing higher-order terms when the  $x_i$  have associated non-zero correlation coefficient [[12,](#page-126-0) clause 5.8]. Other aspects match those in 16.3.

NOTE Calculations with higher-order terms are possible even in the case of correlated *x<sup>i</sup>* [[72](#page-131-1)].

**16.4.1.2** For GUF<sub>1</sub>,  $u(\delta y)$  is evaluated as described in [16.5.5.](#page-120-1) Expression [\(16.17\)](#page-120-2) in that subclause gives, for  $x_2 = 0$ ,

 $u^2(\delta y) = 4x_1^2$  $u^2(u^2(x_1))$ .

Consequently,  $u(\delta y)$  does not depend on *r* and GUF with first-order terms gives identical results to those presented in [16.3.](#page-110-3) In particular, for the case  $x_1 = 0$ ,  $u(\delta y)$  is (incorrectly) computed as zero, as in [16.3.2.1.](#page-111-1)

**16.4.1.3** MCM was implemented by sampling randomly from *X* characterized by a bivariate normal PDF with the given expectation and covariance matrix (expressions [\(16.3\)](#page-110-1)). The procedure in [12, Annex C.5] was used.

NOTE Apart from the requirement to draw from a multivariate distribution, the implementation of MCM for input quantities that are correlated is no more complicated than when the input quantities are uncorrelated.

### **16.4.2** Input estimates  $x_1 = 0$ , 0.010 and 0.050

<span id="page-116-2"></span>**16.4.2.1** Table [16.2](#page-116-1) contains the results obtained. Those from MCM indicate that although  $\delta y$  is unaffected by the correlation between the  $X_i$ ,  $u(\delta y)$  is so influenced, more so for small *x*<sup>1</sup> , compared with the equivalent values in Table [16.1.](#page-113-0) The 95 % coverage intervals are influenced accordingly.

<span id="page-116-1"></span>Table 16.2: Comparison loss results for input estimates with associated non-zero correlation coefficient ( $r = 0.9$ ) obtained analytically (A), and using  $GUF_1$  ( $G_1$ ) and MCM (M) [\(16.4.2.1\)](#page-116-2)

$x_1$	Estimate $\delta y / 10^{-6}$			Std. uncertainty $u(\delta y)/10^{-6}$			Shortest 95 % coverage interval for $\delta Y / 10^{-6}$		
		G1	M	A	$G_1$	M	A	G <sub>1</sub>	М
0.000	50	$\Omega$	50	67		67		[0, 0]	[0, 185]
0.010	150	100		150 121	100	121		$[-96, 296]$	[13, 398]
0.050	2550	2500	2551	505	- 500	504		[1520, 3480]	[1628, 3555]

<span id="page-116-4"></span>**16.4.2.2** Figures [16.4](#page-116-3) and [16.5](#page-117-2) show the results provided by  $GUF_1$  (bell-shaped curves) and MCM (scaled frequency distributions) in the cases  $x_1 = 0.010$  and  $x_1 = 0.050$ , respectively. The endpoints of the shortest 95 % coverage interval provided by the two approaches are also shown, as dotted vertical lines for GUF and continuous vertical lines for MCM.

NOTE Strictly, the conditions under which  $\delta Y$  can be characterized by a normal PDF do not hold following an application of GUF in this circumstance (see [12, clause 5.8] and [[6,](#page-126-1) Annex G.6.6]). However, this PDF and the endpoints of the corresponding 95 % coverage interval are shown because such a characterization is commonly used.

<span id="page-116-5"></span>

<span id="page-116-3"></span>Figure 16.4: Results for the model of comparison loss in the case  $x_1 = 0.010$ ,  $x_2 = 0$ , with  $u(x_1) = u(x_2) = 0.005$  and  $r = 0.9$  [\(16.4.2.2,](#page-116-4) [16.4.3\)](#page-116-0)

## <span id="page-116-0"></span>**16.4.3 Discussion**

In the case  $x_1 = 0.010$  (figure [16.4\)](#page-116-3), the effect of the non-zero correlation coefficient has been to change noticeably the results returned by MCM (compare with figure [16.2\)](#page-113-1). Not

<span id="page-117-3"></span>

<span id="page-117-2"></span>Figure 16.5: As figure [16.4](#page-116-3) except that  $x_1 = 0.050$  (16.4.2.2[,16.4.3\)](#page-116-0)

only has the shape of (the approximation to) the PDF changed, but the corresponding coverage interval no longer has its left-hand endpoint at zero. In the case  $x_1 = 0.050$ (figure [16.5\)](#page-117-2), the differences between the results for the cases where the input quantities are uncorrelated and correlated (compare with figure 16.3) are less obvious.

#### <span id="page-117-0"></span>**16.5 Analytic calculations**

#### **16.5.1 General**

This clause provides some analytic results, partly based on the considerations in [[30](#page-128-0)]. Subclause [16.5.2](#page-117-1) provides expectation and standard deviation of δ*Y* [\(16.3.1.2\)](#page-111-2). Sub-clause [16.5.3](#page-118-0) provides the PDF for  $\delta Y$  when  $x_1 = x_2 = r(x_1, x_2) = 0$  [\(16.3.1.2\)](#page-111-2). Subclause [16.5.4](#page-120-0) applies GUF for uncorrelated input quantities [\(16.3.1.3\)](#page-111-3) and correlated input quantities [\(16.4.1.1\)](#page-115-0).

#### <span id="page-117-1"></span>**16.5.2 Expectation and standard deviation**

**16.5.2.1** The variance of a random variable *X* can be expressed in terms of expectations as [[97,](#page-133-1) page 124]

$$
V(X) = E(X^2) - [E(X)]^2.
$$

Thus,

$$
E(X^{2}) = [E(X)]^{2} + V(X) = x^{2} + u^{2}(x),
$$

where *x* is the estimate of *X* and  $u(x)$  the associated standard uncertainty.

For model (16.4), *i.e.*  $\delta Y = X_1^2 + X_2^2$ ,

$$
\delta y = E(\delta Y) = x_1^2 + x_2^2 + u^2(x_1) + u^2(x_2).
$$

This result applies

1. regardless of the PDFs assigned to  $X_1$  and  $X_2$ , and

2. whether  $X_1$  and  $X_2$  are independent or not.

The standard uncertainty associated with δ*y* can be obtained from

<span id="page-118-1"></span>
$$
u^{2}(\delta y) = u^{2}(x_{1}^{2}) + u^{2}(x_{2}^{2}) + 2u(x_{1}^{2}, x_{2}^{2}),
$$

where, for  $i = 1$  and  $i = 2$ ,  $u^2(x_i^2) = V(X_i^2)$ , and  $u(x_1^2, x_2^2) = Cov(X_1^2, X_2^2)$ . Then, applying Price's Theorem for normal distributions [[86,](#page-132-1) [92](#page-133-2)],

$$
u^{2}(\delta y) = 4u^{2}(x_{1})x_{1}^{2} + 4u^{2}(x_{2})x_{2}^{2} + 2u^{4}(x_{1}) + 2u^{4}(x_{2}) + 4u^{2}(x_{1}, x_{2}) + 8u(x_{1}, x_{2})x_{1}x_{2}.
$$
 (16.9)

When  $x_2 = 0$  and  $u(x_2) = u(x_1)$ , and replacing  $u(x_1, x_2)$  by  $r(x_1, x_2)u^2(x_1)$ ,

$$
u(\delta y) = 2\left\{x_1^2 + [1 + r^2(x_1, x_2)]u^2(x_1)\right\}^{1/2}u(x_1).
$$

**16.5.2.2** When  $X_1$  and  $X_2$  are uncorrelated, that is,  $u(x_1, x_2) = 0$ , expression [\(16.9\)](#page-118-1) becomes

<span id="page-118-2"></span>
$$
u^{2}(\delta y) = 4u^{2}(x_{1})x_{1}^{2} + 4u^{2}(x_{2})x_{2}^{2} + 2u^{4}(x_{1}) + 2u^{4}(x_{2}).
$$
\n(16.10)

Expression [\(16.10\)](#page-118-2) can be verified by applying formula (10) in [6, clause 5.1.2] and the immediately following note.

### <span id="page-118-0"></span>**16.5.3 Analytic solution for zero estimate of the voltage reflection coefficient having associated zero covariance**

<span id="page-118-3"></span>2. whether X<sub>1</sub> and X<sub>2</sub> are independent or not.<br>
The standard uncertainty associated with Sy can be obtained from<br>  $u^2(\xi y) = u^2(x_1^2) + u^2(x_2^2) + 2u(x_1^2, x_2^2)$ , and  $dx_1^2, x_2^2) = \text{Cov}(X_1^2, \frac{\pi \xi}{2})$ . Then, applying<br> **16.5.3.1** For the case  $x_1 = x_2 = r(x_1, x_2) = 0$  and  $u(x_1) = u(x_2)$ , the PDF  $g_Y(\eta)$  for *Y* can be obtained analytically. It is valuable to have such a solution for further validation purposes. In the above circumstances,

$$
\delta Y = u^2(x_1) \left[ \frac{X_1^2}{u^2(x_1)} + \frac{X_2^2}{u^2(x_2)} \right].
$$

**16.5.3.2** The term in square brackets is the sum, *Z*, say, of the squares of two independent quantities, each of which is distributed as a standard normal PDF. Thus, the sum is distributed as chi-squared with two degrees of freedom [[97,](#page-133-1) page 177], so that

$$
\delta Y = u^2(x_1)Z,
$$

where *Z* has PDF

$$
g_{z}(\zeta) = \chi_{2}^{2}(\zeta) = e^{-\zeta/2}/2.
$$

**16.5.3.3** The application of a general formula [[97,](#page-133-1) pages 57–61] for the PDF  $g_Y(\eta)$  of a differentiable and strictly decreasing function of a variable (here *Z*) with a specified PDF yields

$$
g_{\nu}(\eta) = \frac{1}{u^2(x_1)} \chi_2^2\left(\frac{\eta}{u^2(x_1)}\right) = \frac{1}{2u^2(x_1)} \exp\left(-\frac{\eta}{2u^2(x_1)}\right), \qquad \eta \ge 0.
$$

**16.5.3.4** The expectation of  $\delta Y$  is

$$
\delta y = E(\delta Y) = \int_0^\infty \eta g_Y(\eta) d\eta = 2u^2(x_1)
$$

and the variance

$$
u^2(\delta y) = V(\delta Y) = \int_0^\infty (\eta - \delta y)^2 g_Y(\eta) d\eta = 4u^4(x_1),
$$

that is, the standard deviation is  $2u^2(x_1)$ , results that are consistent with those in [16.5.2.](#page-117-1)

**16.5.3.5** By integration, the corresponding distribution function is

<span id="page-119-0"></span>
$$
G_{\mathcal{Y}}(\eta) = 1 - \exp\left(-\frac{\eta}{2u^2(x_1)}\right), \qquad \eta \ge 0. \tag{16.11}
$$

**16.5.3.6** Let  $\delta y_\alpha$  be that  $\eta$  in expression (16.11) corresponding to  $G_\gamma(\eta) = \alpha$  for any  $\alpha$ satisfying  $0 \le \alpha \le 1 - p$ . Then

 $δy<sub>α</sub> = -2u<sup>2</sup>(x<sub>1</sub>) \ln(1 - α)$ 

and a 100*p* % coverage interval for δ*Y* [12, clause 7.7] is

$$
[\delta y_{\alpha}, \ \delta y_{p+\alpha}] \equiv -2u^2(x_1)[\ln(1-\alpha), \ \ln(1-p-\alpha)] \tag{16.12}
$$

with length

$$
H(\alpha) = -2u^2(x_1)\ln\left(1 - \frac{p}{1-\alpha}\right).
$$

<span id="page-119-1"></span>**16.5.3.4** The expectation of SY is<br>  $\Rightarrow y = E(bY) = \int_0^\infty \eta g_0(y) d\eta = 2a^2(x_1)$ <br>
and the variance<br>  $u^2(5y) = V(5Y) = \int_0^\infty (\eta - 5y)^2 g_0(y) d\eta = 4a^2(x_1)$ ,<br>
that is, the standard deviation is  $2a^2(x_1)$ , results that are consistent **16.5.3.7** The shortest 100*p* % coverage interval is given by determining  $\alpha$  to minimize *H*( $\alpha$ ) [[12,](#page-126-0) clause 5.3.4]). Since *H*( $\alpha$ ) is a strictly increasing function of  $\alpha$  for  $0 \le \alpha \le 1 - p$ , *H*(*α*) is minimized when  $α = 0$ . Thus, the shortest 100*p* % coverage interval for δ*Y* is

 $[0, -2u^2(x_1)\ln(1-p)].$ 

For  $u(x_1) = 0.005$ , the shortest 95 % coverage interval is

[0, 0.000 149 8].

**16.5.3.8** The 95 % probabilistically symmetric coverage interval for δ*Y* is given by setting  $\alpha = (1-p)/2$  [12, clause 5.3.3]:

$$
-2u^2(x_1)[\ln 0.975, \ln 0.025] = [0.000 001 3, 0.000 184 4],
$$

which is 20 % longer than the shortest 95 % coverage interval.

NOTE The above analysis is indicative of an approach that can be applied to some problems of this type. In this case, the results could have been obtained more directly since  $g_{\rm\scriptscriptstyle Y}(\eta)$  is strictly increasing and unimodal, and so the shortest coverage interval is in the region of highest density.

#### <span id="page-120-0"></span>**16.5.4 GUF — Uncorrelated input quantities**

**16.5.4.1** In model [16.4,](#page-110-2)  $X_1$  and  $X_2$  are assigned normal PDFs having expectations  $x_1$ and  $x_2$  and variances  $u^2(x_1)$  and  $u^2(x_2)$ , respectively. The corresponding estimate of  $\delta Y$ is thus

<span id="page-120-3"></span>
$$
\delta y = x_1^2 + x_2^2. \tag{16.13}
$$

The only non-trivially non-zero partial derivatives of the model are, for *i* 

$$
\frac{\partial f}{\partial X_i} = 2X_i, \qquad \frac{\partial^2 f}{\partial X_i^2} = 2.
$$

**16.5.4.2** Hence the application of [6, clause 5.1.2] gives, for the standard uncertainty *u*(δ*y*),

<span id="page-120-4"></span>
$$
u^{2}(\delta y) = \left[ \left( \frac{\partial f}{\partial X_{1}} \right)^{2} u^{2}(x_{1}) + \left( \frac{\partial f}{\partial X_{2}} \right)^{2} u^{2}(x_{2}) \right] \Big|_{X=x} = 4x_{1}^{2}u^{2}(x_{1}) + 4x_{2}^{2}u^{2}(x_{2}), \tag{16.14}
$$

based on a first-order Taylor series approximation of  $f(X)$ . If the non-linearity of f is significant  $[6, 5.1.2$  $[6, 5.1.2$  $[6, 5.1.2$  note], the terms

$$
\frac{1}{2} \left( \frac{\partial^2 f}{\partial X_1^2} \right) u^2(x_1) + \frac{1}{2} \left( \frac{\partial^2 f}{\partial X_2^2} \right) u^2(x_2)
$$

and

$$
\frac{1}{2} \left( \frac{\partial^2 f}{\partial X_1^2} \right)^2 u^4(x_1) + \frac{1}{2} \left( \frac{\partial^2 f}{\partial X_2^2} \right)^2 u^4(x_2)
$$

need to be appended to formulæ (16.13) and (16.14), respectively. These formulæ thus become

$$
\delta y = E(\delta Y) = x_1^2 + x_2^2 + u^2(x_1) + u^2(x_2)
$$
\n(16.15)

and

$$
u^{2}(\delta y) = 4x_{1}^{2}u^{2}(x_{1}) + 4x_{2}^{2}u^{2}(x_{2}) + 2u^{4}(x_{1}) + 2u^{4}(x_{2}).
$$
\n(16.16)

**16.5.4.3** A 95% coverage interval for  $\delta Y$  is given by

δ*y* ± 2*u*(δ*y*),

as a consequence of δ*Y* having a normal PDF.

#### <span id="page-120-1"></span>**16.5.5 GUF — Correlated input quantities**

**16.5.5.1** When the input quantities are correlated, the covariance matrix associated with the estimates of the input quantities is given in formulæ [\(16.3\)](#page-110-1).

<span id="page-120-2"></span>**16.5.5.2** The application of JCGM 100:2008 clause 5.2.2 gives

<span id="page-120-5"></span>16.5.4 GUF — Uncorrelated input quantities  
\n16.5.4.1 In model 16.4, 
$$
X_1
$$
 and  $X_2$  are assigned normal PDFs having expectations  $x_1$  and  $x_2$  and variances  $u^2(x_1)$  and  $u^2(x_2)$ , respectively. The corresponding estimate of  $\delta Y$  is thus  
\n
$$
\delta y = x_1^2 + x_2^2.
$$
\n
$$
\frac{\partial f}{\partial X_i} = 2X_i, \qquad \frac{\partial^2 f}{\partial X_i^2} = 2.
$$
\n16.5.4.2 Hence the application of [6, clause 5.1.2] gives, for the standard uncertainty  $u(\delta y)$ ,  
\n
$$
u^2(\delta y) = \left[ \left( \frac{\partial f}{\partial X_1} \right)^2 u^2(x_1) + \left( \frac{\partial f}{\partial X_2} \right)^2 u^2(x_2) \right] \Big|_{X=x} = 4x_1^2 u^2(x_1) + 4x_2^2 u^2(x_2)
$$
, (16.14)  
\nbased on a first-order Taylor series approximation of  $f(X)$ . If the non-linearity of  $f$  is  
\nsignificant [6, 5.1.2 note], the terms  
\n
$$
\frac{1}{2} \left( \frac{\partial^2 f}{\partial X_1^2} \right) u^2(x_1) + \frac{1}{2} \left( \frac{\partial^2 f}{\partial X_2^2} \right) u^2(x_2)
$$
\nand  
\n
$$
\frac{1}{2} \left( \frac{\partial^2 f}{\partial X_1^2} \right)^2 u^4(x_1) + \frac{1}{2} \left( \frac{\partial^2 f}{\partial X_2^2} \right)^2 u^4(x_2)
$$
\nneed to formulae (16.13) and (16.14), respectively. These formula thus  
\nbecome  
\n
$$
\delta y = E(\delta Y) = x_1^2 + x_2^2 + u^2(x_1) + u^2(x_2)
$$
\n(16.15)  
\nand  
\n
$$
u^2(\delta y) = 4x_1^2 u^2(x_1) + 4x_2^2 u^2(x_2) + 2u^4(x_1) + 2u^4(x_2).
$$
\n(16.16)  
\n16.5.4.3 A 9.5% coverage interval for  $\delta Y$ 

## **16.6 Conclusion**

16.6 Conclusion<b[r](#page-126-0)>
This example shows that in microwere power meter calibration the JCGM 100-2008 undertainty immework (GUP), wither with (GUP) or without (GUP) the inclusion of higher-<br>retrainty immework (GUP), wither with This example shows that in microwave power meter calibration the JCGM 100:2008 uncertainty framework (GUF), either with (GUF<sub>2</sub>) or without (GUF<sub>1</sub>) the inclusion of higherorder terms, can yield unrealistic estimates, standard uncertainties and coverage intervals, as demonstrated by the results presented in the figures and tables in this example. It is therefore recommended that, for situations such as these, a Monte Carlo method [12] is used to evaluate a coverage interval.

# **A Overview of the parts of the GUM**

Table A.1: Overview of the parts (published and planned) of the Guide to the expression of uncertainty in measurement

GUM part	Old reference	Title				
<b>JCGM GUM-1</b>	$JCGM$ 104 <sup>a</sup>	Guide to the expression of uncertainty in measurement - Part 1: Introduction				
JCGM GUM-2 <sup>b</sup>		Guide to the expression of uncertainty in measurement - Part 2: Concepts				
JCGM GUM-3 <sup>c</sup>	<b>JCGM 100</b>	Guide to the expression of uncertainty in measurement - Part 3: GUM:1995 with minor corrections				
JCGM GUM-4 <sup>d</sup>	<b>JCGM 106</b>	Guide to the expression of uncertainty in measurement - Part 4: The role of measurement uncertainty in confor- mity assessment Guide to the expression of uncertainty in measurement - Part 5: Examples Guide to the expression of uncertainty in measurement - Part 6: Developing and using measurement models				
JCGM GUM-5 <sup>b</sup>	$JCGM$ 110 <sup>h</sup>					
JCGM GUM-6 <sup>e</sup>	JCGM $103h$					
JCGM GUM-7 <sup>f</sup>	<b>JCGM 101</b>	Guide to the expression of uncertainty in measurement - Part 7: Propagation of distributions using a Monte Carlo method Guide to the expression of uncertainty in measurement - Part 8: Extension to any number of output quantities				
JCGM GUM-8 <sup>8</sup>	<b>JCGM 102</b>					
JCGM GUM-9 <sup>b</sup>		Guide to the expression of uncertainty in measurement - Part 9: Statistical models and data analysis for interlabo- ratory studies				
JCGM GUM-10 <sup>b</sup>		Guide to the expression of uncertainty in measurement - Part 10: Applications of the least squares method				
JCGM GUM-11 <sup>b</sup>		Guide to the expression of uncertainty in measurement - Part 11: Bayesian methods				
JCGM GUM-12 <sup>b</sup>		Guide to the expression of uncertainty in measurement - Part 12: Basic methods for uncertainty propagation				
<sup>a</sup> Replaced by JCGM GUM-1:2023 [9] <b>b</b> Planned		e Available as JCGM GUM-6:2020 [10] Available as JCGM 101:2008 [12] <sup>8</sup> Available as JCGM 102:2011 [7]				
$c$ Available as JCGM 100:2008 [6] $d$ Available as JCGM 106:2012 [8]		<sup>h</sup> Reference of the circulated committee draft				

- $c$  Available as JCGM 100:2008 [6]
- <sup>g</sup> Available as JCGM 102:2011 [[7](#page-126-6)]
- $d$  Available as JCGM 106:2012 [8]
- h Reference of the circulated committee draft

## **B Conventions and notation**

**B.1** A generic quantity is denoted by an upper case letter and a measured value (estimate) by a lower case letter [[12,](#page-126-0) Clause 4.8].

EXAMPLE Measured values of the measurand *Y* and of the generic input quantity  $X_i$  are denoted by  $y$  and  $x_i$ , respectively.

**B.2** A specific quantity is denoted by its agreed symbol. The measured value (estimate) is denoted by the same symbol adorned with a hat [[12,](#page-126-0) Clause 4.8].

EXAMPLE 1 The thermodynamic temperature is denoted by *T* and a measured value of this temperature by  $\hat{T}$ , for example  $\hat{T} = 295.4$  K.

EXAMPLE 2 The mass of a weight is denoted by  $m$  and the measured value by  $\widehat{m}$ .

NOTE 1 In JCGM 106:2012, a measured value is denoted by a subscript 'm'.

NOTE 2 ISO/IEC 80000 and the IUPAC Green Book [20] contain information about symbols for specific physical quantities.

**B.3** The state of knowledge about a quantity is described by the probability density function (PDF) of a random variable denoted by the same symbol used for the quantity [[12,](#page-126-0) Clause 4.1].

**B.4** The state of knowledge about an input quantity  $X$  is described by a PDF  $g_X(\xi)$ , where the subscript *X* denotes a random variable with possible values *ξ* [12, Clause 4.3].

**B.5** Similarly, the state of knowledge about the measurand  $Y$  is described by a PDF  $g_Y(\eta)$ and the corresponding probability distribution by  $G_Y(\eta).$  *Y* denotes a random variable with possible values *η* [12, Clause 4.7].

**B.6** A PDF for more than one input quantity is often called joint even if all input quantities are mutually independent.

**B.7** The PDF provided by the JCGM 100:2008 uncertainty framework (GUF) is either a normal distribution with mean  $\gamma$  and standard deviation  $u(\gamma)$  or a shifted and scaled Student *t*-distribution with mean *y*, scale *u*( *y*) and effective degrees of freedom *ν*.

**31** Conventions and notation<br>
131 A agencic quantity is denoted by an upper case letter and a measured value (certaintex)<br>
31 a lower case letter [12, Clause 4.8].<br>
EXMIVE Measured value of the neassurand Y and of the ge **B.8** The law of propagation of uncertainty (LPU) applies to the use of a first-order Taylor series approximation to the model. The term is qualified accordingly when a higher-order approximation is used [[12,](#page-126-0) Clause 4.9].

**B.9** The notation  $u(y)$  is to be read as the standard uncertainty associated with *y*. It represents the standard uncertainty about the true value of *Y* .

**B.10** Noting that a covariance matrix *V* is symmetric, only the upper triangular part is shown. The same applies for the correlation matrix *R*.

**B.11** The terms "coverage interval" and "coverage probability" are used throughout this document. JCGM 100:2008 uses the term "level of confidence" as a synonym for coverage probability, drawing a distinction between "level of confidence" and "confidence level" [6, 6.2.2], because the latter has a specific definition in statistics. Since, in some languages, the translation from English of these two terms yields the same expression, the use of these terms is avoided here [[12,](#page-126-0) Clause 4.11].

**B.12** The subscript "c" [[6,](#page-126-1) Clause 5.1.1] in  $u_c(y)$  for the combined standard uncertainty is deemed redundant and not used. Thus, the standard uncertainty associated with *y* is denoted generally as  $u(y)$ . Likewise, the qualifier "combined" is also regarded as superfluous [[12,](#page-126-0) Clause 4.10].

**B.13** Unless otherwise qualified, values are expressed in a manner that indicates the number of meaningful significant decimal digits [12, Clause 4.13].

**B.14** Some symbols have more than one meaning in this document. See Annex [C.](#page-125-0) The context clarifies the usage.

**B.15** According to Resolution 10 of the 22nd CGPM (2003) "... the symbol for the decimal marker shall be either the point on the line or the comma on the line . . . ". The Joint Committee for Guides in Metrology (JCGM) has decided to adopt, in its documents in English, the point on the line as decimal marker [12, Clause 4.12].

B.11 The ter[m](#page-126-0)s "coverage intertwill" and "coverage probability" are used throughout this contract the production in the definition in statistics. Since, in some foregroups that the production is a specific diffultion in s

## <span id="page-125-0"></span>**C Glossary of symbols**

### **C.1 Generic variables and functions**

- *C* sensitivity matrix
- Cov covariance
- *E* expectation
- *f* measurement model of the form  $Y = f(X)$
- *f* measurement model having multiple output quantities of the form  $Y = f(X)$
- *h* measurement model of the form  $h(Y, X) = 0$
- C Glossary of symbols<br>
C enterivations and functions<br>
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C exercision<br>
measurement model of the form  $V = f(X)$ <br>
measurement model of the form  $N(Y,X) = 0$ <br>
R measurement model of the form *h* measurement model having multiple output quantities of the form  $h(Y, X) = 0$
- *k* coverage factor
- *N* number of input quantities
- *U* expanded uncertainty
- *u* standard uncertainty
- *r* coefficient of linear correlation
- *R* correlation matrix
- *V* covariance matrix
- *V*, *σ* 2 variance
- *X* input quantity
- $X$  column vector of input quantities  $X_1, \ldots, X_N$
- *x* measured value (estimate) of input quantity *X*
- *Y* measurand (output quantity of a measurement model)
- *<i>Y* column vector of output quantities  $Y_1, \ldots, Y_N$
- *y* measured value (estimate) of measurand *Y*

## **C.2 Probability distributions**

- Cauchy folded scaled Cauchy distribution
- CTrap curvilinear trapezoidal distribution
- Gamma gamma distribution
- LN lognormal distribution
- N normal distribution
- R rectangular distribution (uniform distribution)
- U U-shaped distribution (arc sine distribution)
- *t* Student *t*-distribution

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## **Abbreviations**

- **AIC** Akaike's Information Criterion. 76–78
- **AICc** corrected Akaike's Information Criterion. 76, 77

**ANOVA** analysis of variance. 82, 83, 86

**BIC** Bayesian Information Criterion. 76–78

**BIPM** Bureau International des Poids et Mesures. vi

**CLT** Central Limit Theorem. 46, 102, 105, 107

**CRM** certified reference material. 3, 14, 15, 82

**CSIR** Council for Scientific and Industrial Research, South Africa. 97

**EPA** Environmental Protection Agency. 66

- **GHG** greenhouse gas. 32, 33, 37, 39, 40
- **GUF** JCGM 100:2008 uncertainty framework. iii, v, vii, 14, 15, 18, 19, 40, 42–49, 52, 54, 56, 57, 63, 99–102, 104, 105, 107–110, 114, 116
- **GUM** Guide to the expression of uncertainty in measurement. 1, 32, 39
- **IPCC** Intergovernmental Panel on Climate Change. 32, 34, 35
- **IUPAC** International Union of Pure and Applied Chemistry. 2, 11, 22, 27
- **JCGM** Joint Committee for Guides in Metrology. 32, 40, 117
- **LPU** law of propagation of uncertainty. 1, 2, 10, 16, 20, 21, 23, 25–27, 31, 34, 36, 39–41, 47, 58, 88, 94, 95, 99, 116
- **MCM** Monte Carlo method. iii, v, vii, 1, 10, 14, 15, 17–19, 40–47, 49–54, 56, 57, 79, 80, 99–102, 104–109
- **MCMC** Markov Chain Monte Carlo method. 62, 69, 85, 86
- **OLS** ordinary least squares. 59, 61, 63
- **Abbreviations**<br>
MC Akade's information Criterion. 76-78<br>
MC corrected Akathes' information Criterion. 76-77<br>
MCC corrected Akathes' information Criterion. 76-77<br>
MCC despesian Information Criterion. 76-78<br>
MC Bayesian In **PDF** probability density function. 19, 27, 30, 41–46, 48–51, 56, 57, 61–63, 68, 69, 71, 83, 85, 100, 103–111, 113, 116
- **SI** International System of Units. 2, 11
- **TSPM** total suspended particulate matter. 66–68, 71

**UNFCCC** United Nations Framework Convention on Climate Change. 32

**WLS** weighted least squares. 59, 61, 63, 64